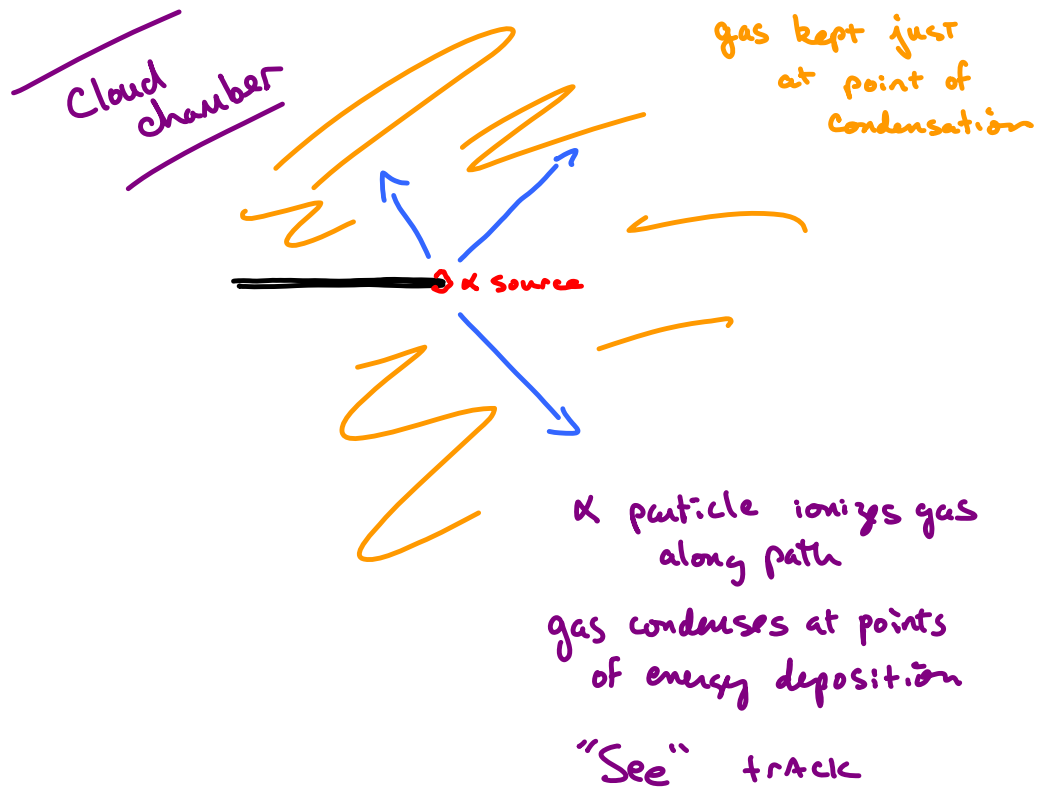
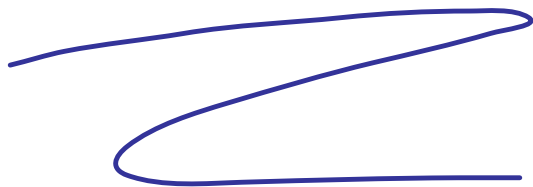


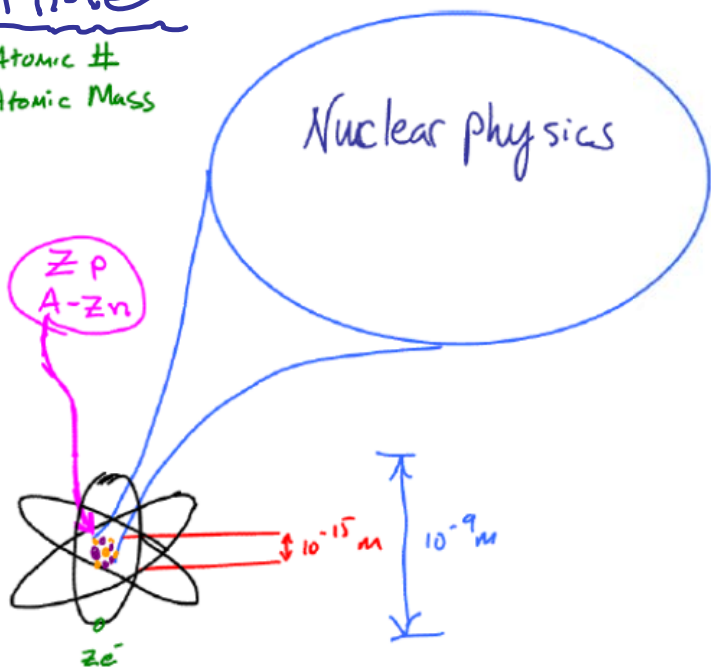
Physics 123 - April 29, 2013

Continuing
with
Nuclear
Physics

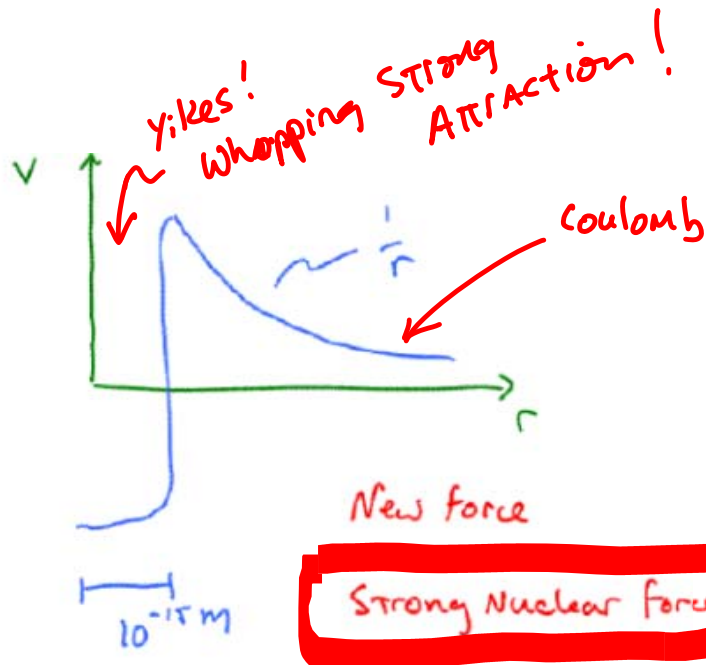


Last Time

$Z \equiv$ Atomic #
 $A \equiv$ Atomic Mass



Particle	chg	mass (mev/c ²)	Spin
Proton, p	+e	938.28	1/2
neutron, n	0	939.37	1/2
electron, e	-e	0.511	1/2



Nuclear radius $\sim 1-10$ Fermi ($10^{-15} m$)

Nuclear density is CONSTANT

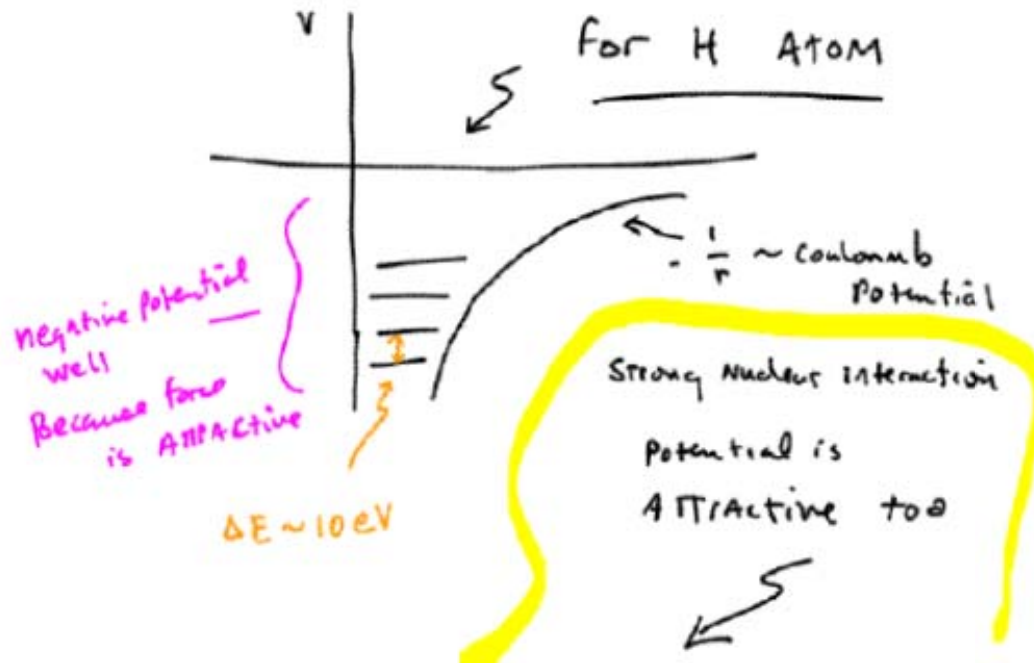
$$\frac{A}{\frac{4}{3}\pi R^3} \sim \text{const.}$$

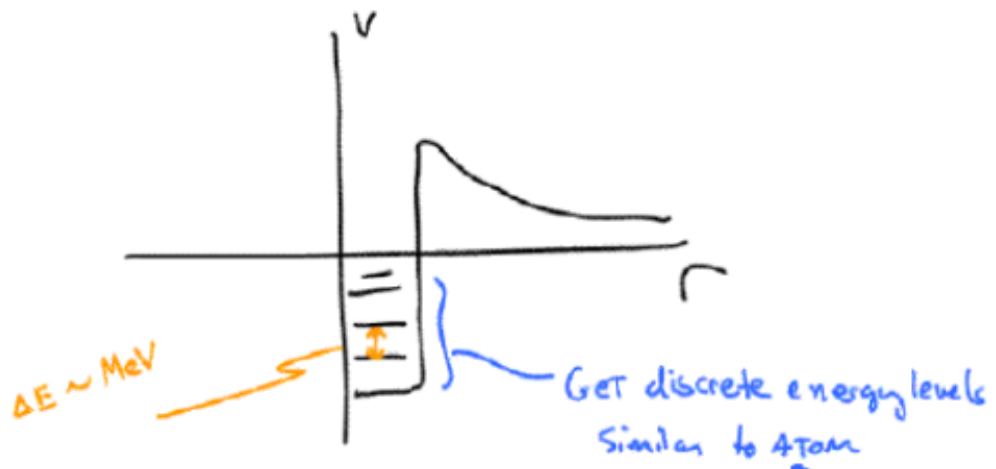
$$A \sim R^3 \quad R \sim A^{1/3}$$

$$\rho \sim 10^{17} \text{ kg/m}^3$$

At this density

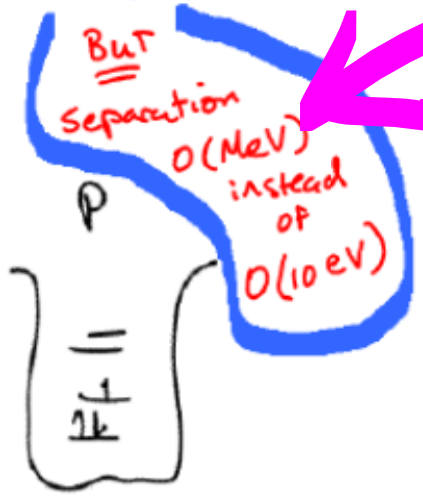
$M_{\text{earth}} \rightarrow$ ball 140 m radius





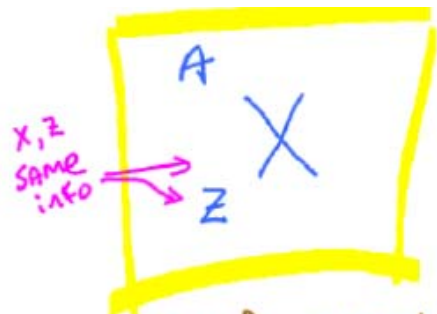
p, n Spin 1/2

fill shells w/ protons + neutrons in much the same way we fill atomic shells



BUT separation $O(\text{MeV})$ instead of $O(10\text{eV})$

Why do you care?



X ≡ Atomic Symbol

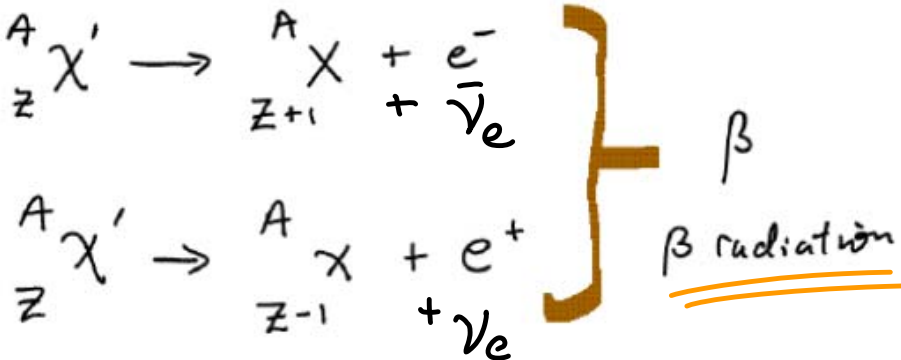
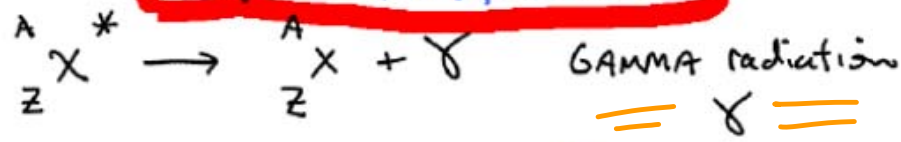
Z ≡ # protons

Atomic #

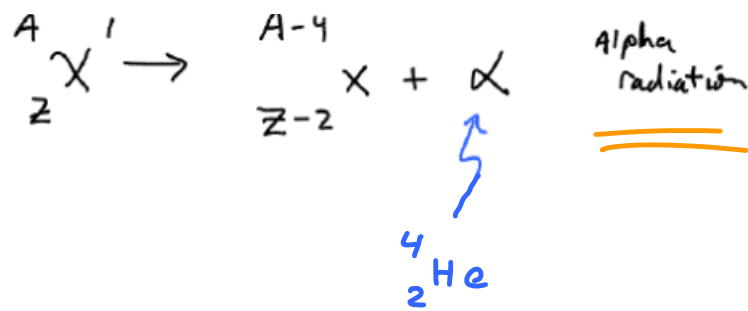
A ≡ Atomic Mass

Symbol for a nucleus $N = A - Z \equiv \# \text{ neutrons}$

Isotope → SAME Z, Different A

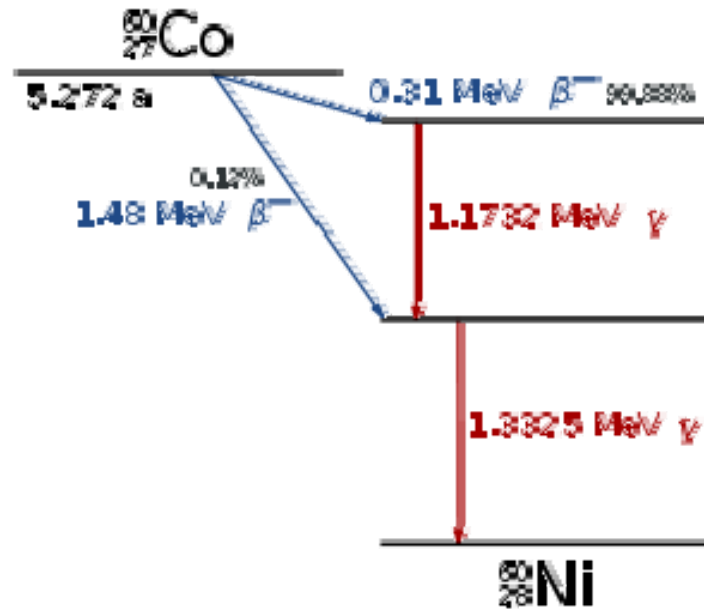


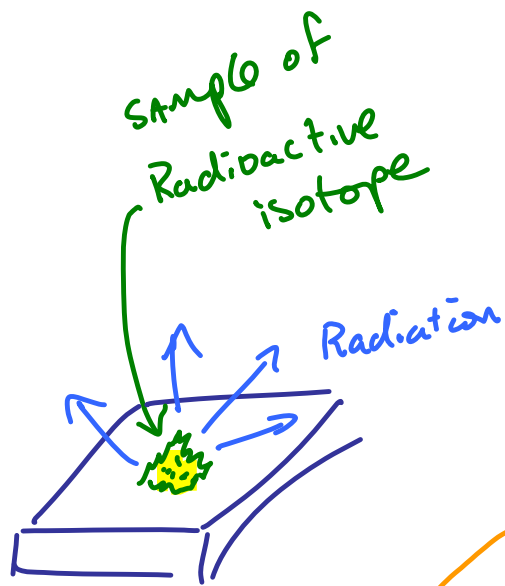
protons + neutrons collectively called "Nucleons"



Transitions between States Very Energetic

Example
nuclear
decay
scheme





decay CONSTANT

$$\frac{dN}{dt} = -\lambda N$$

$N \equiv$ # radioactive nuclei in sample

Activity \equiv # nuclei decaying per second

measured in Curies
 $1 \text{ Ci} = 3.7 \times 10^{10} \frac{\text{dec.}}{\text{s}}$

$$\frac{dN}{N} = -\lambda dt$$

$$\int_{N_0}^N \frac{dN}{N} = - \int_0^t \lambda dt$$

$$\ln \frac{N}{N_0} = -\lambda t$$

$$\frac{N}{N_0} = e^{-\lambda t}$$

$$\rightarrow N = N_0 e^{-\lambda t}$$

$\lambda \equiv$ decay constant

nuclei: drop by $\frac{1}{e}$ when $t = \frac{1}{\lambda}$

"fast decay" \rightarrow large λ

$$t_{\frac{1}{2}} = \frac{0.693}{\lambda} \equiv \text{half-life} \equiv \text{time for } \frac{1}{2} \text{ the nuclei to decay}$$



N atoms in sample
at time = 0

$$\text{Activity} \equiv \frac{\# \text{decays}}{\text{second}} = \frac{\Delta N}{\Delta t} = \lambda N$$

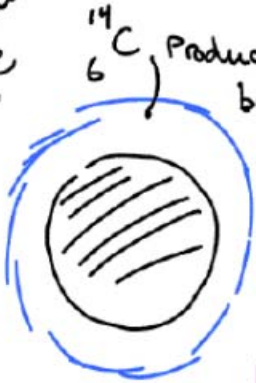
↑
decay CONSTANT

half-life $\equiv t_{1/2}$ = time for $\frac{1}{2}$ sample to decay

$$t_{1/2} = \frac{0.693}{\lambda}$$

Radioactive Dating

Normal ^{12}C
6



^{14}C Produced by cosmic rays hitting atmosphere

^{14}C is naturally radioactive β -emitter
 $t_{1/2} = 5730$ years

- ^{14}C incorporated into living tissue
- Stops at death
- $^{14}\text{C}/^{12}\text{C}$ ratio gives estimate of time since death

^{14}C concentration in atmosphere varies, calibrate w/ tree rings

A sample of bone

^{14}C activity to be 25% of that in living material. How old is the bone?

$$^{14}\text{C} \quad t_{1/2} = 5730$$

radio carbon dating
calibration
Data

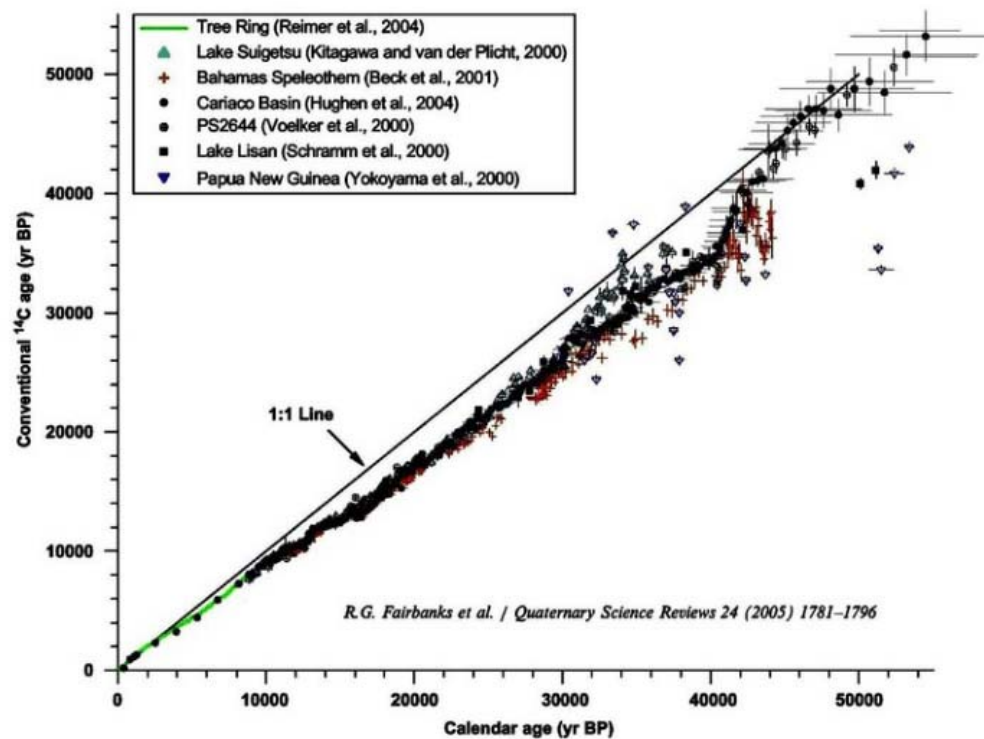


Fig. 1. Composite figure showing the range of radiocarbon calibration data from a variety of archives, including Bahamian speleothems (Beck et al., 2001), marine sediments (Hughen et al., 2000, 2004; Voelker et al., 2000; lake sediments (Schramm et al., 2000; Kitagawa and van der Plicht, 2000), corals (Yokoyama, 2000), and tree rings (Reimer et al., 2004; Friedrich et al., 2004).

Nuclear Binding Energy

consider atomic binding energy (ionization energy)

$$\underbrace{m_e c^2 + m_p c^2}_{\text{Energy of separate parts}} - \underbrace{M_H c^2}_{\text{energy of Bound beast}} = 13.6 \text{ eV}$$

Binding energy

For atoms ... ionization energy is what we call it.

For a nucleus "x" with A, Z:

$$(A-Z)M_n c^2 + ZM_p c^2 - M_x c^2 = \text{Total Nuclear Binding Energy}$$

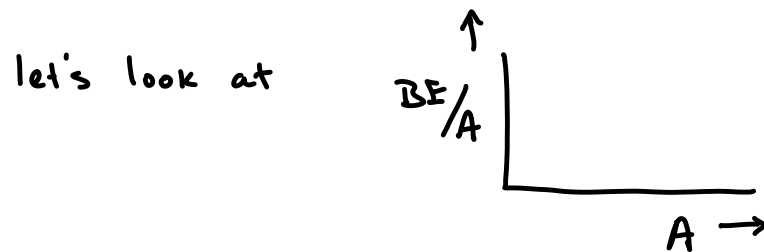
Nuclear Binding Energy

$$\text{Binding Energy per Nucleon} = \frac{\text{Total BE}}{A}$$

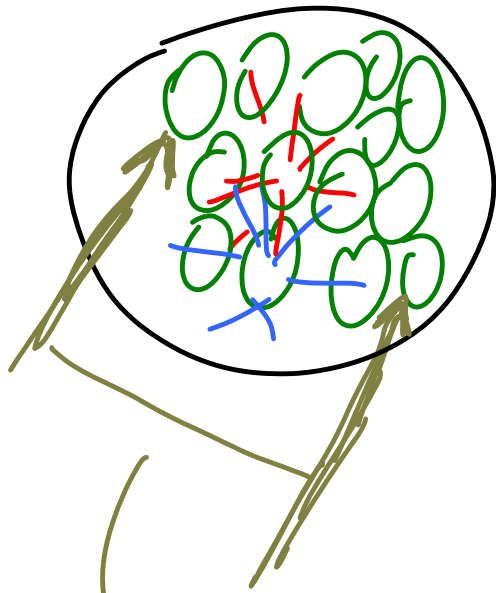
\equiv Amount of E to remove a nucleon if you will
"ionization energy for a nucleon" sort of

High $\frac{\text{BE}}{A} \rightarrow$ Very stable nucleus

Low $\frac{\text{BE}}{A} \rightarrow$ less stable nucleus

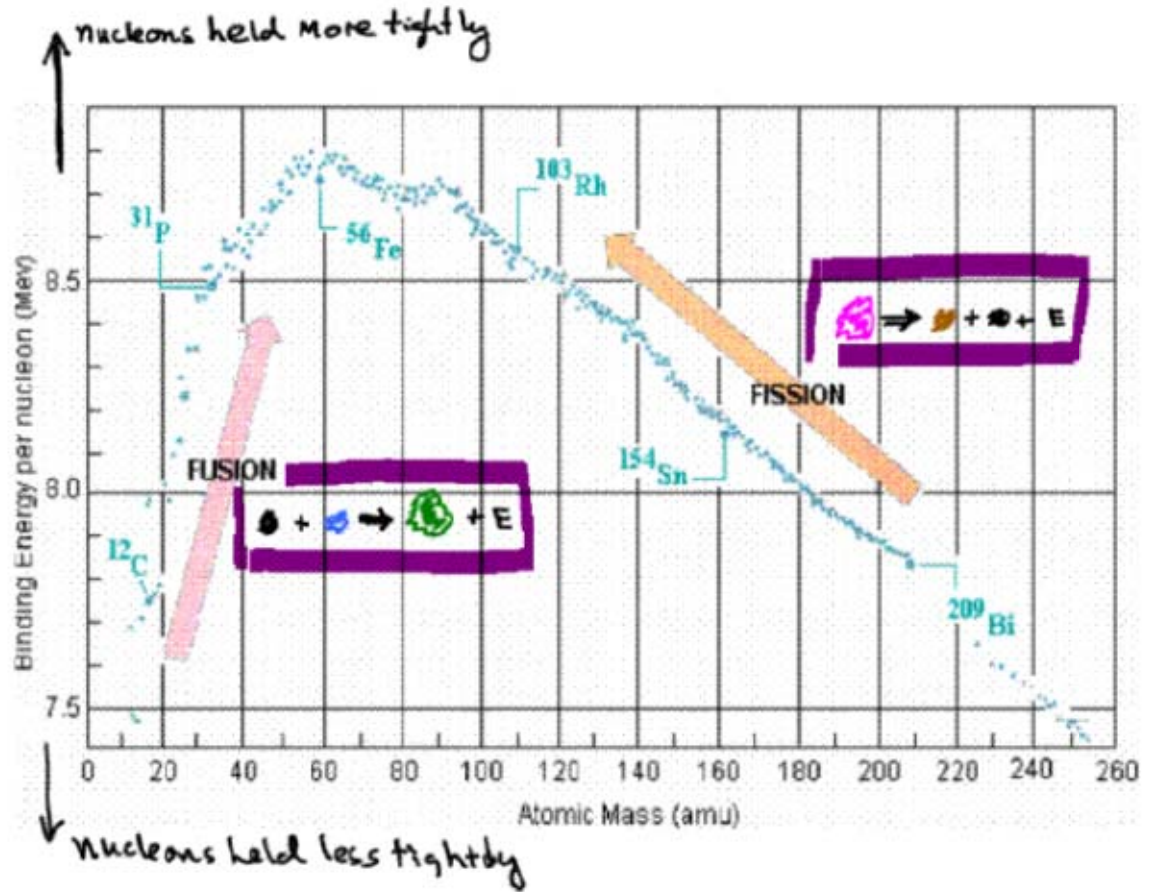


recall range of str. force (10^{-15} m)



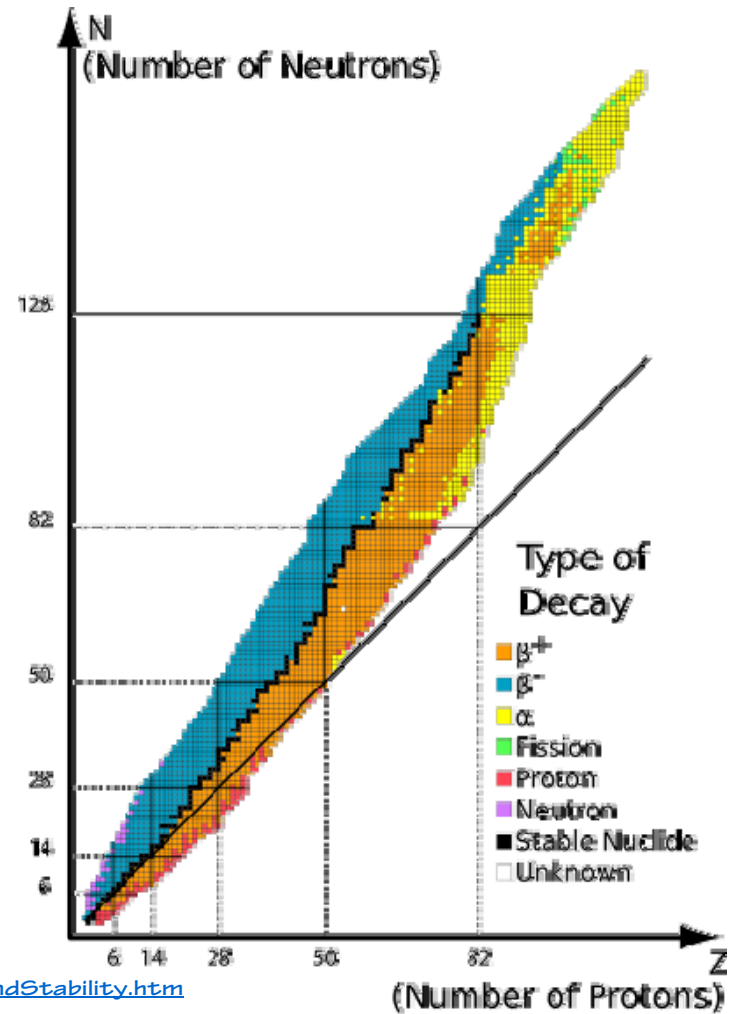
over large distances
→ coulomb

Inherent Nuclear stability as function of nuclear size



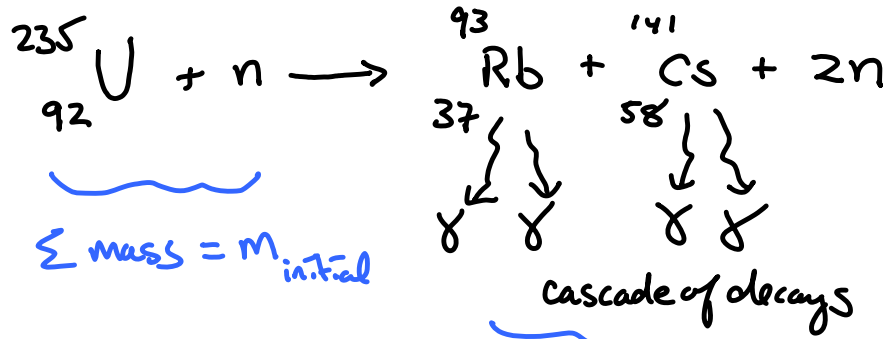
line of Stability

$$\#n > \#p$$



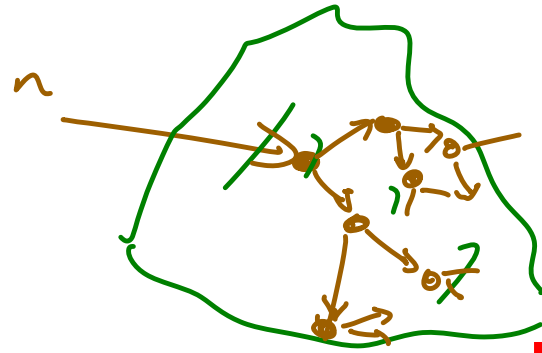
Nice figure from <http://www.kentchemistry.com/links/Nuclear/BandStability.htm>

A very important nuclear process ...



$$M_{\text{initial}} - M_{\text{final}} > 0 \qquad \sum \text{mass} = M_{\text{final}}$$

Fun with Fission

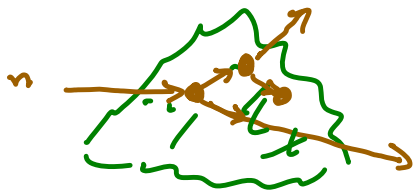


Supercritical
> 1 split per split

Watchout!

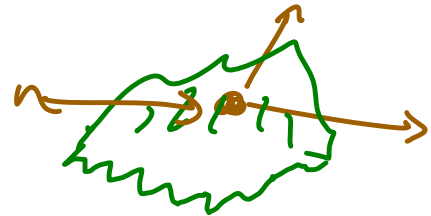
Kablooey!

or
meltdown
+
fine



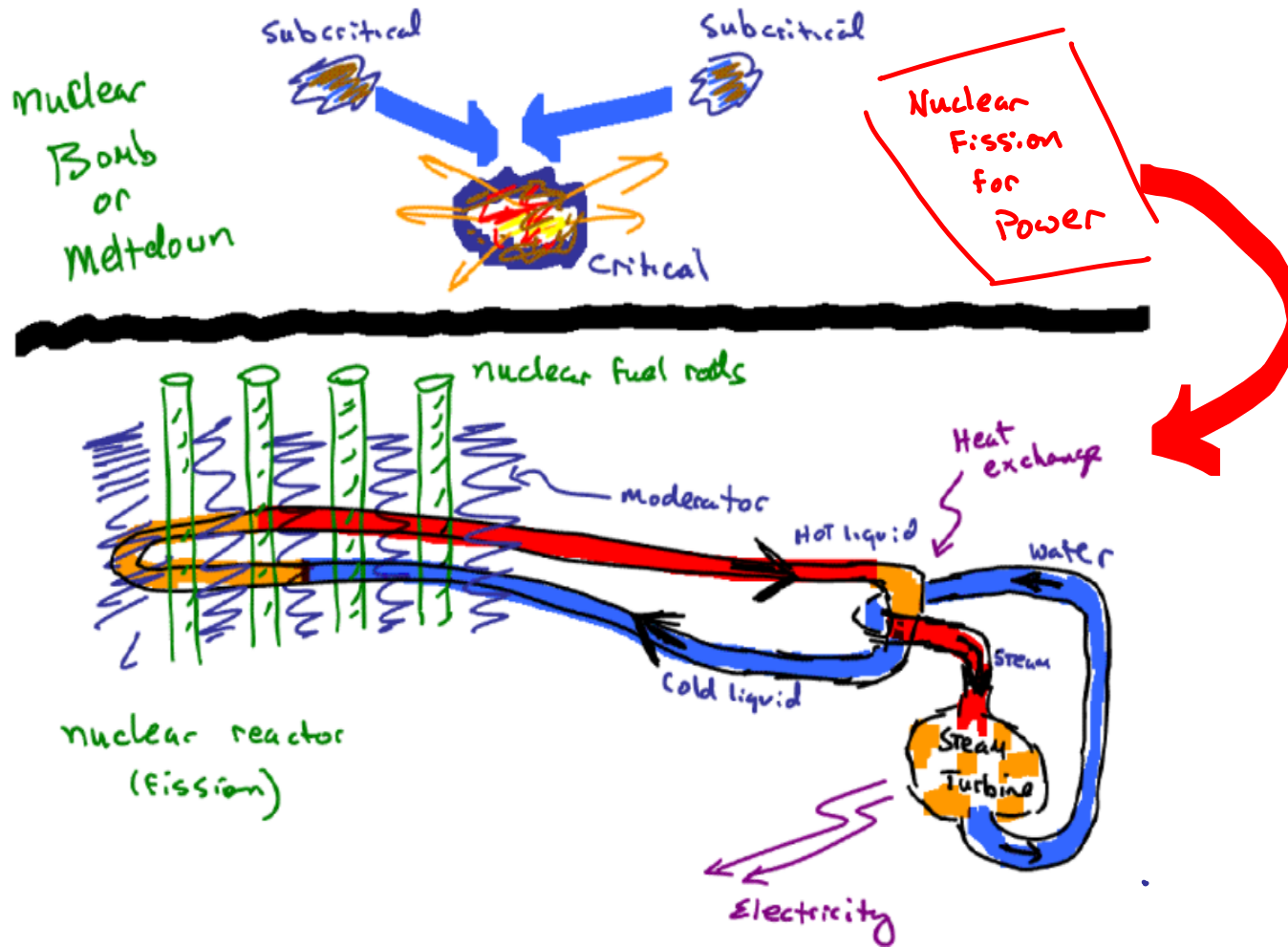
Critical

1 Split per Split

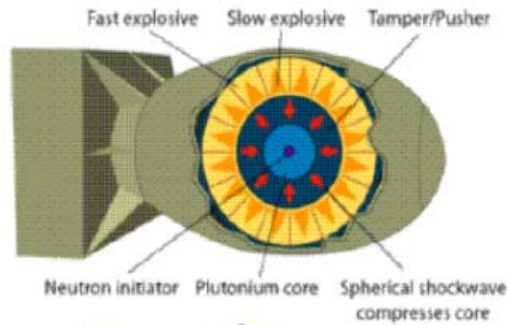


Subcritical

< 1 Split per Split

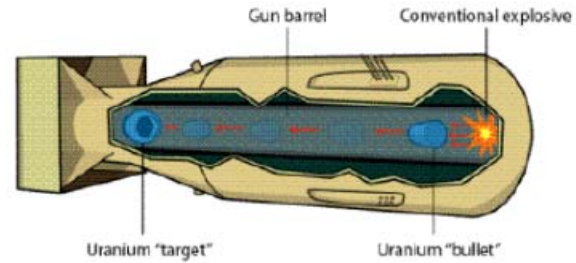


Implosive-type
Technically harder
Than
Gun-type



Similar to Fat Man
Used on Nagasaki
Aug. 9, 1945

Similar to "Little Boy"
used on Hiroshima
August 6, 1945



Diagrams
From
Wikipedia

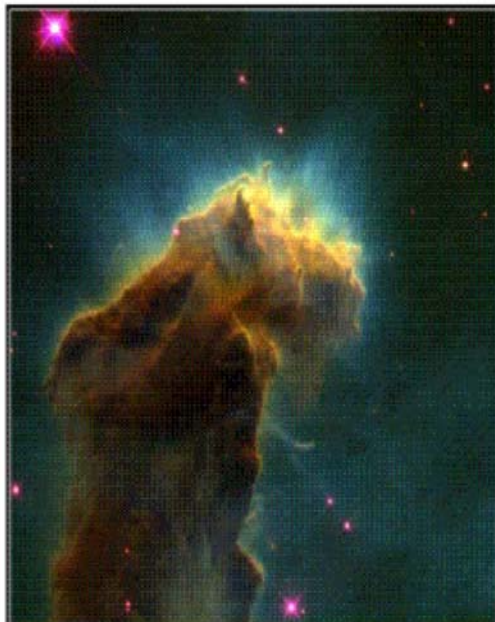


Natural U
~99% ^{238}U
Must enrich
in ^{235}U to
use as
Nuclear fuel
or bomb

Gun-type

isotopes
very difficult
to impossible
to separate
chemically

Stars - from dust to dust



Star-Birth Clouds - M16 HST - WFPC2
PRC83-46b - ST ScI OPO - November 2, 1995
J. Hester and P. Scowen (AZ State Univ.), NASA

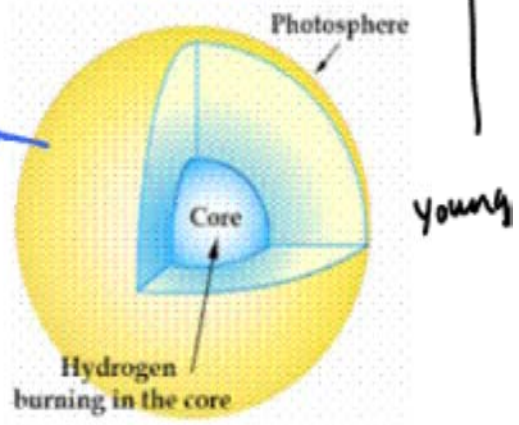
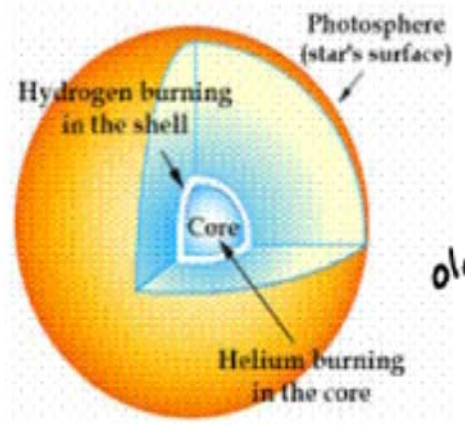
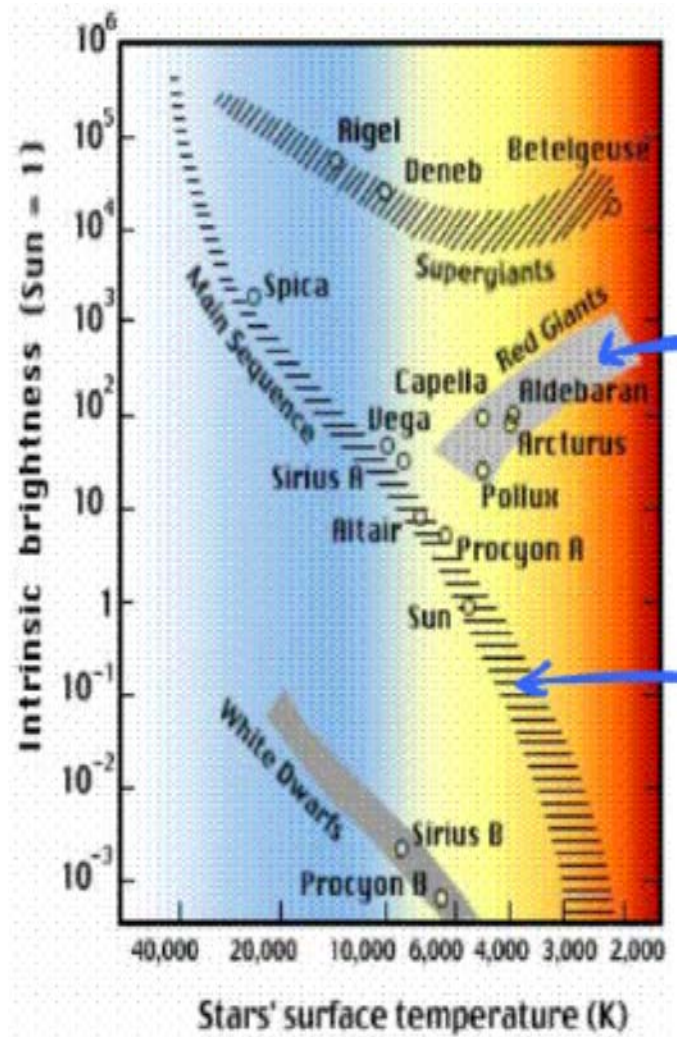
Stars form from
condensation of gas/dust
due to gravitation

mostly hydrogen gas



The Pleiades

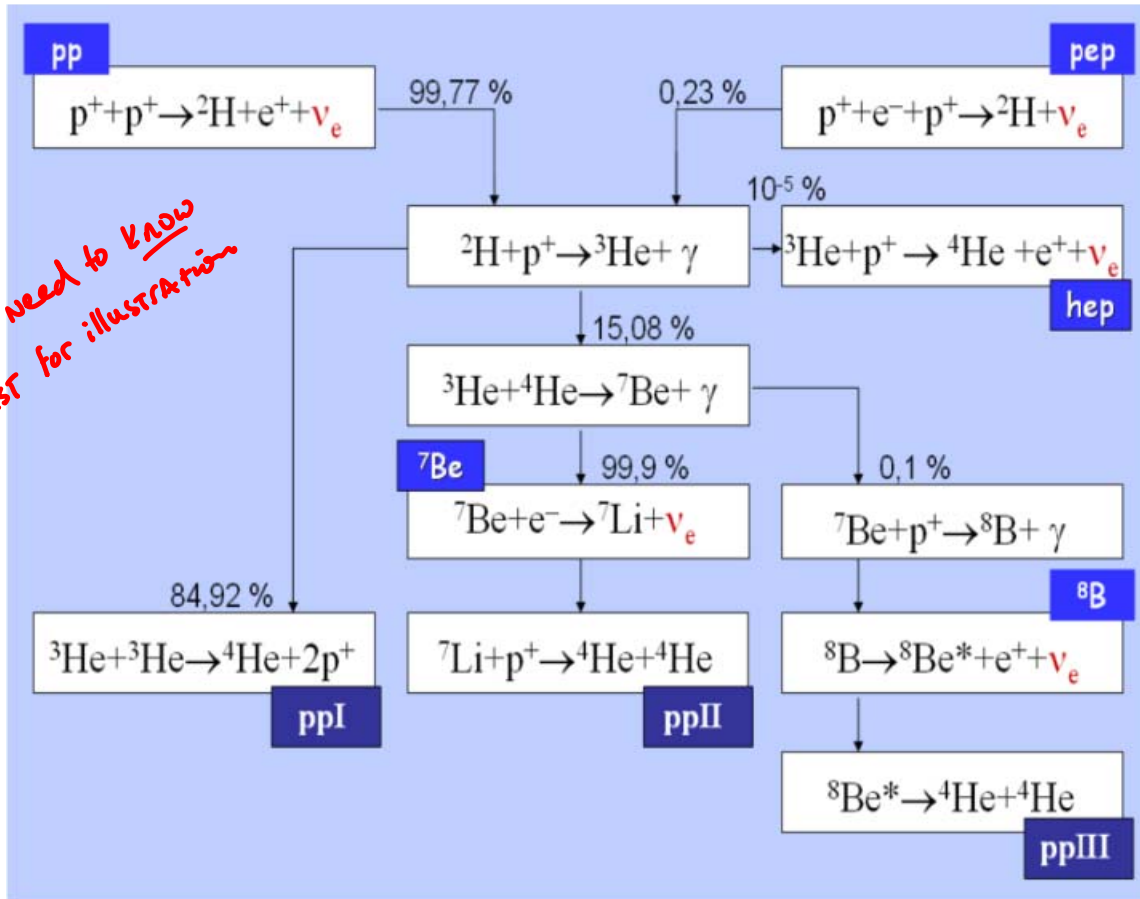
Young stars residual dust
surrounding them

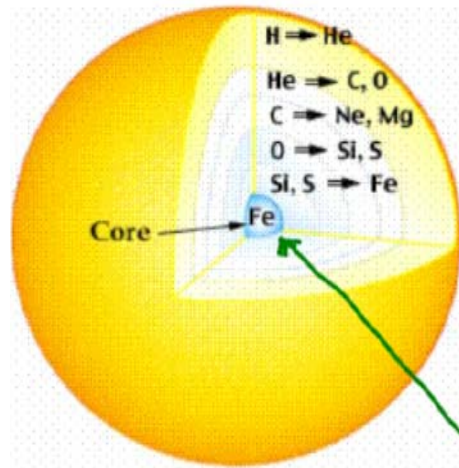


older
↑
Young

Primary Fusion Processes in the Sun

No need to know just for illustration

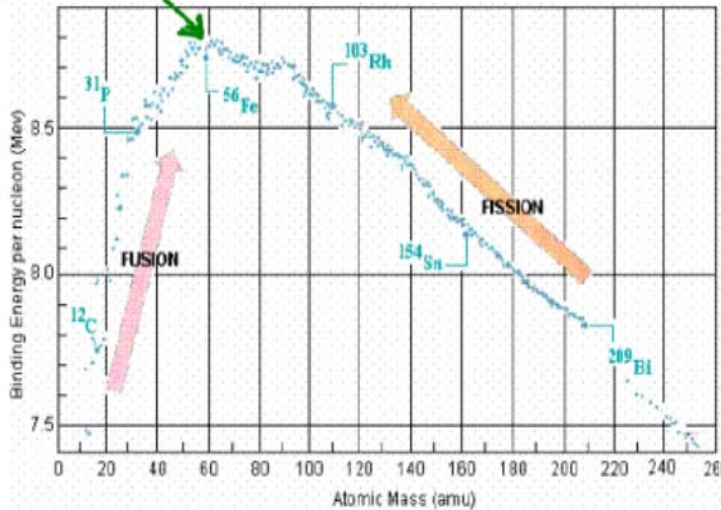


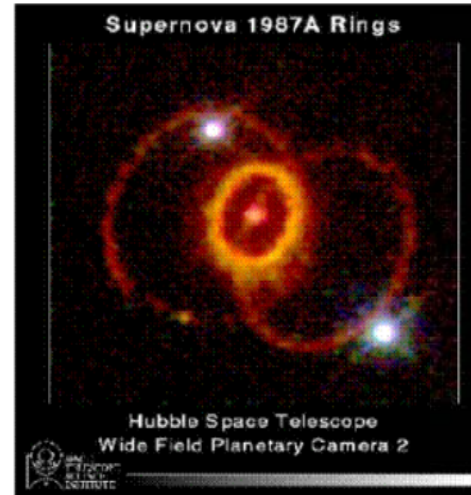
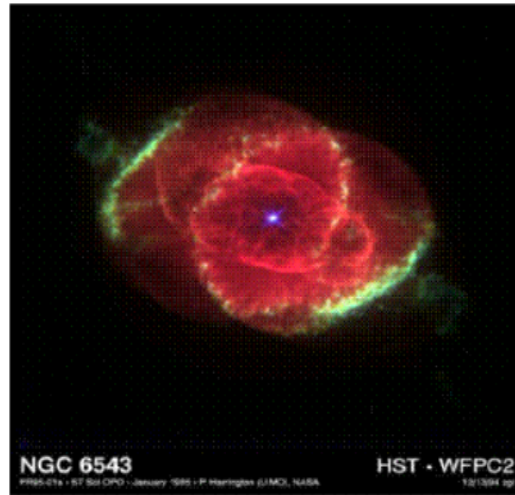


late life massive star

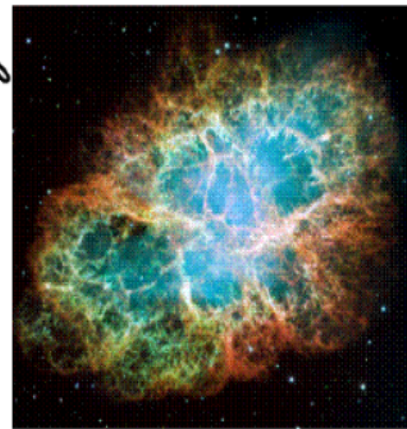
fusion process into nuclei larger than ^{56}Fe takes energy rather than releasing energy

- Early universe almost entirely Hydrogen
- normal stellar evolution \rightarrow fusion A up to ~ 57
- Supernova processes \rightarrow fusion A > 57





*STAR went Supernova in
 1054 - observed during day
 by Chinese and Arab
 Astronomers*



*Crab
 Nebula
 star went
 Supernova in
 1054*