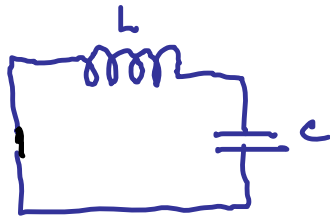


Physics 142 - November 11, 2014

■ Only 1 request for project group meeting so far ...

Last
Time

Looked at LC circuit (DC)



fully charged capacitor at $t=0$
Look at time dependence of energy flow

Find $0 = L \frac{d^2 q}{dt^2} + \frac{q}{C}$

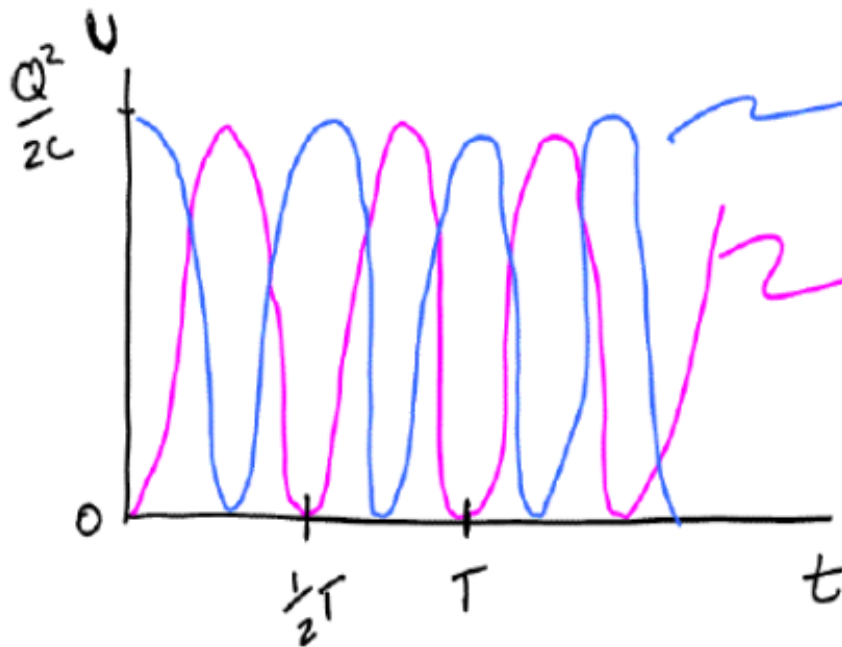
SHM for q on capacitor!

$$q(t) = Q \cos(\omega t + \phi)$$

$\omega^2 = \frac{1}{LC}$

$$i(t) = -\omega Q \sin(\omega t + \phi)$$

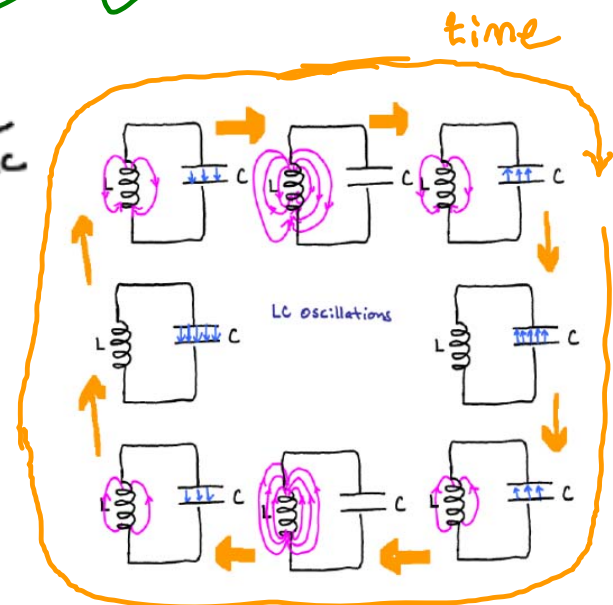
LC Oscillations



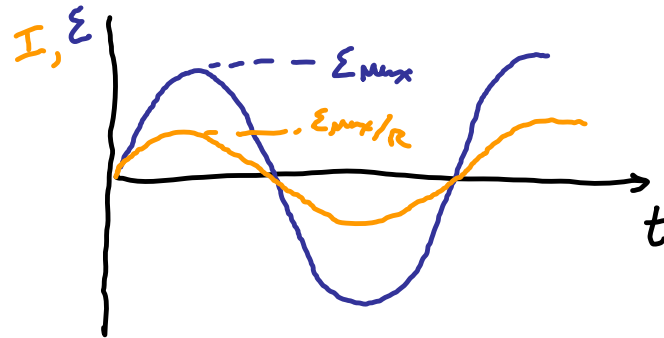
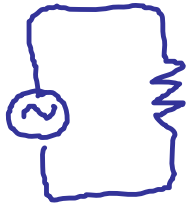
$$\frac{q^2(t)}{2C}$$

$$\frac{L i^2(t)}{2}$$

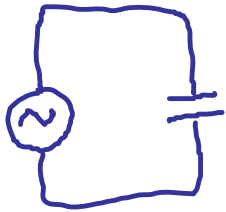
$$T = \frac{2\pi}{\omega} = 2\pi\sqrt{LC}$$



AC circuits

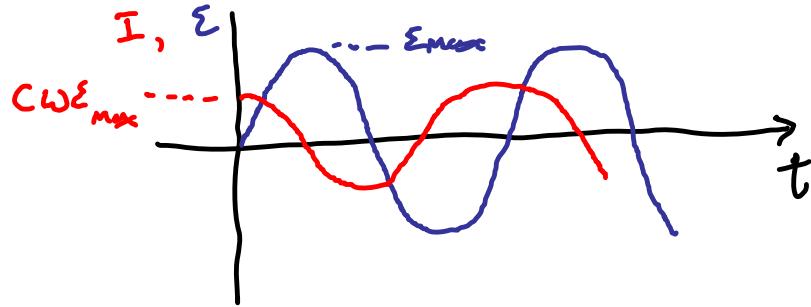


$\epsilon = IR$
 so I in phase w/ ϵ



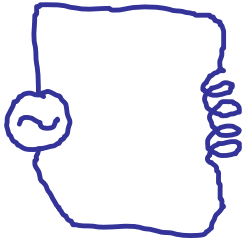
$$Q = C\epsilon = C\epsilon_{max} \sin \omega t$$

$$I = \frac{dQ}{dt} = C\omega \epsilon_{max} \cos \omega t$$



$$I \frac{1}{\omega C} \sim \epsilon$$

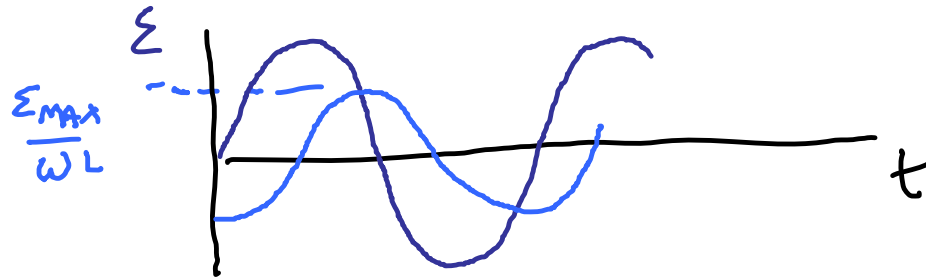
"Resistance" - Capacitive Reactance X_c



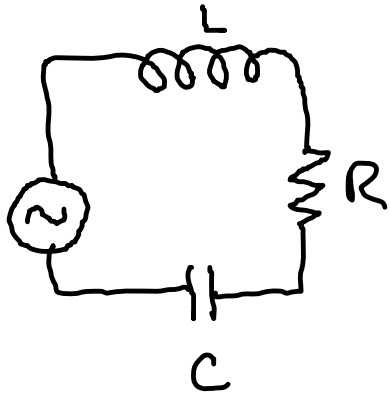
$$\mathcal{E} = -L \frac{di}{dt}$$

$$I = - \frac{\mathcal{E}_{\max}}{\omega L} \cos \omega t$$

$I \omega L \sim \mathcal{E}$
Inductive Reactance X_L



LRC circuit



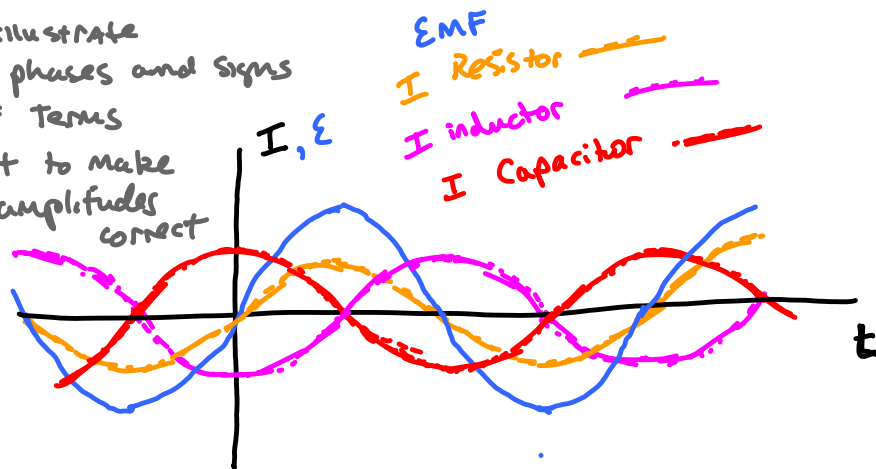
$$\mathcal{E} = \mathcal{E}_{\max} \sin \omega t$$

Assume $I = I_{\max} \sin(\omega t + \phi)$

(Note: In the original image, pink brackets and question marks are drawn under I_{\max} and ϕ respectively.)

Graph to illustrate
relative phases and signs
of EMF terms

No attempt to make
relative amplitudes
correct



$$\mathcal{E} = \underbrace{\Delta V_R}_{R I_{\max} \sin(\omega t + \phi)} + \underbrace{\Delta V_C}_{X_C I_{\max} \cos(\omega t + \phi)} + \underbrace{\Delta V_L}_{-X_L I_{\max} \cos(\omega t + \phi)}$$

often people use a graphical analysis
with I and \mathcal{E} vectors rotating
in a plane \rightarrow Phasors [see Ohanian p.1047]

I think earlier edition of Ohanian did a nice treatment
using Math + Trig ID's \rightarrow get

$$\tan \phi = \frac{\chi_L - \chi_C}{R}$$

$$I_{\max} = \frac{\Sigma_{\max}}{\sqrt{R^2 + (\chi_L - \chi_C)^2}}$$

Impedance $\equiv Z$

Plays the
role of
Resistance in
LRC circuit

⊗ → EMF has ω

what happens if $\omega \rightarrow \omega_0 = \frac{1}{\sqrt{LC}}$

$$X_L \approx \omega L$$

$$X_C = \frac{1}{\omega C}$$

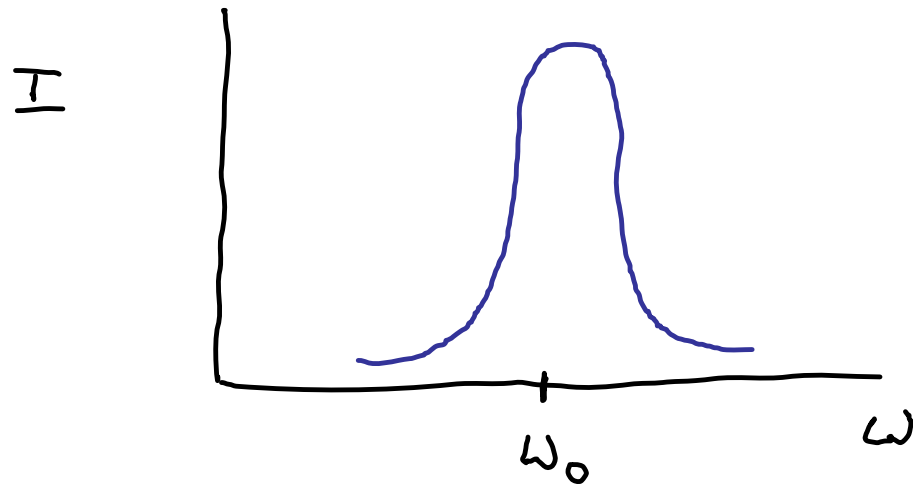
$$\sqrt{R^2 + (\omega L - \frac{1}{\omega C})^2}$$

$$\frac{1}{\sqrt{LC}}$$

$$\frac{L}{\sqrt{LC}} = \sqrt{\frac{L}{C}}$$

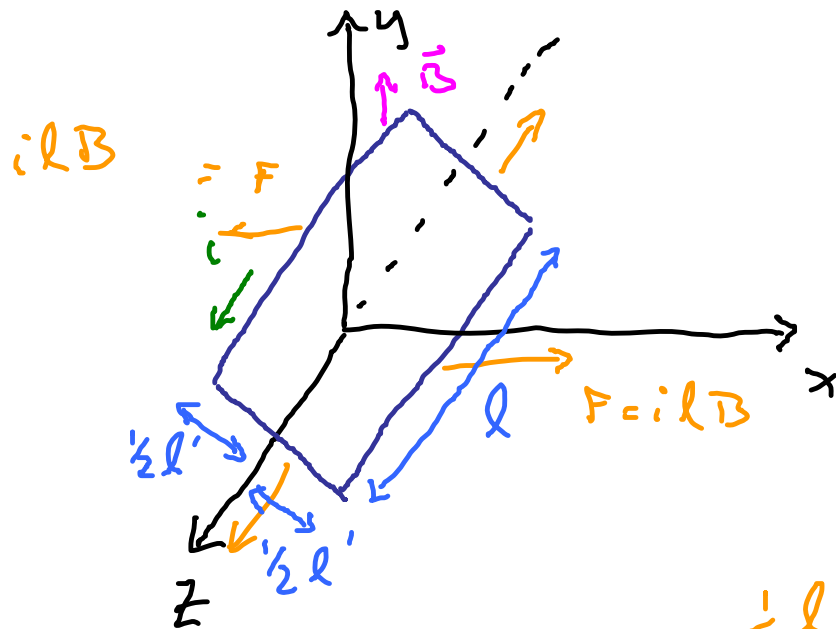
$$\frac{\sqrt{LC}}{C} = \sqrt{\frac{L}{C}}$$

→ 0 as
 $\omega \rightarrow \omega_0$

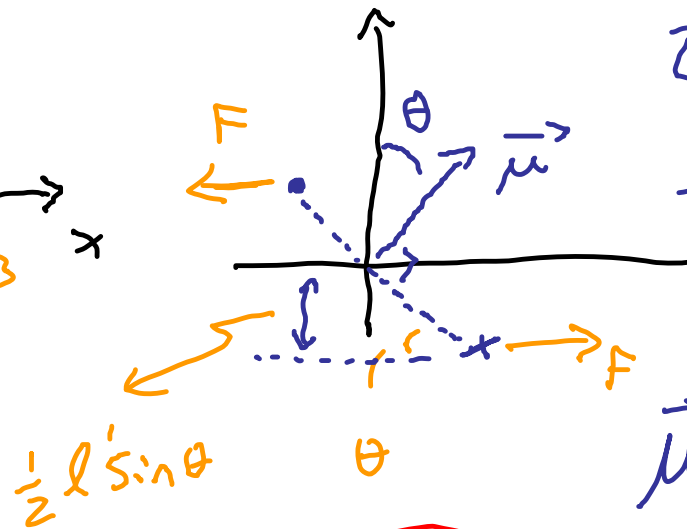


Resonance

Magnetic fields in Matter



Rectangular current loop
free to rotate at z axis



$$\tau = 2F \frac{1}{2} l' \sin \theta$$

$$\tau = i l l' \sin \theta B$$

A

$$\vec{\mu} = i A \hat{n}$$

$$\mu = i A$$

$$\tau = \vec{\mu} \times \vec{B}$$



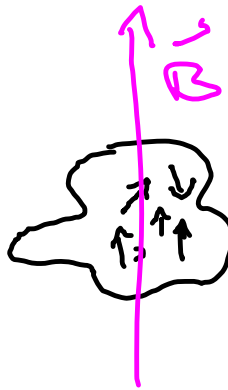
Magnetism in Materials

Paramagnetic materials



natural
Dipoles

Random
direction



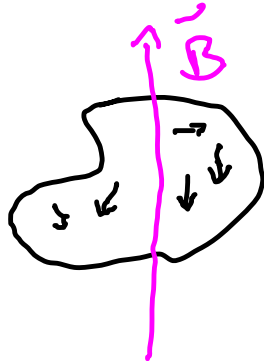
Torques
cause
Alignment

\vec{B} increased

Dia magnetic materials



NO
inherent
dipoles



induced dipoles
oppose \vec{B}

\vec{B} weakened

Ferromagnetic Materials



local domains
Strong dipoles



\vec{B} increased

$$B_{\text{in Material}} = \mu_0 (1 + \chi_m) B_{\text{external or "free"}}$$

\equiv Magnetic Susceptibility

Relative Permeability
 $\equiv \mu_r$

$\mu_0 \mu_r$
 \equiv Permeability, μ

