

## Physics 123 - February 25, 2013

- Exam 1 - Thursday 8 AM - B+L 106
- Calculator
- 1 side 8.5 x 11 inch sheet w/ formulas/notes
- Q + A session

LAST TIME

Fourier Analysis

$$f(x) = \frac{A_0}{2} + \sum_{m=1}^{\infty} A_m \cos(mkx) + \sum_{m=1}^{\infty} B_m \sin(mkx)$$

↳ periodic function

↳ more detail means  
Higher frequency components  
Have more relative importance

as  $\lambda$  becomes large  
function becomes anharmonic (e.g. a wave pulse)

$$F(x) = \frac{1}{\pi} \left[ \int_0^{\infty} A(k) \cos(kx) dx + \int_0^{\infty} B(k) \sin(kx) dx \right]$$

with

$$A(k) = \int_{-\infty}^{\infty} F(x) \cos(kx) dx$$

cosine transform

Fourier integrals

$$B(k) = \int_{-\infty}^{\infty} F(x) \sin(kx) dx$$

sine transform

Describe  $F(x)$  in frequency space  $k$   
 $F(t)$   $\omega$



Narrower functions / sharper details in position or time  
require broader functions in  $k$  or  $\omega$

Very powerful - will see this again + again

... Data compression  
for example

Expect you to do a basic Fourier problem  
and just be aware of integral + Transforms for now

Essence of the uncertainty principle in quantum mechanics

⋮

# Electromagnetic Waves

Review ...

Derivation of EM plane waves

Will go thru fast

Want you to see it and

Believe!!



Not an  $\epsilon + M$  course

↳ you're responsible to  
understand / use end result

**Buckle your seat belts**

## Integral form of Maxwell's equations

Gauss  $\oint_S \vec{E} \cdot d\vec{A} = \frac{Q_{enc}}{\epsilon_0}$

No magnetic monopoles  $\oint_S \vec{B} \cdot d\vec{A} = 0$

Ampere  $\oint_C \vec{B} \cdot d\vec{l} = \mu_0 \int_S \vec{j} \cdot d\vec{A} + \mu_0 \epsilon_0 \frac{d}{dt} \int_S \vec{E} \cdot d\vec{A}$

Faraday  $\oint_C \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{a}$

new term -  
"Maxwell's  
Displacement  
current"

Will now derive "differential" form ...

# Divergence of vector field $\vec{V}$ (cartesian coordinates)

$$\vec{\nabla} \cdot \vec{V} = \frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} + \frac{\partial V_z}{\partial z}$$

measures divergence or convergence of field

*can prove to yourself*



+ Divergence



- Divergence



0 Divergence

$$\vec{\nabla} = \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k}$$

Degree of "Divergence"  
or "convergence"

$$\vec{\nabla} \cdot \vec{V} = \left( \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k} \right) \cdot \left( V_x \hat{i} + V_y \hat{j} + V_z \hat{k} \right)$$

$$\text{use } \hat{i} \cdot \hat{i} = \hat{j} \cdot \hat{j} = \hat{k} \cdot \hat{k} = 1$$

$$\hat{i} \cdot \hat{j} = \hat{i} \cdot \hat{k} = \hat{j} \cdot \hat{k} = 0$$

vector calculus

Gauss' Theorem  
Green's Theorem  
divergence theorem

$$\int_{Vol} (\nabla \cdot \vec{V}) dv = \int_{Surf} \vec{V} \cdot d\vec{A}$$

for vector field  $\vec{V}$

Gauss

$$\oint_S \vec{E} \cdot d\vec{A} = \frac{Q_{enc}}{\epsilon_0} \rightarrow \int_S \vec{E} \cdot d\vec{A} = \int_V (\nabla \cdot \vec{E}) dv = \frac{\int_V \rho dv}{\epsilon_0}$$

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

No magnetic monopoles

$$\oint_S \vec{B} \cdot d\vec{A} = 0 \rightarrow \int_S \vec{B} \cdot d\vec{A} = \int_V (\nabla \cdot \vec{B}) dv = 0$$

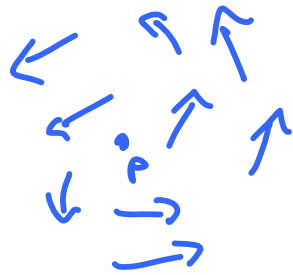
$$\nabla \cdot \vec{B} = 0$$

# Curl of vector field

$$\text{curl } \vec{V} \equiv \vec{\nabla} \times \vec{V} \longrightarrow \underline{\underline{\text{Vector}}}$$

Cartesian

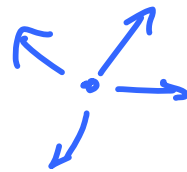
$$\vec{\nabla} \times \vec{V} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ v_x & v_y & v_z \end{vmatrix} = \hat{i} \left( \frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z} \right) - \hat{j} \left( \frac{\partial v_z}{\partial x} - \frac{\partial v_x}{\partial z} \right) + \hat{k} \left( \frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial y} \right)$$



large  
curl



0 curl



degree of "circulation" of vector field  
about point

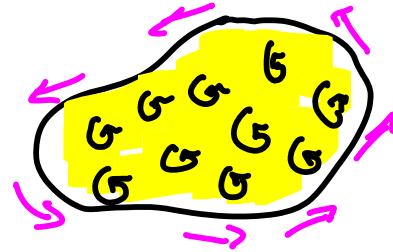


# Stoke's Theorem

$$\oint_C \vec{v} \cdot d\vec{l} = \int_S (\vec{\nabla} \times \vec{v}) \cdot d\vec{A}$$

Faraday (induction)

$$\oint_C \vec{E} \cdot d\vec{l} = -\frac{d\Phi_M}{dt} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{A}$$



$$\int_S (\vec{\nabla} \times \vec{E}) \cdot d\vec{A} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{A}$$

$$\vec{\nabla} \times \vec{E} = -\frac{d\vec{B}}{dt}$$

$$\oint_{\mathcal{L}} \vec{B} \cdot d\vec{l} = \mu_0 \oint_{\mathcal{S}} \vec{j} \cdot d\vec{A} + \mu_0 \epsilon_0 \frac{\partial}{\partial t} \oint_{\mathcal{S}} \vec{E} \cdot d\vec{A}$$

Ampere's Law  
plus  
Maxwell's  
displacement  
current

$$\oint_{\mathcal{S}} (\vec{\nabla} \times \vec{B}) \cdot d\vec{A} = \mu_0 \oint_{\mathcal{S}} \vec{j} \cdot d\vec{A} + \mu_0 \epsilon_0 \frac{\partial}{\partial t} \oint_{\mathcal{S}} \vec{E} \cdot d\vec{A}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{j} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Note Title

$$\oint_S \vec{E} \cdot d\vec{A} = \frac{Q_{enc}}{\epsilon_0} \rightarrow \int_V (\vec{\nabla} \cdot \vec{E}) dV = \int_V \rho dV$$

Summary  
summary

$$\vec{\nabla} \cdot \vec{E} = \rho$$

$$\oint_S \vec{B} \cdot d\vec{A} = 0 \rightarrow \int_V (\vec{\nabla} \cdot \vec{B}) dV = 0$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\oint_C \vec{B} \cdot d\vec{l} = \mu_0 \int_S \vec{j} \cdot d\vec{A} + \mu_0 \epsilon_0 \frac{d}{dt} \int_S \vec{E} \cdot d\vec{A}$$

$$\hookrightarrow \int_S (\vec{\nabla} \times \vec{B}) \cdot d\vec{A} = \mu_0 \int_S \vec{j} \cdot d\vec{A} + \mu_0 \epsilon_0 \frac{d}{dt} \int_S \vec{E} \cdot d\vec{A}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{j} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\oint_C \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{A}$$

$$\hookrightarrow \int_S (\vec{\nabla} \times \vec{E}) \cdot d\vec{A} = -\frac{\partial}{\partial t} \int_S \vec{B} \cdot d\vec{A} \rightarrow \vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Differential form  
of Maxwell's eqns

region w/ no current

$$\vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Take curl of Both sides

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{B}) = \mu_0 \epsilon_0 \frac{\partial (\vec{\nabla} \times \vec{E})}{\partial t}$$

and it was  
said ...

$$\vec{\nabla} \times \vec{\nabla} \times \vec{B} = \vec{\nabla} (\vec{\nabla} \cdot \vec{B}) - \vec{\nabla} \cdot (\vec{\nabla} \vec{B})$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

Known  
as

the Laplacian

$$\nabla^2 \vec{B}$$

Laplacian of scalar field T

$$\nabla^2 T = \vec{\nabla} \cdot (\vec{\nabla} T) = \left( \hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) \cdot \left( \hat{i} \frac{\partial T}{\partial x} + \hat{j} \frac{\partial T}{\partial y} + \hat{k} \frac{\partial T}{\partial z} \right)$$

$$\nabla^2 T = \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2}$$

Laplacian of vector field  $\vec{v}$

$$\nabla^2 \vec{v} = (\nabla^2 v_x) \hat{i} + (\nabla^2 v_y) \hat{j} + (\nabla^2 v_z) \hat{k}$$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{B}) = \mu_0 \epsilon_0 \frac{\partial (\vec{\nabla} \times \vec{E})}{\partial t}$$

$$\vec{\nabla} \cdot (\vec{\nabla} \vec{B}) =$$

$$\vec{\nabla}^2 \vec{B} = \mu_0 \epsilon_0 \frac{\partial (\nabla \times \vec{E})}{\partial t} = \mu_0 \epsilon_0 \frac{\partial}{\partial t} \left( -\frac{\partial \vec{B}}{\partial t} \right)$$

from another  
Maxwell  
eqn

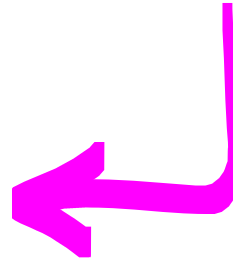
$$\nabla^2 \vec{B} = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

$$\nabla^2 B_x = -\mu_0 \epsilon_0 \frac{\partial^2 B_x}{\partial t^2}$$

$$\nabla^2 B_y = -\mu_0 \epsilon_0 \frac{\partial^2 B_y}{\partial t^2}$$

$$\nabla^2 B_z = -\mu_0 \epsilon_0 \frac{\partial^2 B_z}{\partial t^2}$$

$$\frac{\partial^2 \vec{B}_x}{\partial x^2} + \frac{\partial^2 \vec{B}_x}{\partial y^2} + \frac{\partial^2 \vec{B}_x}{\partial z^2} = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{B}_x}{\partial t^2}$$



$$\vec{\nabla} \times (\vec{\nabla} \times \vec{B}) = \mu_0 \epsilon_0 \frac{\partial (\vec{\nabla} \times \vec{E})}{\partial t}$$

Note Title

11/18/2008

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B}) =$$

$$\vec{\nabla}^2 \vec{B} = \mu_0 \epsilon_0 \frac{\partial (\vec{\nabla} \times \vec{E})}{\partial t} = \mu_0 \epsilon_0 \frac{\partial}{\partial t} \left( -\frac{\partial \vec{B}}{\partial t} \right)$$

from another  
Maxwell  
eqn

$$\nabla^2 \vec{B} = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

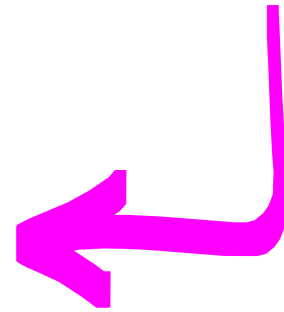
$$\nabla^2 B_x = -\mu_0 \epsilon_0 \frac{\partial^2 B_x}{\partial t^2}$$

$$\nabla^2 B_y = -\mu_0 \epsilon_0 \frac{\partial^2 B_y}{\partial t^2}$$

$$\nabla^2 B_z = -\mu_0 \epsilon_0 \frac{\partial^2 B_z}{\partial t^2}$$



$$\frac{\partial^2 B_x}{\partial x^2} + \frac{\partial^2 B_x}{\partial y^2} + \frac{\partial^2 B_x}{\partial z^2} = -\mu_0 \epsilon_0 \frac{\partial B_x}{\partial t^2}$$



Similarly

$$\nabla^2 \vec{B} = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \left\{ \nabla \times \right.$$

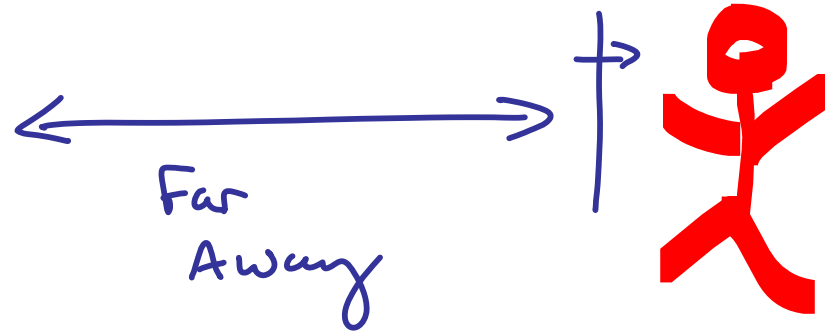
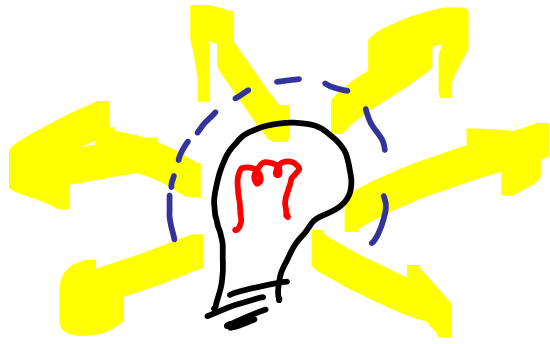
$$\nabla^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

1d wave prop in x

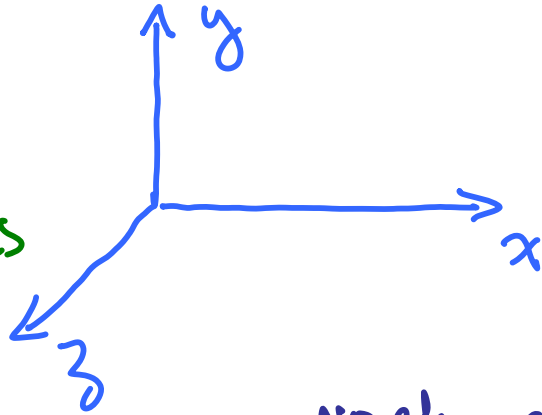
$$\frac{\partial^2 y}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 y}{\partial t^2}$$

wave eqns for  $\vec{E}, \vec{B}$   
w/ velocity of  
propagation

$$\sqrt{\frac{1}{\mu_0 \epsilon_0}} = c$$



Impose  
Boundary  
Conditions



$$\vec{E} = \vec{E}(x, t)$$

no charge  $\rightarrow \rho = 0$

$$\vec{\nabla} \cdot \vec{E} = \rho / \epsilon_0$$

$$\vec{\nabla} \cdot \vec{E}_0 = 0$$

$$\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$$

$\swarrow \quad \swarrow$   
 $0 \quad 0$

So  $E_x$  is const for all  $x \rightarrow E_x = 0$

$E_y$  or  $E_z$  might be nonzero

**E** is TRANSVERSE to Direction of Propagation

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Equivalent to this

$E = f(x, t)$  only

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = -\frac{\partial B_x}{\partial t}$$

Impose Boundary Condition

$$\vec{E} = E(x, t) \hat{j}$$

"polarization"

$$-\left( \frac{\partial E_z}{\partial x} - \frac{\partial E_x}{\partial z} \right) = -\frac{\partial B_y}{\partial t}$$

$$\left( \frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right) = -\frac{\partial B_z}{\partial t}$$

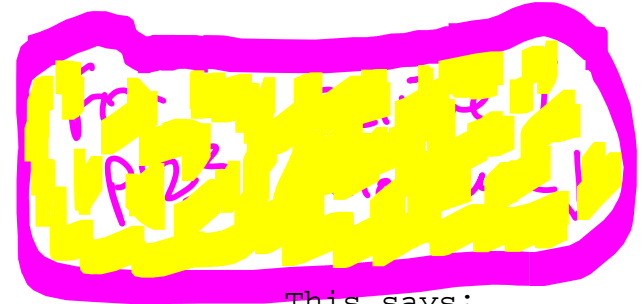
$B_x$  is CONSTANT in time

$B_y$  is CONSTANT in time

Time dependent B field

$\perp$  to  $\vec{E}$

$\vec{E}$ ,  $\vec{B}$  are Transverse  
and mutually  $\perp$



This says:  
for P123, know how to do  
and use below. above is so  
you have seen it before

$$E_y(x,t) = E_{0y} \cos(kx - \omega t + \phi)$$

$$\frac{2\pi}{\lambda}$$

$$\frac{2\pi}{T}$$

Phase  
(initial condition)

$$B_z = - \int \frac{\partial E_y}{\partial x} dt$$

$$B_z = \int k E_{0y} \sin(kx - \omega t) dt$$

$$B_z = \frac{k}{\omega} E_{0y} \cos(kx - \omega t)$$

set  
 $\phi = 0$   
for now

$$\frac{\omega}{k} = \frac{\frac{2\pi}{\lambda}}{\frac{2\pi}{\lambda}} = \frac{1}{\lambda} = \frac{1}{c}$$

$$B_z = \frac{1}{c} E_y$$

$E, B$  are coupled, time dependent

Mutually  $\perp$

$$|\vec{E}| = c |\vec{B}|$$

in phase in

Time + space

