

Exam given 1900 - 2200 hr, 30 October 2014

THIS IS A CLOSED BOOK EXAM. Show all steps to get full credit. Answer questions 1,2 in Book 1: questions 3, 4, in Book 2.

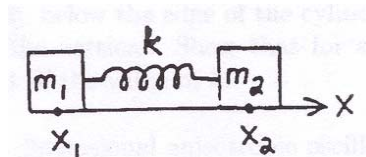
Book 1

1:(20pts) Consider a rocket fired vertically upwards from the ground in a uniform vertical gravitation field $g = 9.81m/s^2$. The rocket, with initial mass m_I , ejects uniformly α kg/sec of propellant at an exhaust velocity u m/sec.

- Derive the equation giving the time dependence of the vertical velocity of the rocket
- If the initial to final mass ratio of the rocket is $\frac{m_I}{m} = 10$, the exhaust velocity $u = 4500m/s$ and the fuel burns at a constant rate α for 300 seconds, calculate the maximum velocity of the rocket.
- If the rocket burns half the fuel in 150 seconds, then ejects the empty fuel tank of mass $0.05m_I$ and then burns the remaining half of the fuel in the second 150 seconds, calculate the maximum velocity of the rocket.

2;(25pts) Two masses m_1 and m_2 slide freely on a horizontal frictionless plane surface and are connected by a spring whose force constant is k and unstretched length is l .

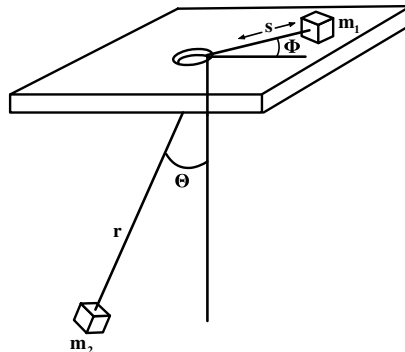
- Derive the equations of motion
- Derive the frequency of oscillation
- Show that the system oscillation can be represented by motion of a single equivalent mass.
- Which observables are conserved?



Book 2

3; (30pts) A mass m_1 sliding freely on a horizontal frictionless plane, is attached by a string of length b to a second mass m_2 that hangs through a hole in the frictionless plane. Assume that mass m_1 is rotating about the hole with an angular velocity $\frac{d\phi}{dt} = \omega$ at a distance s from the hole while mass m_2 is swinging in a vertical plane at an instantaneous angle θ to the vertical at a distance r from the hole. That is, $r + s = b$.

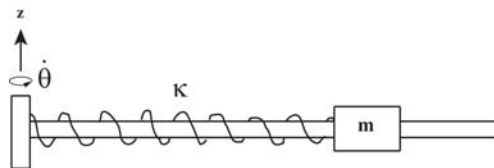
- Derive the Lagrangian for the system
- Use the Lagrangian to derive the equations of motion of the system for r, θ , and ϕ .
- Derive the Hamiltonian
- Determine all the constants of motion
- For the simple case where $\dot{\theta} = \theta = 0$, show that there is a distance r_0 for which the system is stable to external perturbations in r .
- Determine the angular frequency for oscillations about this stable position.



Mass m_2 , hanging from a rope that is connected to m_1 , which slides on a frictionless plane.

4; (25pts) A rigid straight, frictionless, massless, rod rotates about the z axis at an angular velocity $\dot{\theta}$. A mass m slides along the frictionless rod and is attached to the rod by a massless spring of spring constant κ and an unstretched length d

- Derive the Lagrangian and the Hamiltonian
- Derive the equations of motion in the stationary frame using Hamiltonian mechanics.
- What are the constants of motion?
- If the rotation is constrained to have a constant angular velocity $\dot{\theta} = \omega$ then is the non-cyclic Routhian $R_{noncyclic} = H - p_{\dot{\theta}}\dot{\theta}$ a constant of motion, and does it equal the total energy?
- Use the non-cyclic Routhian $R_{noncyclic}$ to derive the radial equation of motion in the rotating frame of reference for the cranked system with $\dot{\theta} = \omega$.



Useful formulae

Damped harmonic oscillator:

$$\frac{d^2x}{dt^2} + \Gamma \frac{dx}{dt} + \omega_0^2 x = 0$$

Solution for $\frac{\Gamma}{2} < \omega_0$ is

$$x = Ae^{-\frac{\Gamma t}{2}} \cos(\omega_1 t - \delta)$$

where

$$\omega_1^2 = \omega_0^2 - \frac{\Gamma^2}{4}$$

Sinusoidal-driven damped harmonic oscillator

$$\frac{d^2x}{dt^2} + \Gamma \frac{dx}{dt} + \omega_0^2 x = A \cos(\omega t)$$

Steady-state solution is

$$x_{ss} = \frac{A}{\sqrt{(\omega_0^2 - \omega^2)^2 + \Gamma^2 \omega^2}} \cos(\omega t - \delta)$$
$$\tan \delta = \frac{\Gamma \omega}{(\omega_0^2 - \omega^2)}$$

Newton's law of Gravitation

$$\vec{\mathbf{F}}_m = -G \frac{mM}{r^2} \hat{\mathbf{r}}$$
$$\vec{\nabla} \cdot \vec{\mathbf{g}} = -4\pi G \rho$$

Virial Theorem

$$\langle T \rangle = -\frac{1}{2} \left\langle \sum_i \vec{\mathbf{F}}_i \cdot \vec{\mathbf{r}}_i \right\rangle$$

Lagrange equations:

Lagrange equations including undetermined multipliers and generalized forces

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} - \frac{\partial L}{\partial q_j} = Q_j^{EXC} + \sum_{k=1}^m \lambda_k(t) \frac{\partial g_k}{\partial q_j} \quad (i = 1, 2, 3, \dots, s)$$

Generalized momentum

$$p_j = \frac{\partial L}{\partial \dot{q}_j}$$

Hamiltonian

$$H(q_i, p_i, t) = \sum_i p_i \dot{q}_i - L(q_i, \dot{q}_i, t)$$

Hamilton's equations of motion

$$\begin{aligned}\dot{q}_k &= \frac{\partial H}{\partial p_k} \\ \dot{p}_j &= -\frac{\partial H(\mathbf{q}, \mathbf{p}, t)}{\partial q_j} + \left[\sum_{k=1}^m \lambda_k \frac{\partial g_k}{\partial q_j} + Q_j^{EXC} \right] \\ \frac{dH(\mathbf{q}, \mathbf{p}, t)}{dt} &= \sum_j \left(\left[\sum_{k=1}^m \lambda_k \frac{\partial g_k}{\partial q_j} + Q_j^{EXC} \right] \dot{q}_j \right) - \frac{\partial L(\mathbf{q}, \dot{\mathbf{q}}, t)}{\partial t} \\ &= -\frac{\partial L}{\partial t} = \frac{\partial H}{\partial t}\end{aligned}$$

Routhian reduction

$$\begin{aligned}R_{cyclic}(q_1, \dots, q_n; \dot{q}_1, \dots, \dot{q}_s; p_{s+1}, \dots, p_n; t) &= \sum_{cyclic}^m p_i \dot{q}_i - L = H - \sum_{noncyclic}^s p_i \dot{q}_i \\ R_{noncyclic}(q_1, \dots, q_n; p_1, \dots, p_s; \dot{q}_{s+1}, \dots, \dot{q}_n; t) &= \sum_{noncyclic}^s p_i \dot{q}_i - L = H - \sum_{cyclic}^m p_i \dot{q}_i\end{aligned}$$

Cylindrical coordinates

$$\begin{aligned}L = T - U &= \frac{m}{2} (\dot{\rho}^2 + \rho^2 \dot{\phi}^2 + \dot{z}^2) - U(\rho, z, \phi) \\ H = T + U &= \frac{1}{2m} \left(p_\rho^2 + \frac{p_\phi^2}{\rho^2} + p_z^2 \right) + U(\rho, z, \phi)\end{aligned}$$

Spherical coordinates

$$\begin{aligned}L = T - U &= \frac{m}{2} (\dot{r}^2 + r^2 \dot{\theta}^2 + r^2 \sin^2 \theta \dot{\phi}^2) - U(r, \theta, \phi) \\ H = T + U &= \frac{1}{2m} \left(p_r^2 + \frac{p_\theta^2}{r^2} + \frac{p_\phi^2}{r^2 \sin^2 \theta} \right) + U(r, \theta, \phi)\end{aligned}$$

Vectors;

$$\begin{aligned}\vec{\mathbf{C}} = \vec{\mathbf{A}} \times \vec{\mathbf{B}} &= \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{vmatrix} \\ \vec{\mathbf{A}} \times \vec{\mathbf{A}} &= \vec{0} \\ \vec{\mathbf{A}} \cdot (\vec{\mathbf{A}} \times \vec{\mathbf{B}}) &= 0 \\ \vec{\mathbf{A}} \cdot (\vec{\mathbf{B}} \times \vec{\mathbf{C}}) &= (\vec{\mathbf{A}} \times \vec{\mathbf{B}}) \cdot \vec{\mathbf{C}} \\ \vec{\mathbf{A}} \times (\vec{\mathbf{B}} \times \vec{\mathbf{C}}) &= \vec{\mathbf{B}} (\vec{\mathbf{A}} \cdot \vec{\mathbf{C}}) - (\vec{\mathbf{A}} \cdot \vec{\mathbf{B}}) \vec{\mathbf{C}}\end{aligned}$$