

# Relativistic Quantum Mechanics

## Homework 5 (solution)

November 1, 2007

### 1

#### 1.1

$$\psi(x) = \sum_{h_{\pm 1/2}} \int \frac{d^3 k}{(2\pi)^{3/2}} (a(k, h) e^{-ik \cdot x} u(k, h) + b^*(k, h) e^{ik \cdot x} v(k, h)) \quad (1)$$

$$\psi^\dagger(x) = \sum_{h_{\pm 1/2}} \int \frac{d^3 k}{(2\pi)^{3/2}} (a^*(k, h) e^{ik \cdot x} u^\dagger(k, h) + b(k, h) e^{-ik \cdot x} v^\dagger(k, h)) \quad (2)$$

$$\begin{aligned} \psi^\dagger(x)\psi(x) = \sum_{h, h'_{\pm 1/2}} \int \frac{d^3 k d^3 k'}{(2\pi)^3} [ & a^*(k', h') a(k, h) e^{i(k'-k) \cdot x} u^\dagger(k', h') u(k, h) + \\ & a^*(k', h') b^*(k, h) e^{i(k'+k) \cdot x} u^\dagger(k', h') v(k, h) + \\ & b(k', h') b^*(k, h) e^{i(k-k') \cdot x} v^\dagger(k', h') v(k, h) + \\ & b(k', h') a(k, h) e^{-i(k+k') \cdot x} v^\dagger(k', h') u(k, h)] \quad (3) \end{aligned}$$

In the massive normalization

$$u^\dagger(k, h) u(k', h') = \frac{k^0}{m} \delta_{k'k} \delta_{h'h} = v^\dagger(k, h) v(k', h') \quad (4)$$

and

$$u^\dagger v = 0 = v^\dagger u \quad (5)$$

So Eq.3 becomes

$$\psi^\dagger(x)\psi(x) = \sum_{h_{\pm 1/2}} \int \frac{d^3 k}{(2\pi)^3} (|a(k, h)|^2 + |b(k, h)|^2) \frac{k^0}{m} \quad (6)$$

Requiring that

$$\int d^3 x \psi^\dagger(x)\psi(x) = 1$$

we get

$$\sum_{h_{\pm 1/2}} \int d^3 k (|a(k, h)|^2 + |b(k, h)|^2) = \frac{8\pi^3 m}{k^0} \quad (7)$$

## 1.2

In the massless normalization

$$u^\dagger(k, h)u(k', h') = k^0 \delta_{k'k} \delta_{h'h} = v^\dagger(k, h)v(k', h') \quad (8)$$

and we get

$$\sum_{h_{\pm 1/2}} \int d^3k (|a(k, h)|^2 + |b(k, h)|^2) = \frac{8\pi^3}{k^0} \quad (9)$$

## 2

### 2.1

Dirac Hamiltonian in the presence of a constant magnetic field (minimal coupling) becomes

$$H = \vec{\alpha} \cdot (\vec{p} - e\vec{A}) + \beta m \quad (10)$$

Defining

$$\vec{\Pi} = \vec{p} - e\vec{A} \quad (11)$$

we have

$$H = \vec{\alpha} \cdot \vec{\Pi} + \beta m \quad (12)$$

and

$$\begin{aligned} [\Pi_i, H] &= [p_i + eA_i, \alpha_j(p_j - eA_j) + \beta m] \\ &= -\alpha_j [p_i + eA_i, p_j - eA_j] \end{aligned} \quad (13)$$

We also see that

$$[\Pi_i, \Pi_j] \stackrel{[p_i, A_j] \sim \partial_i A_j}{\sim} \epsilon_{ijk} B_k \neq 0 \quad (14)$$

since

$$\begin{aligned} [p_i - eA_i, p_j - eA_j] &= \\ (p_i - eA_i)(p_j - eA_j) - (p_j - eA_j)(p_i - eA_i) &= \\ p_i p_j - e p_i A_j - e A_i p_j + e^2 A_i A_j - p_j p_i + e p_j A_i + e A_j p_i - e^2 A_j A_i &= \\ -e[p_i, A_j] + e[p_j, A_i] &= ie(\partial_i A_j - \partial_j A_i) = -ie\epsilon_{ijk} B_k \neq 0 \end{aligned} \quad (15)$$

And  $\vec{\Pi}$  isn't conserved (similarly  $\vec{p}$  doesn't commute with H).

Total angular momentum  $\vec{J}$  is conserved

$$\begin{aligned}
[J_i, H] &= [\epsilon_{ijk}x_j\Pi_k + \frac{1}{2}\tilde{\alpha}_i, H] = \\
&= [\epsilon_{ijk}x_j\Pi_k + \frac{1}{2}\tilde{\alpha}_i, \alpha_m\Pi_m + \beta m] = \\
&= [\epsilon_{ijk}x_j\Pi_k, \alpha_m\Pi_m] + [\frac{1}{2}\tilde{\alpha}_i, \alpha_m\Pi_m] + [\epsilon_{ijk}x_j\Pi_k, \beta m] + [\frac{1}{2}\tilde{\alpha}_i, \beta m] = \\
&= \epsilon_{ijk}x_j\alpha_m[\Pi_k, \Pi_m] + \epsilon_{ijk}\alpha_m[x_j, \Pi_m]\Pi_k + \frac{1}{2}[\tilde{\alpha}_i, \alpha_m]\Pi_m = \\
&= \epsilon_{ijk}x_j\alpha_m(-ie)\epsilon_{kml}B_l + i\epsilon_{ijk}\alpha_m\delta_{jm}\Pi_k + i\epsilon_{imn}\alpha_n\Pi_m = \\
&= -ie\epsilon_{ijk}\epsilon_{mlk}x_j\alpha_mB_l + i\epsilon_{ijk}\alpha_j\Pi_k + i\epsilon_{imn}\alpha_n\Pi_m = \\
&= -ie\epsilon_{ijk}\epsilon_{mlk}x_j\alpha_mB_l + i\epsilon_{ijk}\alpha_j\Pi_k + i\epsilon_{ijk}\alpha_k\Pi_j = \\
&= -ie\epsilon_{ijk}\epsilon_{mlk}x_j\alpha_mB_l + i\epsilon_{ijk}\alpha_j\Pi_k - i\epsilon_{ijk}\alpha_j\Pi_k = 0
\end{aligned} \tag{16}$$

and

$$\begin{aligned}
H^2 \{\alpha_i, \beta\}=0 \vec{\Pi}^2 + m^2 \mathbf{1} &= \\
\alpha_i\alpha_j(p_i - eA_i)(p_j - eA_j) + m^2 &= \\
(\frac{1}{2}\{\alpha_i, \alpha_j\} + \frac{1}{2}[\alpha_i, \alpha_j])\Pi_i\Pi_j + m^2 &= \\
(\delta_{ij} + i\epsilon_{ijk}\alpha_k)\Pi_i\Pi_j + m^2 &= \\
\vec{\Pi}^2 + i\epsilon_{ijk}\alpha_k\Pi_i\Pi_j + m^2 &= \\
\vec{\Pi}^2 + i\epsilon_{ijk}\alpha_k\frac{1}{2}[\Pi_i, \Pi_j] + m^2 &= \\
\vec{\Pi}^2 + \epsilon_{ijk}\alpha_k\frac{1}{2}e\epsilon_{ijm}B_m + m^2 &= \\
\vec{\Pi}^2 + \alpha_k\delta_{km}B_m + m^2 &= \\
\vec{\Pi}^2 + \vec{\alpha} \cdot \vec{B} + m^2 &=
\end{aligned} \tag{17}$$

Since  $[H, H^2] = 0$  we see that  $\vec{\Pi}^2 + \vec{\alpha} \cdot \vec{B} + m^2$  is a conserved quantity (such as energy and total angular momentum).

## 2.2

Solving the characteristic equation for the Hamiltonian

$$\begin{vmatrix} (m - E)\mathbf{1} & \vec{\sigma} \cdot (\vec{p} - e\vec{A}) \\ \vec{\sigma} \cdot (\vec{p} - e\vec{A}) & -(m + E)\mathbf{1} \end{vmatrix} = 0 \tag{18}$$

or

$$\begin{aligned}
(\vec{\sigma} \cdot (\vec{p} - e\vec{A}))(\vec{\sigma} \cdot (\vec{p} - e\vec{A})) &= E^2 - m^2 \\
\Rightarrow E^2 - m^2 &= [(\vec{p} - e\vec{A}) \cdot (\vec{p} - e\vec{A}) - e\vec{\sigma} \cdot \vec{B}]
\end{aligned} \tag{19}$$

## 3

We need to show that

$$\bar{u}(p_2)\gamma^\mu u(p_1) = \frac{(p_1 + p_2)^\mu}{2m}\bar{u}(p_2)u(p_1) - \frac{i(p_1 - p_2)_\nu}{2m}\bar{u}(p_2)\sigma^{\mu\nu}u(p_1) \tag{20}$$

where,

$$(\gamma^\mu p_{1\mu} - m)u(p_1) = 0, \quad \bar{u}(p_2)(\gamma^\mu p_{2\mu} - m) = 0 \quad (21)$$

We start with

$$\begin{aligned} \bar{u}(p_2)\sigma^{\mu\nu}(p_1 - p_2)_\nu u(p_1) &= \\ \bar{u}(p_2)i[(\gamma^\mu\gamma^\nu - \eta^{\mu\nu})p_{1\nu} - (\eta^{\mu\nu} - \gamma^\nu\gamma^\mu)p_{2\nu}]u(p_1) &= \\ \bar{u}(p_2)i[\gamma^\mu\gamma^\nu p_{1\nu} - p_1^\mu - p_2^\mu + \gamma^\nu p_{2\nu}\gamma^\mu]u(p_1) &= \\ \bar{u}(p_2)i[\gamma^\mu m - (p_1 + p_2)^\mu + m\gamma^\mu]u(p_1) & \end{aligned} \quad (22)$$

or that

$$-\frac{i}{2m}\bar{u}(p_2)\sigma^{\mu\nu}(p_1 - p_2)_\nu u(p_1) = -\bar{u}(p_2)\left[\frac{1}{2m}(p_1 + p_2)^\mu - \gamma^\mu\right]u(p_1) \quad (23)$$

So that RHS of Eq.20 gives

$$RHS = \bar{u}(p_2)\gamma^\mu u(p_1) \quad (24)$$