

# VIRIAL THEOREM

I.

Following Clayton §2 and Hansen §1.3

The overall behaviour of a star is largely an outcome of the virial theorem - it is worth going over in detail.

The theorem is a statistical statement about a system of masses/particles that are mutually interacting.

Recall forces can be written as  $\vec{F}_i = \frac{d}{dt} \vec{p}_i$

The system is composed of mass points  $m_i$  at positions  $\vec{r}_i$  susceptible to forces  $\vec{F}_i$

Consider the quantity  $\sum_i \vec{p}_i \cdot \vec{r}_i$  summed over all particles/masses

then 
$$\frac{d}{dt} \sum_i \vec{p}_i \cdot \vec{r}_i = \sum_i \frac{d\vec{p}_i}{dt} \cdot \vec{r}_i + \sum_i \vec{p}_i \cdot \frac{d\vec{r}_i}{dt}$$

$$\begin{aligned} \uparrow \vec{p}_i &= m_i \dot{\vec{r}}_i \\ \frac{d\vec{r}_i}{dt} &= \dot{\vec{r}}_i \end{aligned}$$

in non-relativistic case the RH term is just  $2K = \sum m_i \dot{\vec{r}}_i^2$

where  $K =$  total kinetic energy

[ $K$  is sometimes split into kinetic energy of mass motion " $T$ " and kinetic energy of random motions of atoms/molecules/nuclei " $U$ "]

LH side can be rewritten

$$\frac{d}{dt} \sum_i \vec{p}_i \cdot \vec{r}_i = \frac{d}{dt} \sum_i m_i \dot{\vec{r}}_i \cdot \vec{r}_i = \frac{d}{dt} \sum_i \frac{1}{2} \frac{d}{dt} (m_i r_i^2) = \frac{1}{2} \ddot{I}$$

$I = \sum m_i r_i^2$  is the spherical moment of inertia

So now we have the identity

$$\frac{1}{2} \ddot{I} = 2K + \sum_i \vec{F}_i \cdot \vec{r}_i$$

the sum " $\sum_i \vec{F}_i \cdot \vec{r}_i$ " is called the "virial of Clausius"

Only long-range forces and external forces are important - "collisions" between particles contribute two terms of  $\vec{F}_i \cdot \vec{r}_i$  of equal and opposite sign - i.e. they contribute zero to the sum.

In static configurations,  $I =$  constant,  $\dot{I} = \ddot{I} = 0$ , so

$$K = -\frac{1}{2} \sum_i \vec{F}_i \cdot \vec{r}_i$$

## (VIRIAL THEOREM) II.

considering forces between the particles

$$\sum_i \vec{F}_i \cdot \vec{r}_i = \sum_{\text{pairs}} (\vec{F}_{ij} \cdot \vec{r}_i + \vec{F}_{ji} \cdot \vec{r}_j) \quad \text{For Coulomb + gravity}$$

$$\vec{F}_{ij} = -\vec{F}_{ji}$$

Coulomb forces are strong, but star is neutral overall and interactions at "point" —

i.e.  $\vec{r}_i = \vec{r}_j$ , cancel each other, contributing negligibly to  $\sum_i \vec{F}_i \cdot \vec{r}_i$ . So we are left with gravity:

$$F_{ij} = -\frac{G m_i m_j}{(r_{ij})^3} (\vec{r}_i - \vec{r}_j)$$

so our "static" version of the virial theorem becomes

$$K = -\frac{1}{2} \sum_i \vec{F}_i \cdot \vec{r}_i = -\frac{1}{2} \sum_{\text{pairs}} F_{ij} (\vec{r}_i - \vec{r}_j) = \frac{1}{2} \sum_{\text{pairs}} \frac{G m_i m_j}{r_{ij}}$$

so  $K = -\frac{\mathcal{U}}{2}$

$$\mathcal{U} = -\sum \frac{G m_i m_j}{r_{ij}} =$$

total potential energy of system (star in this case)

i.e. negative of P.E. of pairs

This is what is commonly called "the virial theorem" — appropriate for static configurations ( $I = \text{constant}$ ) with interparticle forces with  $\frac{1}{r^2}$  dependence.

Particle kinetic energy scales with temperature....

Consider the "virial of Clausius" for a box of ideal gas at some pressure. Consider only external pressure force acting on the surface of the container — assume interparticle Coulomb forces contribute negligibly to sum.

$$\sum_i \vec{F}_i \cdot \vec{r}_i = \int_S (-P) d\vec{S} \cdot \vec{r} = -P \int_S \vec{r} \cdot \vec{n} dS \quad \vec{n} \text{ is unit vector normal to surface}$$

Apply Gauss' theorem  $= -P \int_V \nabla \cdot \vec{r} dV = -3PV$

so from before  $K = -\frac{1}{2} \sum_i \vec{F}_i \cdot \vec{r}_i = \frac{3}{2} PV - \frac{1}{2} \sum_{\text{pair}} \frac{G m_i m_j}{r_{ij}}$

for ideal gas

interparticle forces e.g. Coulomb

# VIRIAL THEOREM

III

Ideal gas law  $PV = NkT$   
 $P = nkT$

so

$$K = \frac{3}{2} PV = \frac{3}{2} NkT$$

i.e. Kinetic energy per particle is  $\frac{3}{2} kT$

Note that  $K$  only includes translational kinetic energy thus far. Additional degrees of freedom (rotation, vibration, excitation) and photon energy has not yet been included.

For a star in steady-state / equilibrium, we can write the total energy

$$E = K + \Omega = -\frac{\Omega}{2} + \Omega = \frac{\Omega}{2}$$

|
|
  
 Kinetic      potential

so for star composed of ideal gas

$$K = -\frac{\Omega}{2} = \frac{3}{2} NkT \quad E = \frac{\Omega}{2} \quad \text{since } \Omega \text{ negative } E \text{ is negative}$$

if star formed from a very large configuration (through gravitational contraction) then  $\Omega$  initially was  $\approx 0$ . Since observed stars have very negative  $\Omega$ , and energy is conserved overall, then significant  $E$  must have been radiated away, out of the system, during collapse.

When star changes potential energy  $\Omega$ , i.e. contracts, half of the  $\Omega$  is converted to "heat"  $K$ . Since we have a thermal gradient ( $T \downarrow$  as  $R \uparrow$ ) - energy is transported via radiation or convection to the surface of the star and "lost" as luminosity.

$P$ : Pressure  
 $V$ : Volume  
 $N$ : # molecules

$k$ : Boltzmann constant  
 $1.38 \times 10^{-23} \text{ J/K}$   
 $= R/N_A$

$T$ : Temperature

$n$ : # molecules/unit volume

# VIRIAL THEOREM

## IV

Consequences of virial theorem for adiabatic gas

From 1st law of thermo, for  $dQ = 0$

$$dU = -PdV$$

internal energy      work

For adiabatic gas

$$PV^\gamma = \text{constant}$$

$$= TV^{\gamma-1}$$

$$= P^{1-\gamma} T^\gamma$$

$$\frac{dT}{T} + (\gamma-1) \frac{dV}{V} = 0$$

recall  $\gamma = \frac{C_p}{C_v}$  ratio of specific heats at constant P and V

$$dV = -\frac{dT}{T} \cdot \frac{V}{(\gamma-1)}$$

for completely ionized matter there are only 3 deg. of freedom

$$\gamma = 1 + \frac{2}{f} \quad f=3 \quad \gamma = \frac{5}{3}$$

$$dU = -PdV = \frac{PVdT}{T(\gamma-1)}$$

$\gamma = \frac{5}{3}$  for monatomic ideal gas

$$dU = \frac{1}{\gamma-1} Nk \frac{dT}{T}$$

recall  
 $P = nKT$   
 $PV = NkT$

the total kinetic energy of our ideal gas is

$$\frac{PV}{T} = Nk$$

$$K = \frac{3}{2} NkT$$

↑ total particles

so  $dK = \frac{3}{2} Nk dT$

$$dT = \frac{dK}{\frac{3}{2} Nk} = \frac{2}{3 Nk} dK$$

so  $dU = Nk \frac{dT}{\gamma-1} = \frac{Nk}{\gamma-1} \cdot \frac{2}{3 Nk} dK = \frac{2}{3} \frac{dK}{\gamma-1}$

For adiabatic ideal gas

integrating over the whole star  $U = \frac{2}{3} \frac{K}{\gamma-1} \Rightarrow K = \frac{3}{2} (\gamma-1) U$

Write a more general form of an equation for total energy E

$$E = U + \Omega = U - 2K = U - 3(\gamma-1)U = U(1-3\gamma+3)$$

$$E = -(3\gamma-4)U = \frac{3\gamma-4}{3(\gamma-1)} \Omega$$

total energy - generalized for additional degrees of freedom (for adiabatic index  $\gamma$ ) - restatement of virial theorem

for ideal monatomic gas  $\gamma = \frac{5}{3}$   $E = -U = \frac{\Omega}{2}$