

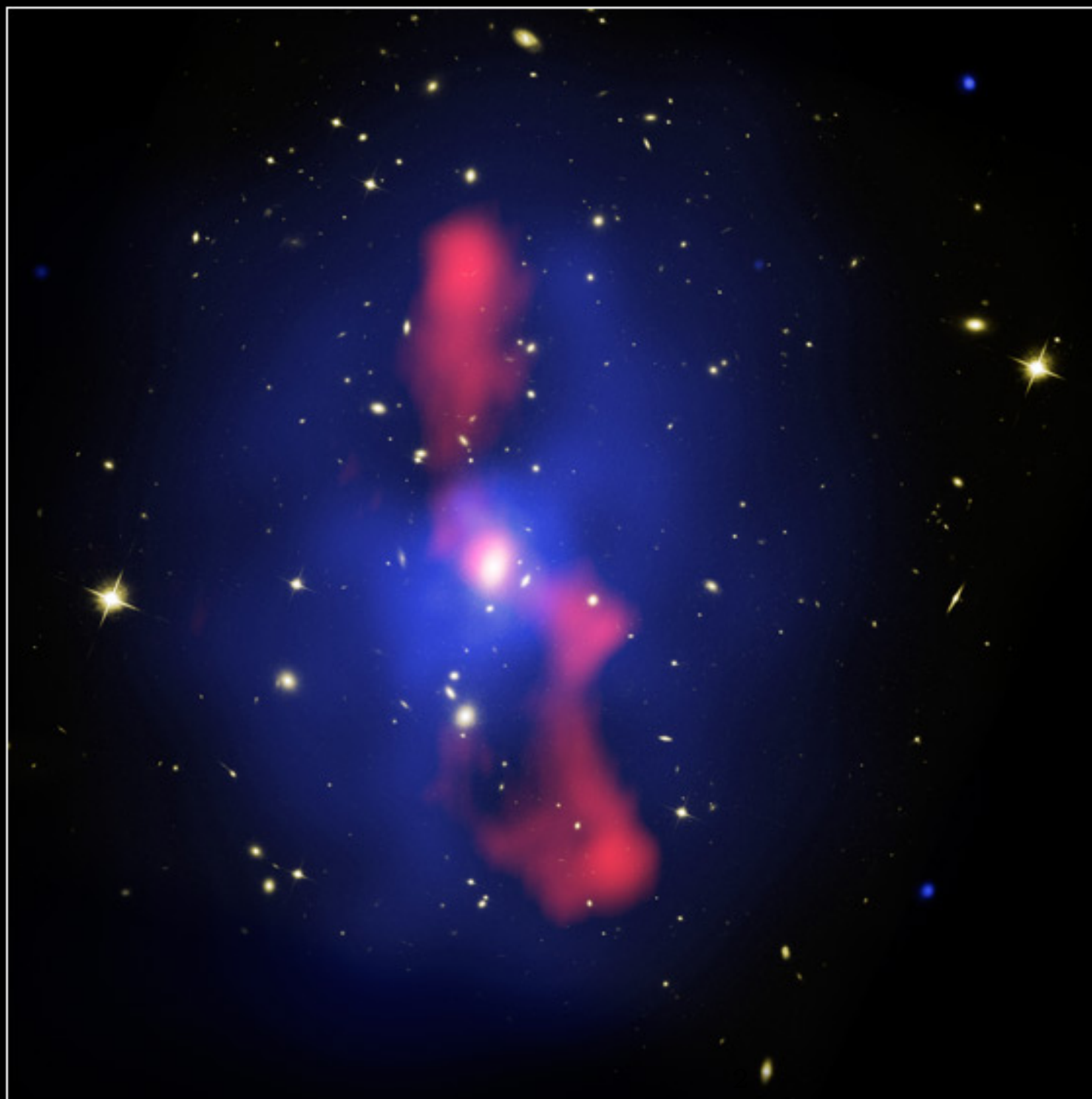
MODELING THE MAGNETIC FIELDS IN RADIO GALAXIES AND QUASARS

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University of Rochester
Department of Physics and Astronomy

15 October '09

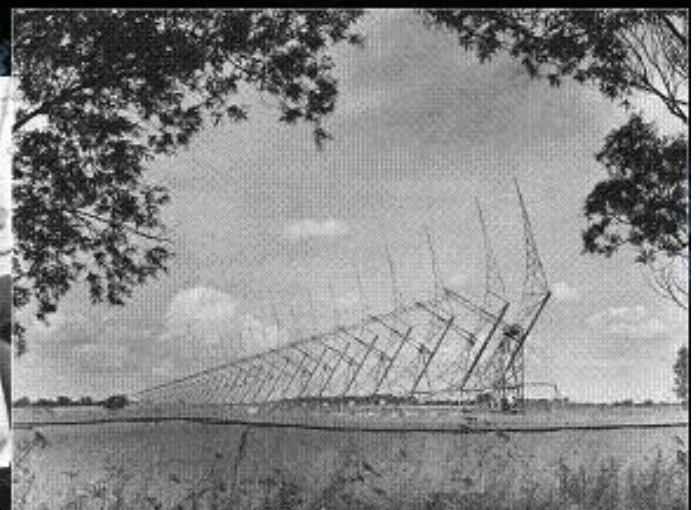
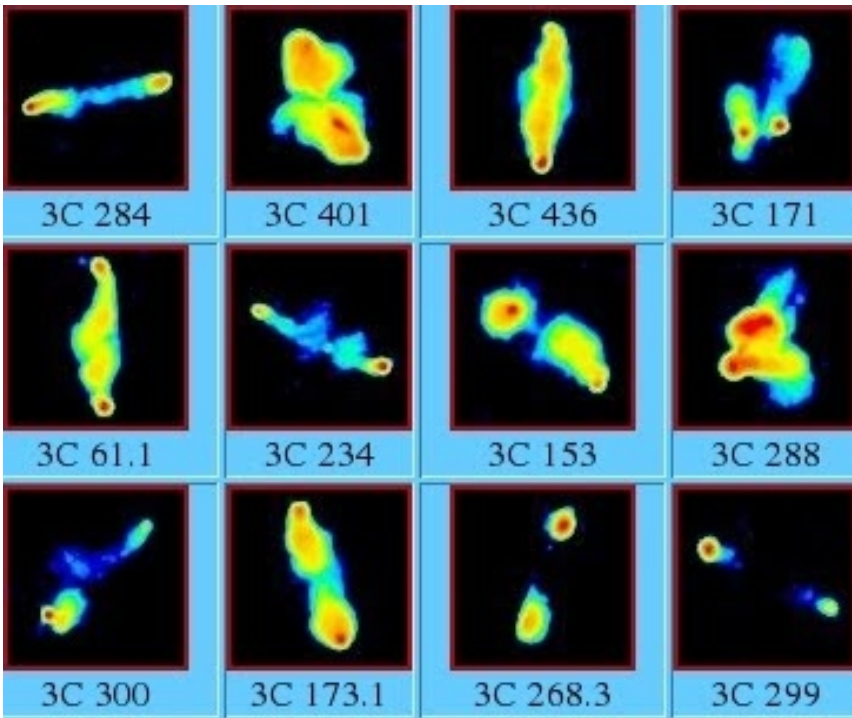


X-ray
Chandra X-Ray Observatory

Visible
Hubble Space Telescope

Radio
Very Large Array

Extragalactic radio sources



Carilli & Barthel 1996

5 GHz

Lobe

Hot Spot

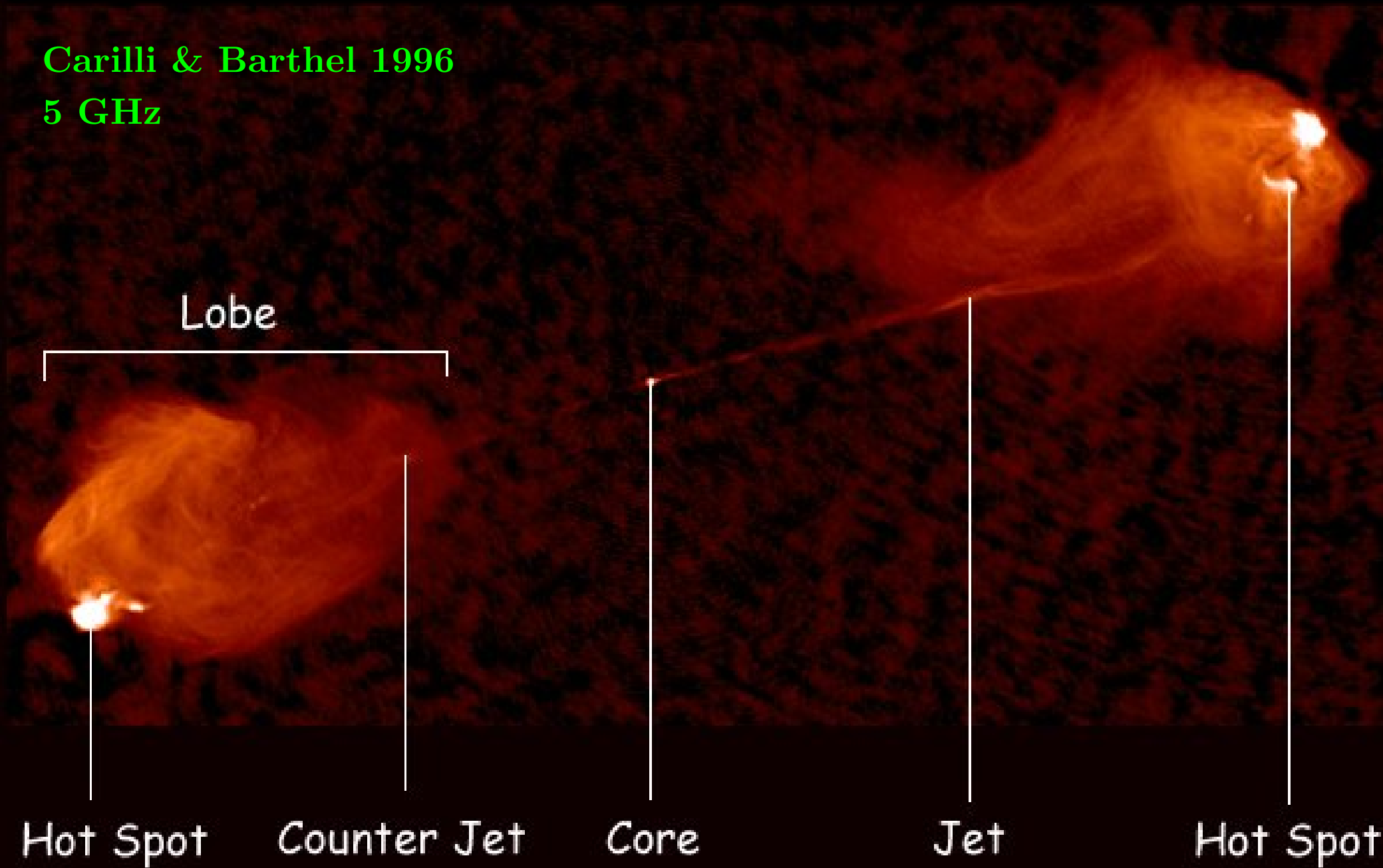
Counter Jet

Core

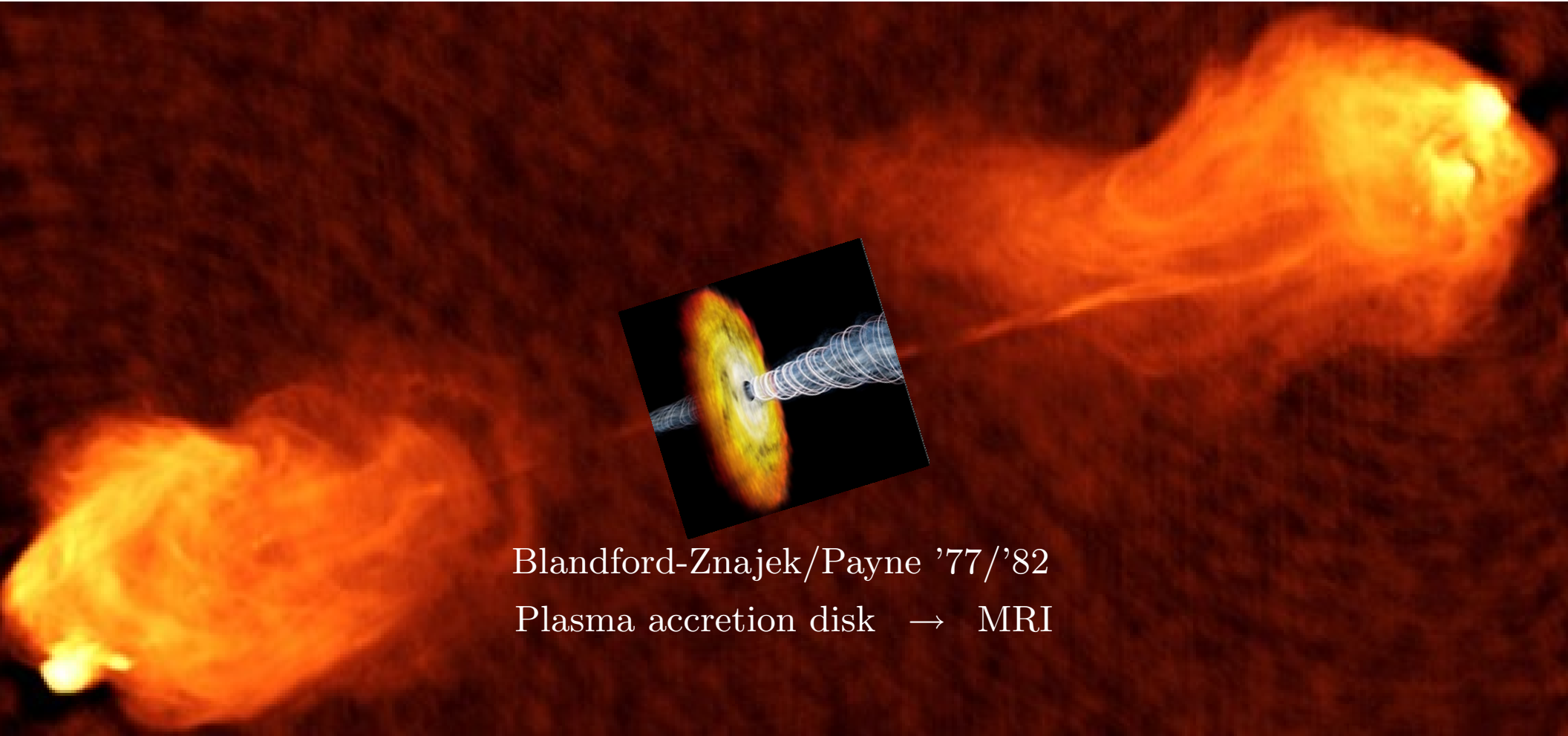
Jet

Hot Spot

Parts of a DRAGN (Cygnus A)



Formation and evolution

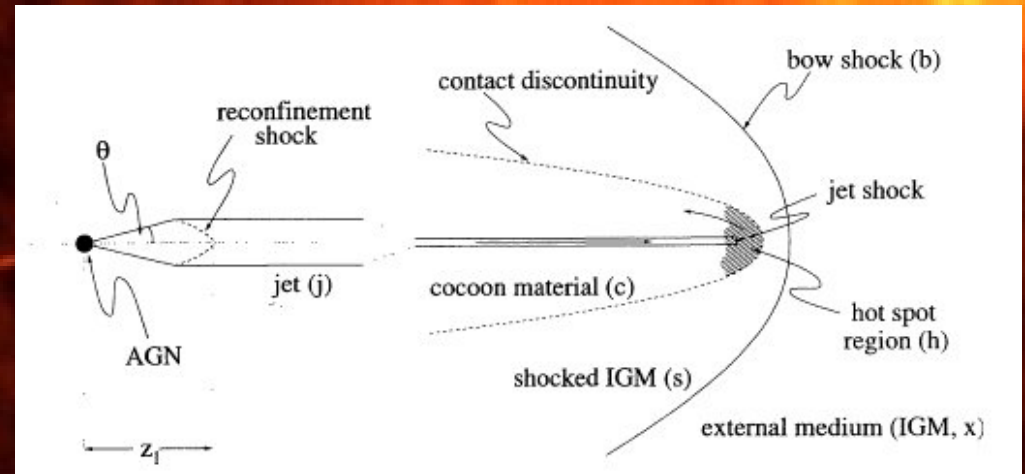


Blandford-Znajek/Payne '77/'82

Plasma accretion disk \rightarrow MRI

Formation and evolution

Sam Falle 1991

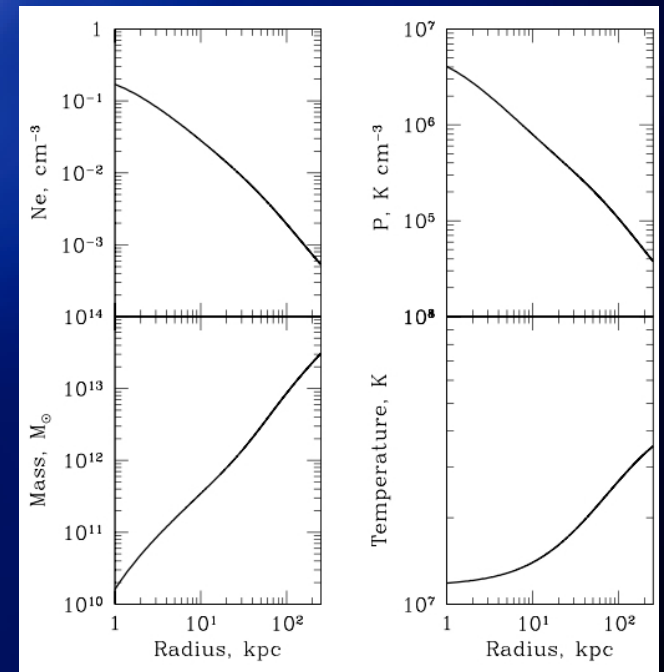


Self-similar model

Kaiser & Alexander '97

The ICM X-ray Bremsstrahlung

$$T \sim 10^8 \text{ K}$$
$$\rho \sim 10^{-2} \text{ cm}^{-3}$$



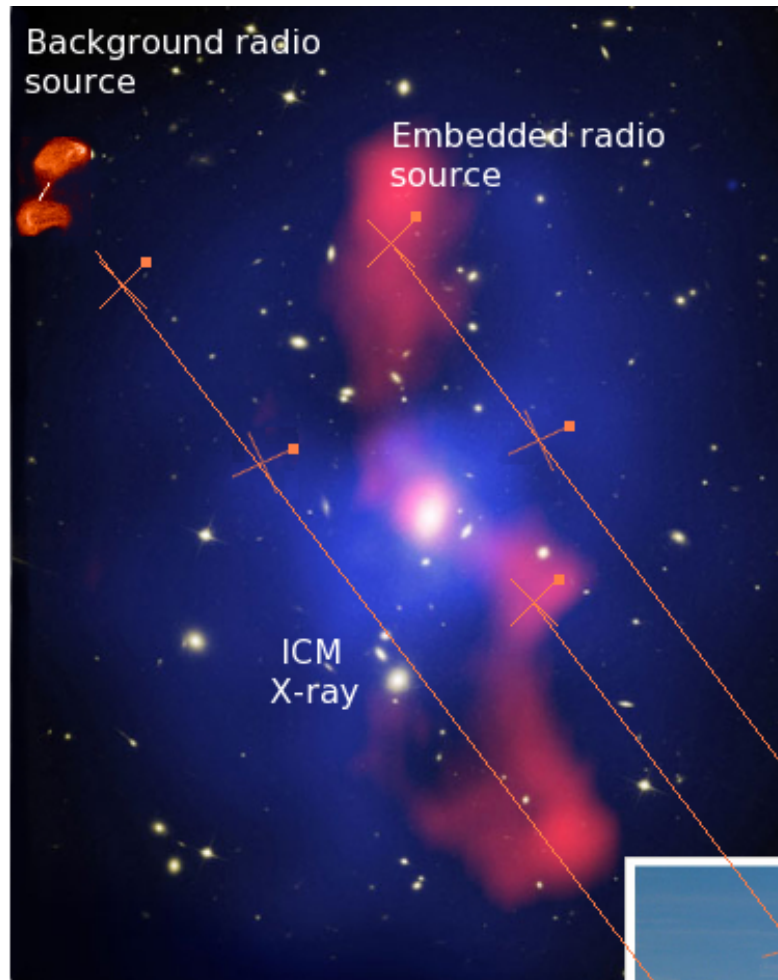
Nulsen & Böhringer 1995

NGC1275
Hydrogen-alpha

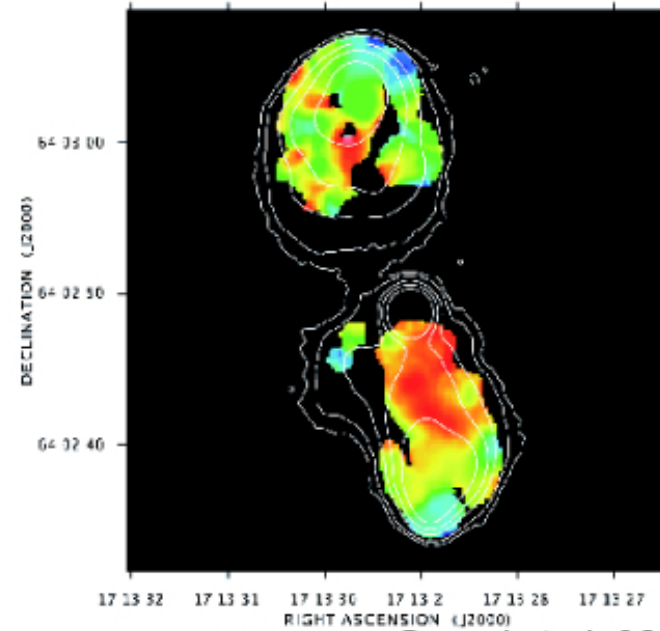
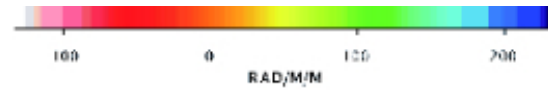
The magnetic fields in galaxy clusters



RM maps



MS 0735.6+7421 NASA composition

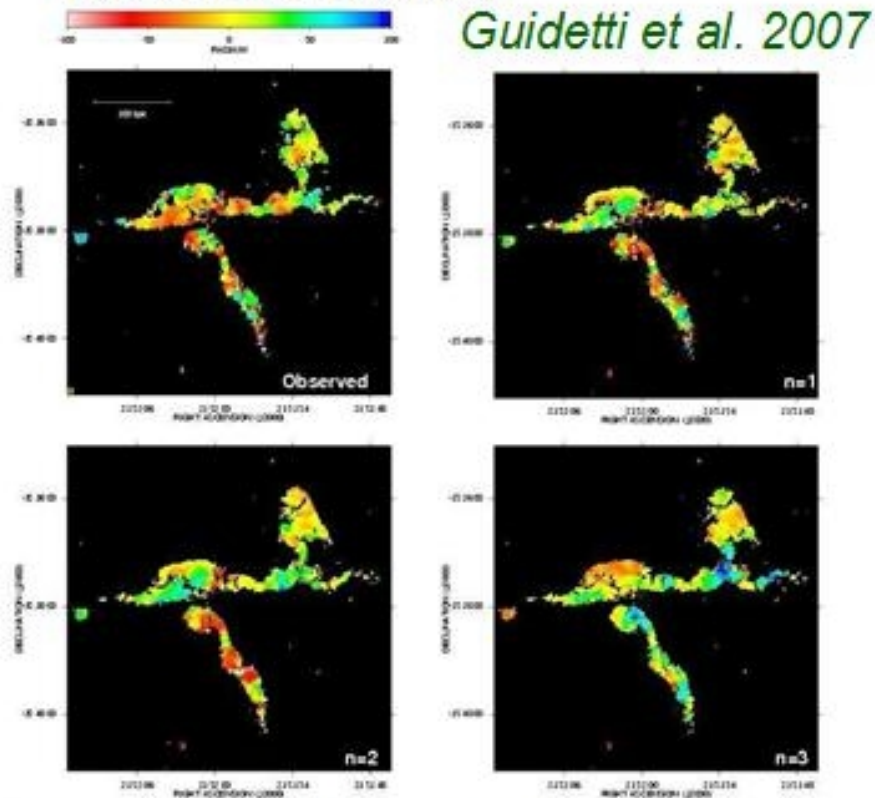


Govoni et al. 06
J1713.5+6402
Abell 2255



RM against embedded radio sources

- A23825 contains 2 extended radio galaxies
- Kolmogorov $n=11/3$ spectral index reproduces the data



Laing Garrington effect (the further radio lobe is more depolarised)

Power-spectrum of magnetic fields

Following Tribble 1991, and Murgia et al. 2004, we generate cubes with random magnetic field distributions:

Start in Fourier space. Choose a power-law power spectrum: $|A_k|^2 \propto k^{-\zeta}$, where ζ takes the values 2, 3, 11/3 (Kolmogorov) or 4.

Draw the amplitudes A and phases ϕ of the magnetic potential vector $\tilde{\mathbf{A}}(\mathbf{k})$ randomly from: $P(A, \phi)dAd\phi = \frac{A}{2\pi|A_k|^2} \exp\left(-\frac{A^2}{2|A_k|^2}\right) dAd\phi$,

Eliminate mean field, i.e., $A(k=0) = 0$,

$$\tilde{\mathbf{A}}(\mathbf{k}) \rightarrow \mathbf{A}(\mathbf{x}),$$

$$\mathbf{A}'(\mathbf{x}) = (\mathbf{A}(\mathbf{x}) * \rho(r)_{ICM})(\mathbf{x}), \text{ where } \rho_{ICM}(r) \propto \frac{1}{1+(r/r_c)^2},$$

$\mathbf{B}(\mathbf{x}) = \nabla \times \mathbf{A}'(\mathbf{x})$, and not the other way around due to noise and

$\mathbf{B}(\mathbf{x})' = B_0 \mathbf{B}(\mathbf{x})$, so $\beta \sim 10$ everywhere.

Ideal MHD to follow the evolution

$$\begin{aligned}
 & \frac{\partial}{\partial t} \begin{bmatrix} \rho \\ \rho v_x \\ \rho v_y \\ \rho v_z \\ \epsilon \\ B_x \\ B_y \\ B_z \end{bmatrix} + \frac{\partial}{\partial x} \begin{bmatrix} \rho v_x \\ \rho v_x^2 + P + B^2/2 - B_x^2 \\ \rho v_y v_x \\ \rho v_z v_x \\ (\epsilon + P + B^2/2)v_x - B_x(\vec{B} \cdot \vec{v}) \\ 0 \\ -E_z \\ E_y \end{bmatrix} + \\
 & \frac{\partial}{\partial y} \begin{bmatrix} \rho v_y \\ \rho v_x v_y \\ \rho v_y^2 + P + B^2/2 - B_y^2 \\ \rho v_z v_y \\ (\epsilon + P + B^2/2)v_y - B_y(\vec{B} \cdot \vec{v}) \\ E_z \\ 0 \\ -E_x \end{bmatrix} + \frac{\partial}{\partial z} \begin{bmatrix} \rho v_z \\ \rho v_x v_z \\ \rho v_y v_z \\ \rho v_z^2 + P + B^2/2 - B_z^2 \\ (\epsilon + P + B^2/2)v_z - B_z(\vec{B} \cdot \vec{v}) \\ -E_y \\ E_x \\ 0 \end{bmatrix} = \vec{S}
 \end{aligned}$$

Faraday's law:

$$\frac{\partial}{\partial t} \vec{B} + \vec{\nabla} \times \vec{E} = 0, \quad (1)$$

hence, initially solenoidal magnetic field topologies remain solenoidal:

$$\vec{\nabla} \cdot \vec{B} = 0.$$

Ohm's law for a perfectly conducting medium,

$$\vec{E} = -\vec{v} \times \vec{B} \quad (2)$$

and the polytropic equation of state for an ideal gas,

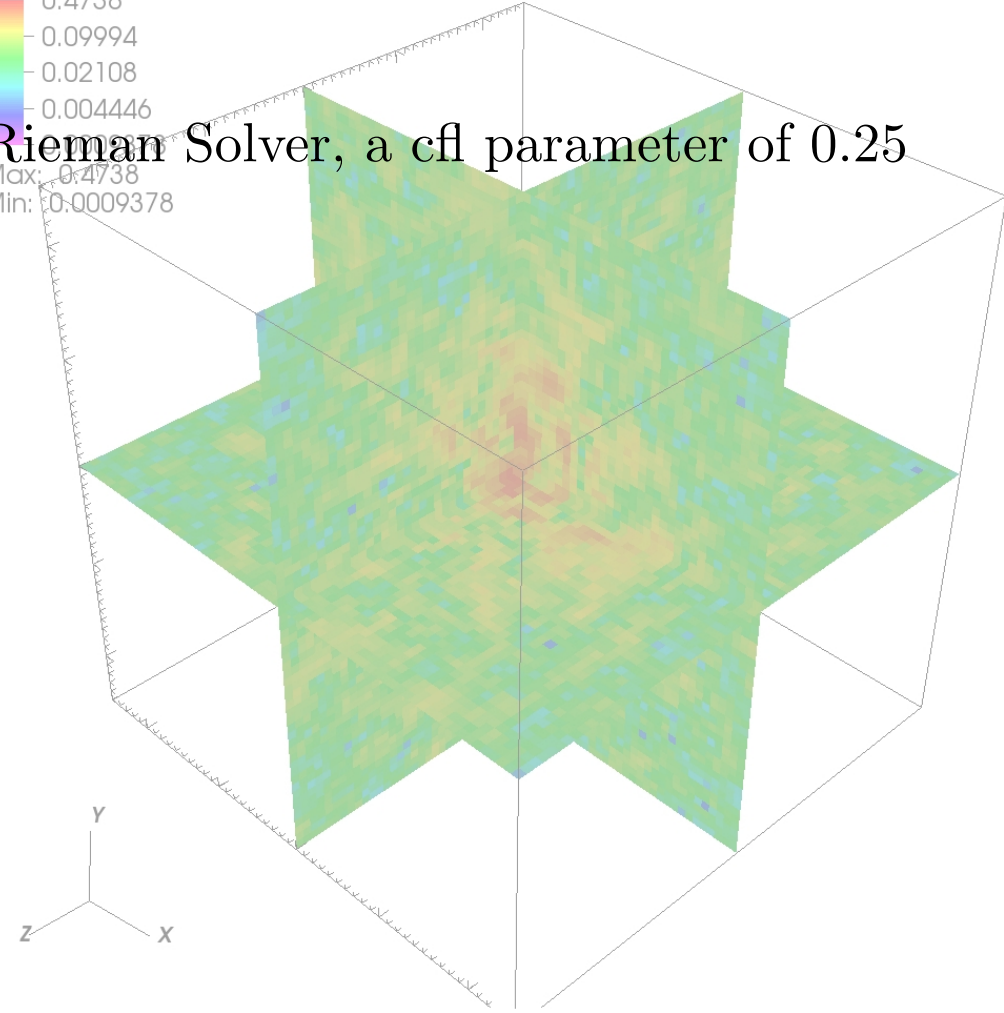
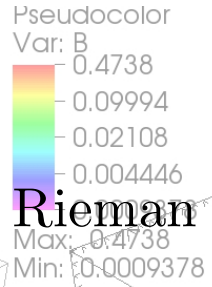
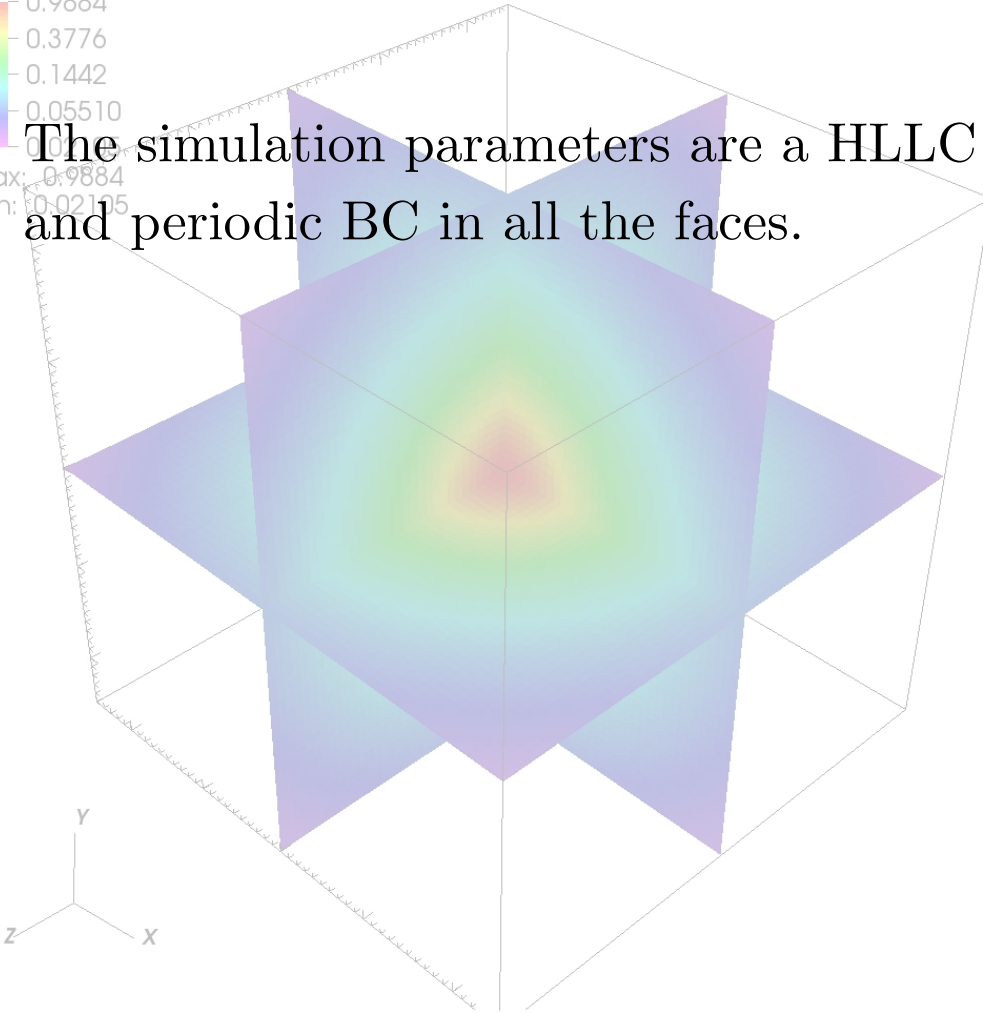
$$P = (\gamma - 1)(\epsilon - \rho \vec{v}^2 / 2 - B^2 / 2). \quad (3)$$

Magnetic flux freezing

$$\frac{\partial}{\partial t} \left(\int_{S(t)} \mathbf{B}(\mathbf{r}, t) \cdot d\mathbf{a} \right) = 0 \quad (4)$$

We use Flash3.1 (Fryxell et al. 2000) which is an MHD, multidimensional, AMR and MPI based code.

The simulation parameters are a HLLC Riemann Solver, a cfl parameter of 0.25 and periodic BC in all the faces.



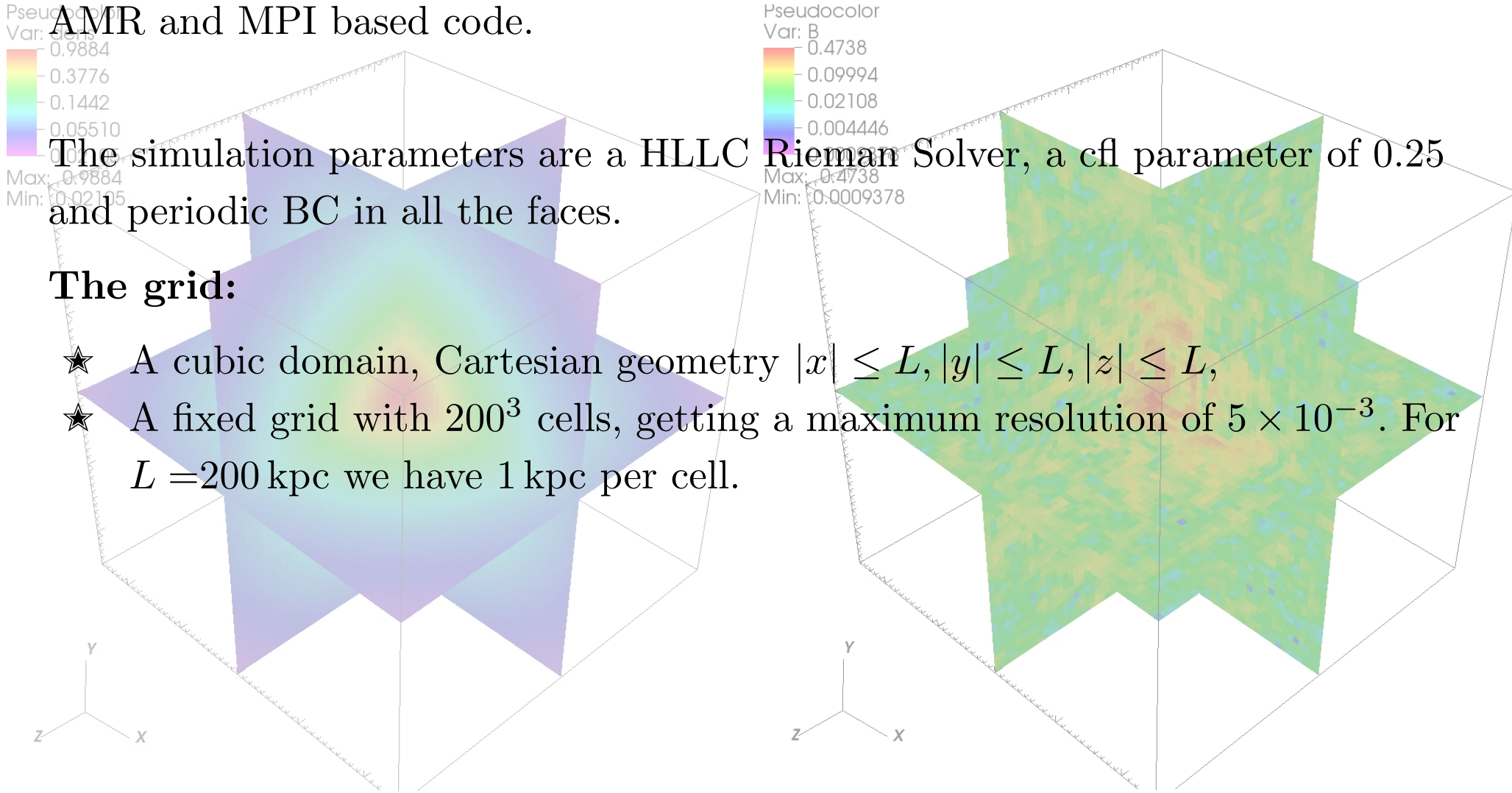
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The grid:

- ★ A cubic domain, Cartesian geometry $|x| \leq L, |y| \leq L, |z| \leq L$,
- ★ A fixed grid with 200^3 cells, getting a maximum resolution of 5×10^{-3} . For $L = 200$ kpc we have 1 kpc per cell.



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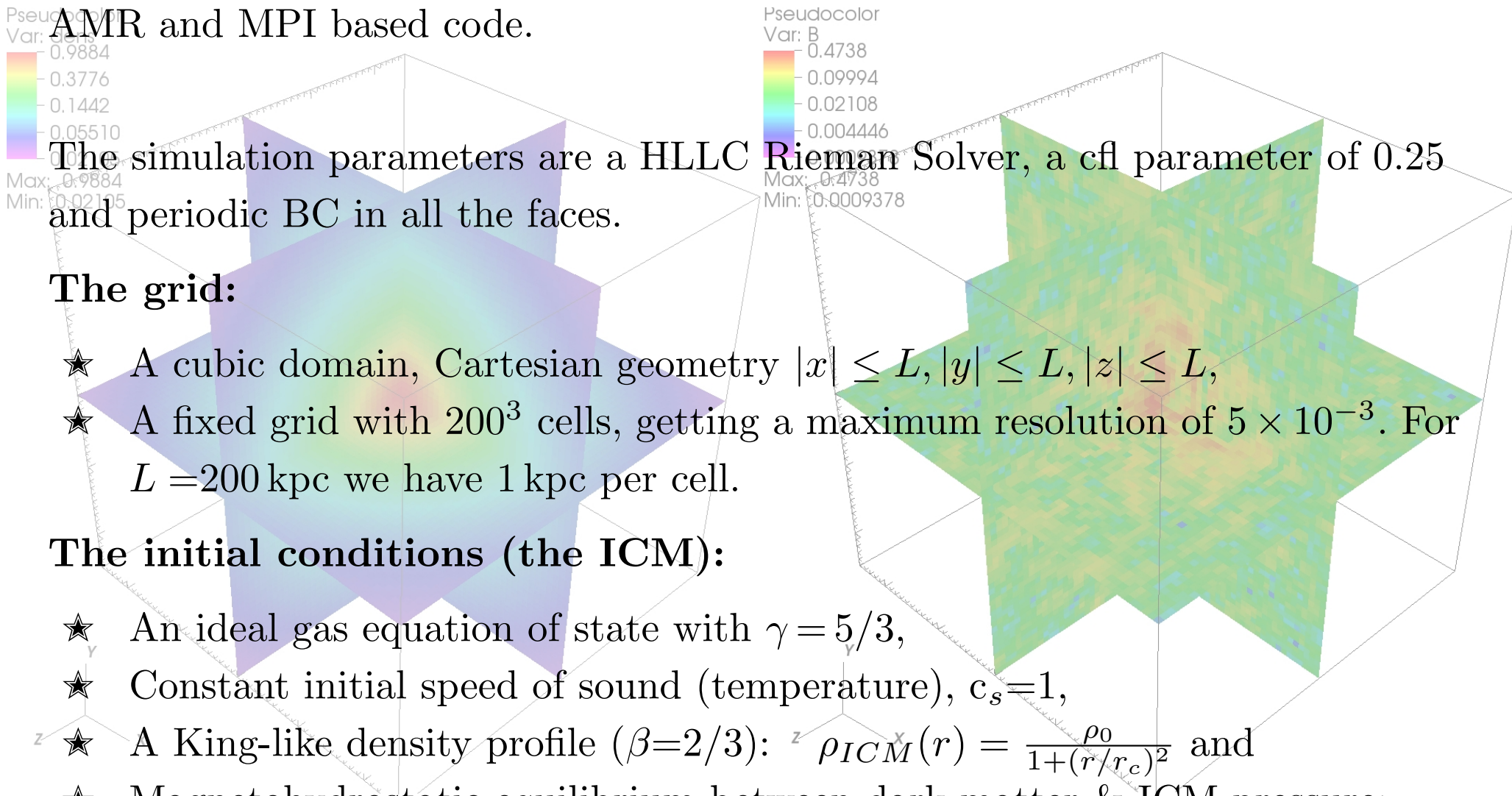
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The initial conditions (the ICM):

- ★ An ideal gas equation of state with $\gamma = 5/3$,
- ★ Constant initial speed of sound (temperature), $c_s = 1$,
- ★ A King-like density profile ($\beta = 2/3$): $\rho_{ICM}(r) = \frac{\rho_0}{1+(r/r_c)^2}$ and
- ★ Magneto hydrostatic equilibrium between dark matter & ICM pressure:

$$g_r = -\frac{2c_s^2}{\gamma a_0^2} \frac{r}{(1+(r/a_0)^2)} \left(1 + \frac{1}{\beta_m}\right).$$



The Jets

We inject mass and energy into a central cylinder in the grid:

$$\begin{aligned}\rho_{x,\text{nozzle}}^{n+1} &= \rho_{x,\text{nozzle}}^n + \dot{\rho}_j dt \\ &= \rho_{x,\text{nozzle}}^n + \delta(\mathbf{x}, t) \frac{\eta\rho_0}{t_m} dt, \\ v_{x,\text{nozzle}}^{n+1} &= v_{x,\text{nozzle}}^n + v_j dt \\ &= v_{x,\text{nozzle}}^n + \delta(\mathbf{x}, t) \frac{v_j}{t_k} dt,\end{aligned}$$

respectively, where

$$\delta(\mathbf{x}, t) = \begin{cases} 1 & \text{for } |\mathbf{x}| \leq l_j \text{ and } t \in (t_s, t_e); \\ 0 & \text{otherwise,} \end{cases}$$

η is the jets vs ambient density contrast, v_j is the jets' velocity, t_m and t_k are the mass and energy injection rates, respectively, and t_j characterises the duration of the injection.

Parameter space:

Jets velocity [Mach]	Jets vs ambient density contrast	Jets power [W]
40	0.001	3.8×10^{38}
40	0.01	1.6×10^{39}
80	0.001	1.5×10^{39}
80	0.01	1.1×10^{40}

$$t_{\text{simul}} \propto \frac{1}{n_{\text{processors}}} \frac{\Delta x}{\vec{v}_{\text{relevant}}} + t_{\text{queue}} \sim 4\text{--}12 \text{ hours} \quad \text{with } 64\text{--}128 \text{ proc.}$$

For the analysis, rescale to units in the high energy extragalactic context:

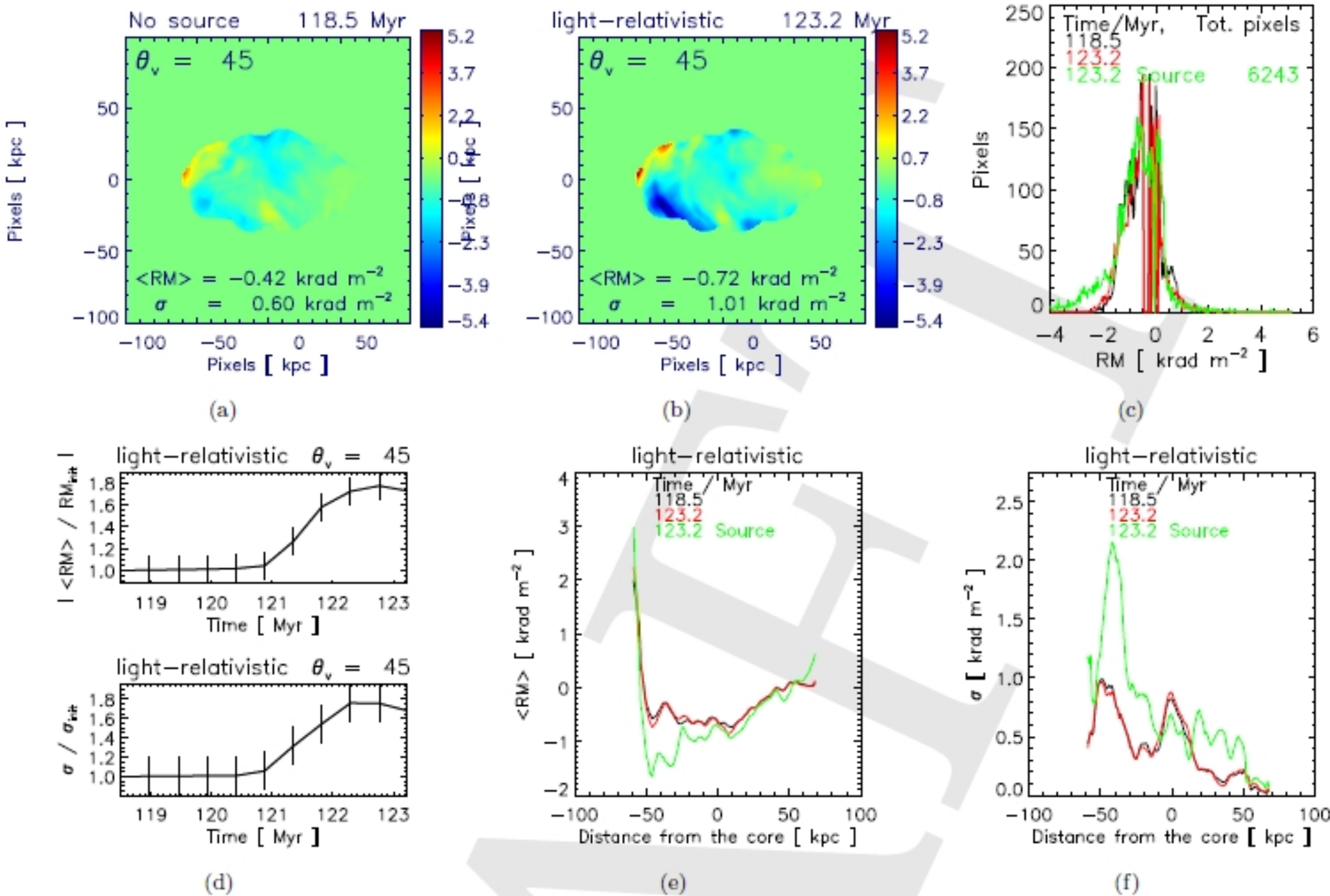
$$\rho_0 = 1 \times 10^{-2} \text{ cm}^{-3},$$

$$L = 200 \text{ kpc},$$

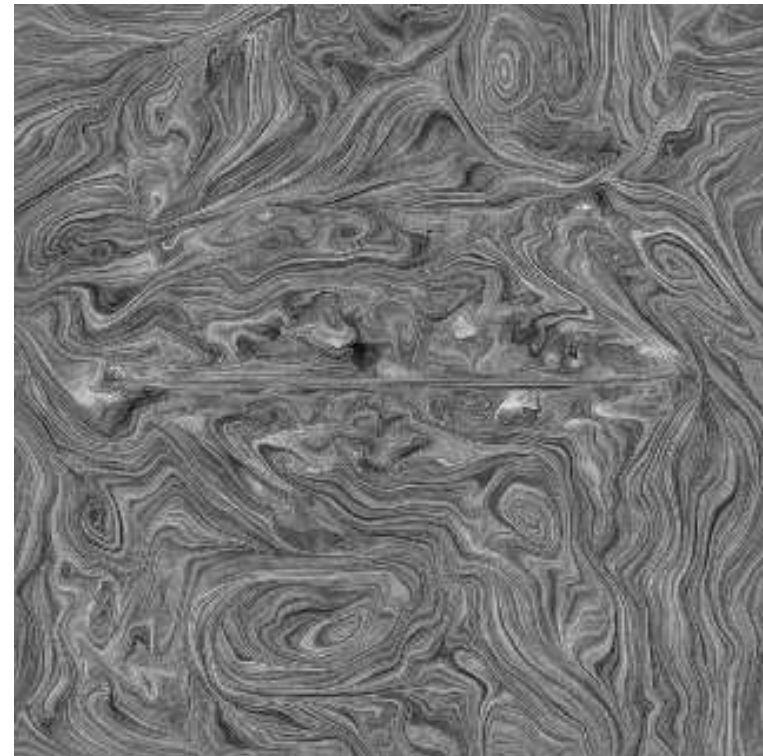
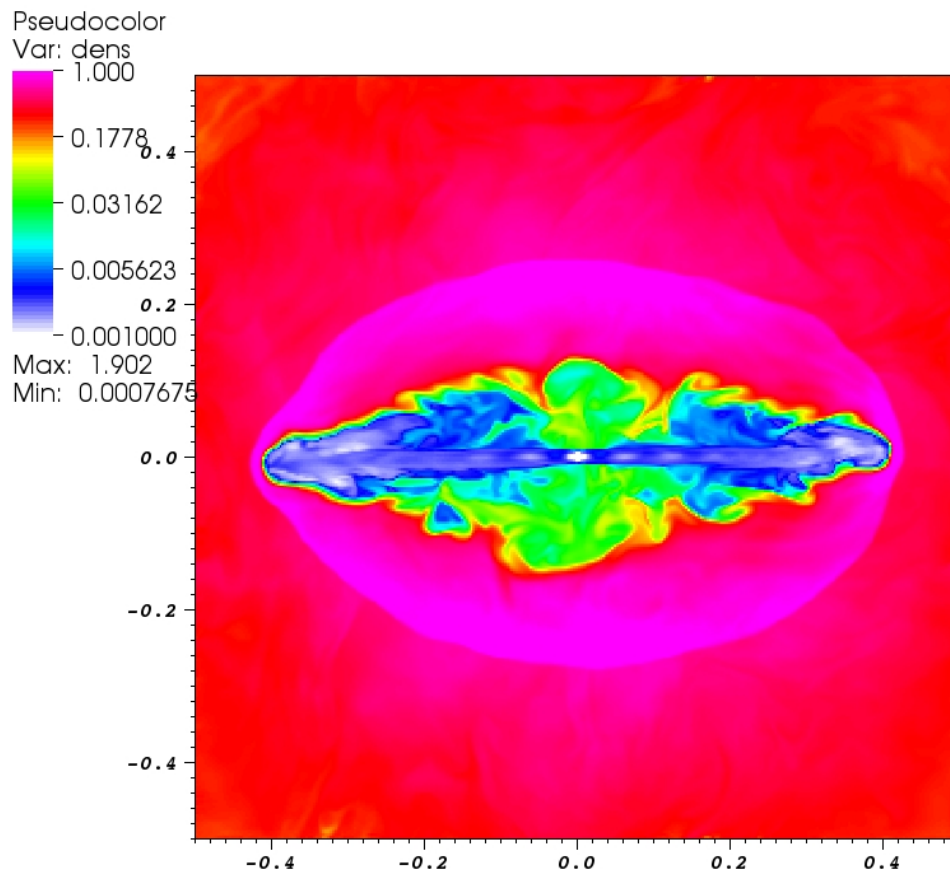
$$\langle \mathbf{B} \rangle \in (1, 100) \mu\text{Gauss and}$$

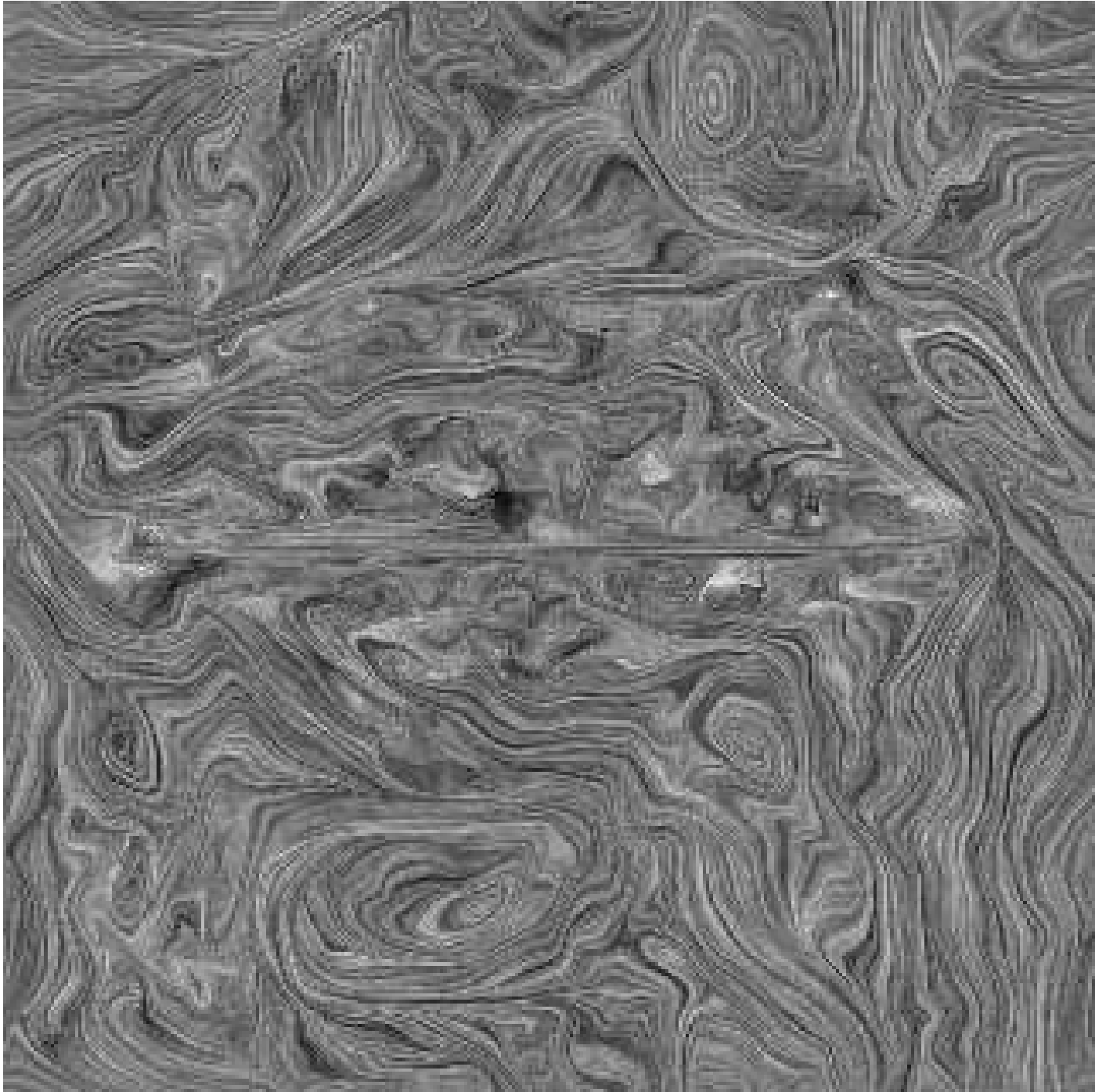
$$c_s \sim 1.6 \times 10^8 \text{ cm s}^{-1}, \text{ for a plasma temperature of } \sim 10^8 \text{ K.}$$

Synthetic RM maps of the AGN feedback on the cluster magnetic fields:



The topology of the magnetic fields lines 12 Myr after the relaxation:





Synthetic maps of non-thermal polarized synchrotron emission.

Data cubes \longrightarrow Stokes parameters (for each cell):

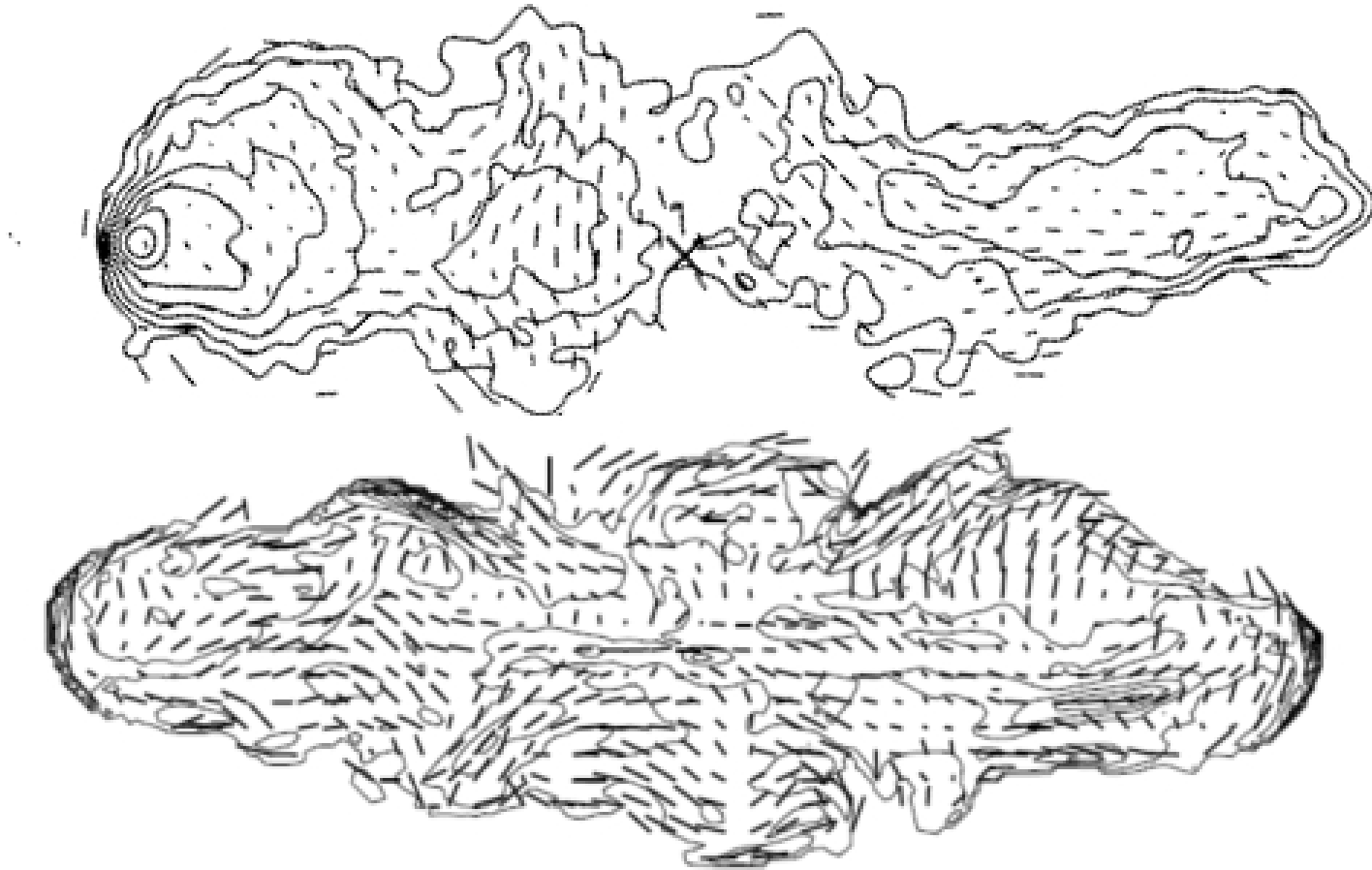
$$I = \langle E_x^2 + E_y^2 \rangle, \quad Q = \langle E_x^2 - E_y^2 \rangle, \quad U = \langle 2E_x E_y \cos(\psi_x - \psi_y) \rangle \quad (5)$$

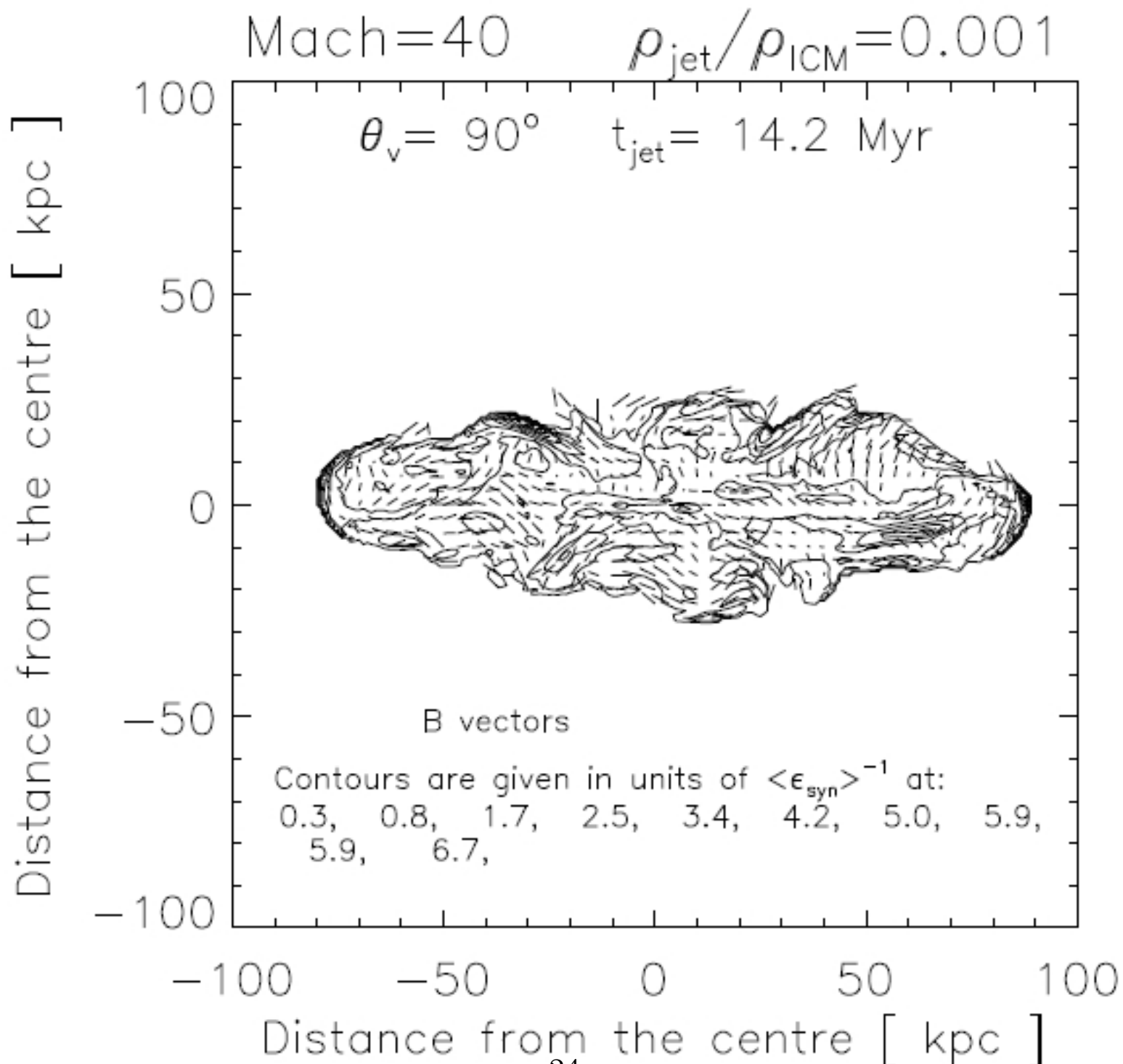
where $\langle \rangle$ denotes ensemble-averaging through the source, along the line of sight, and I represents the polarisation intensity of the radiation.

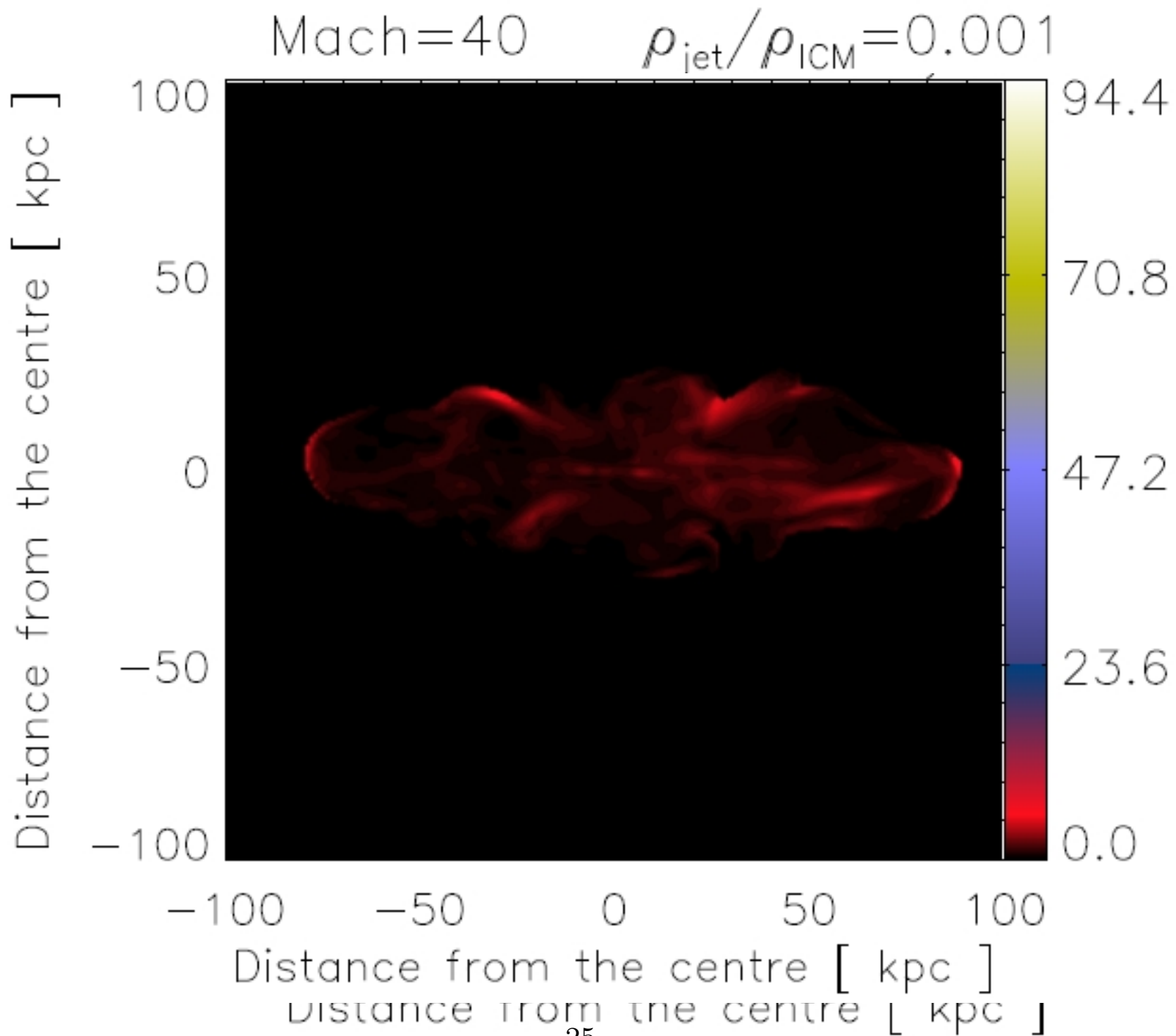
In terms of these parameters:

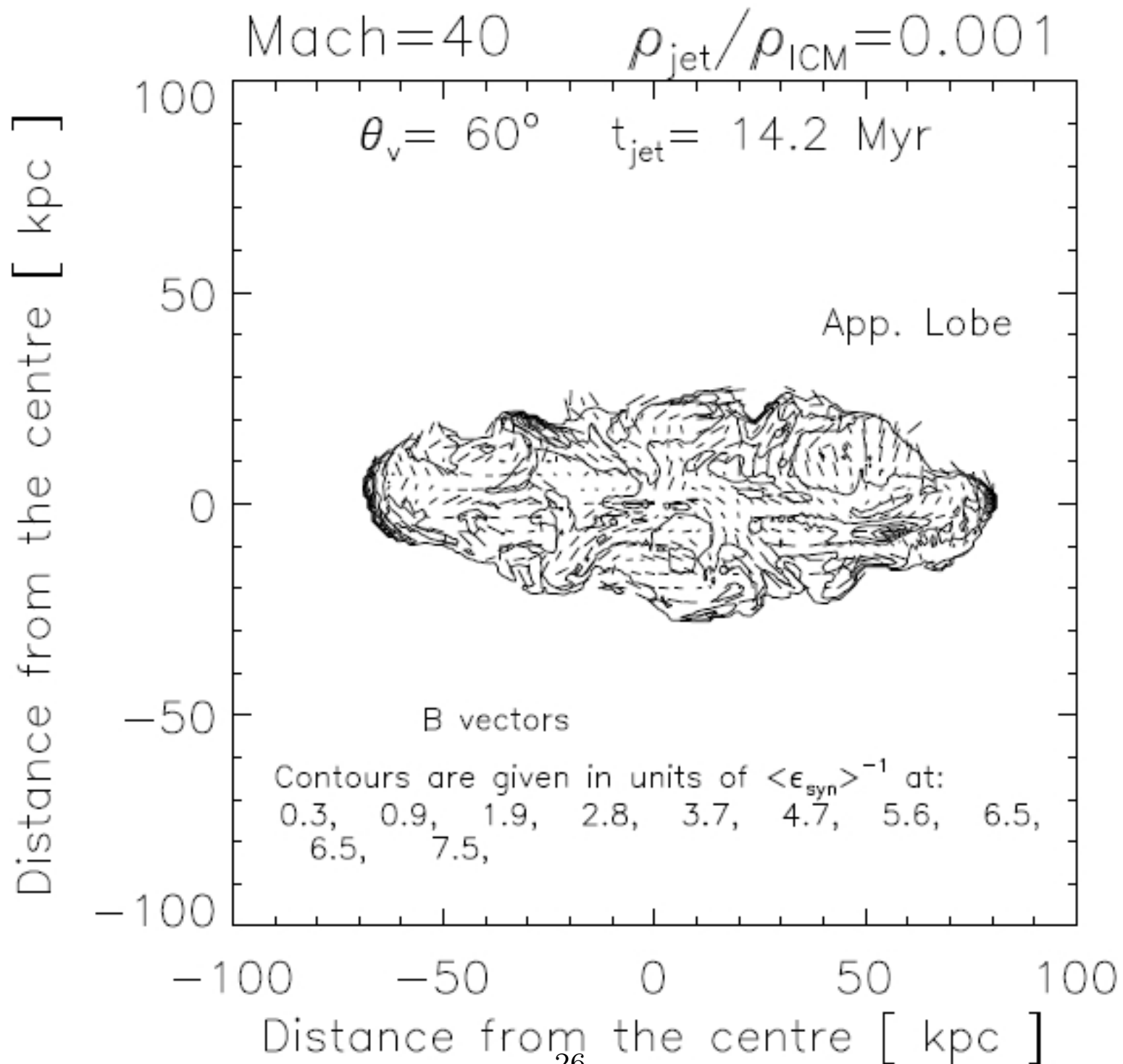
The polarisation angle $\chi_B = \arctan(U/Q)/2 + \pi/2$ and the degree of polarisation $d_p = \sqrt{U^2 + Q^2}/I$.

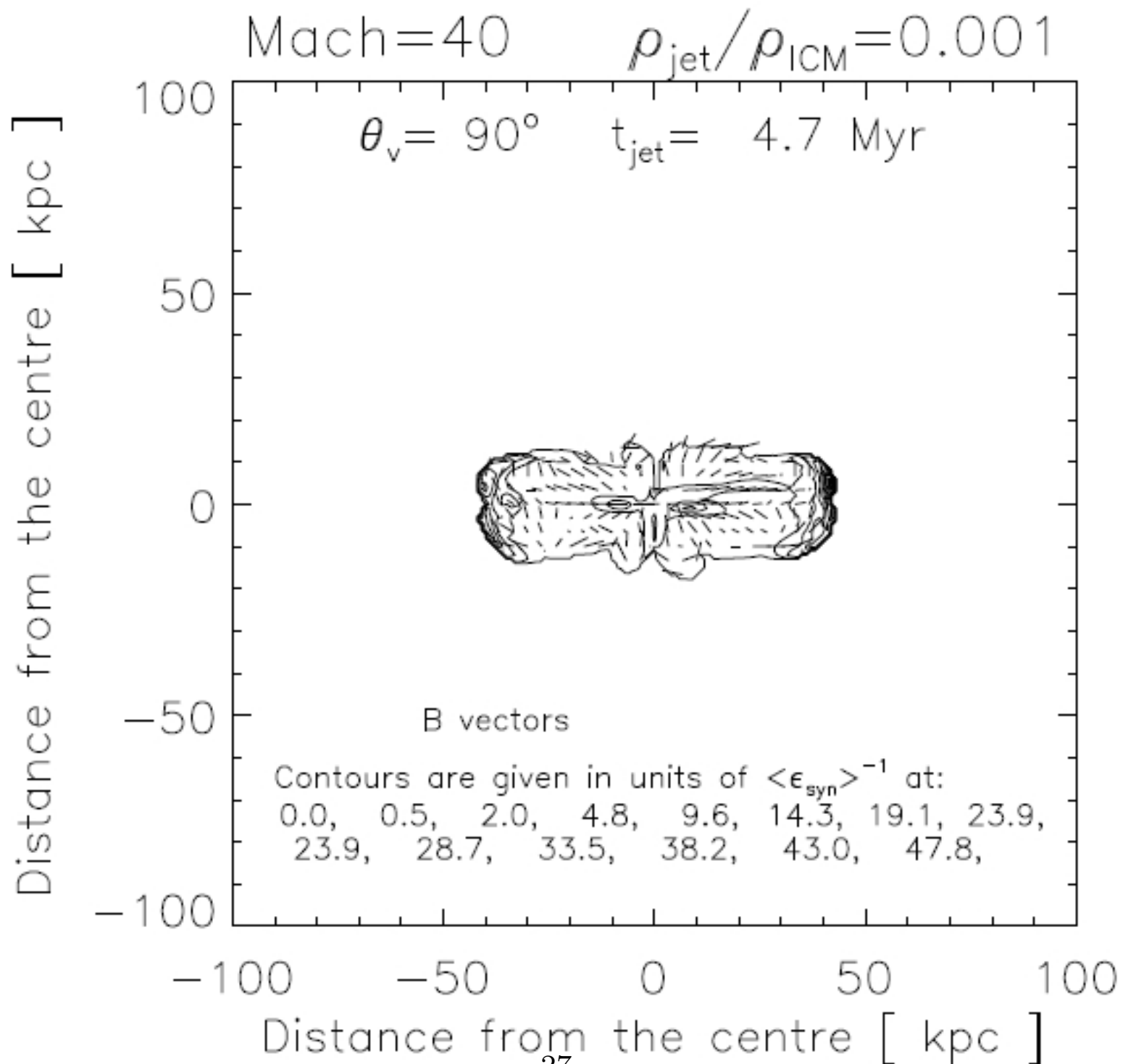
Which one is the observation?





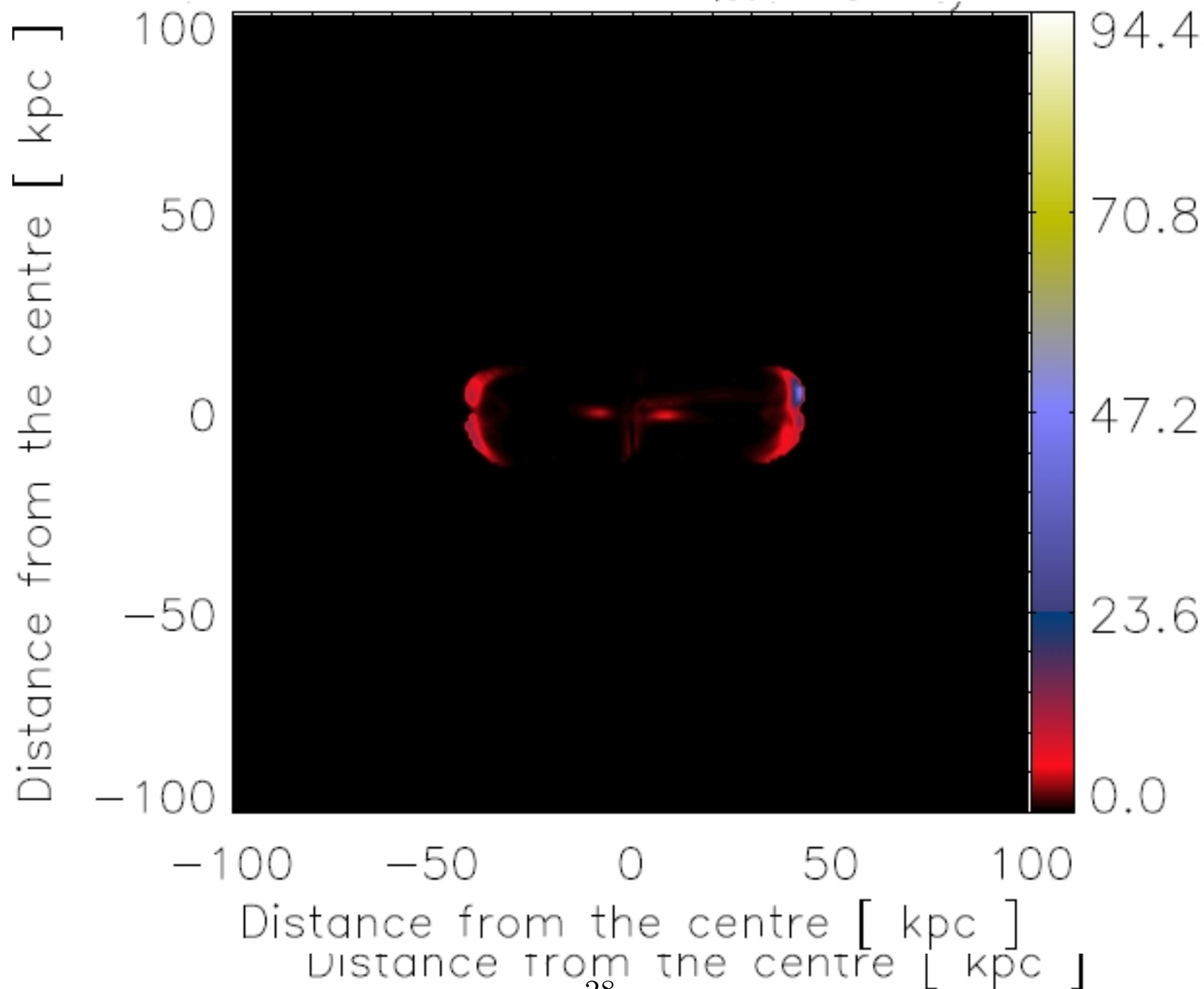


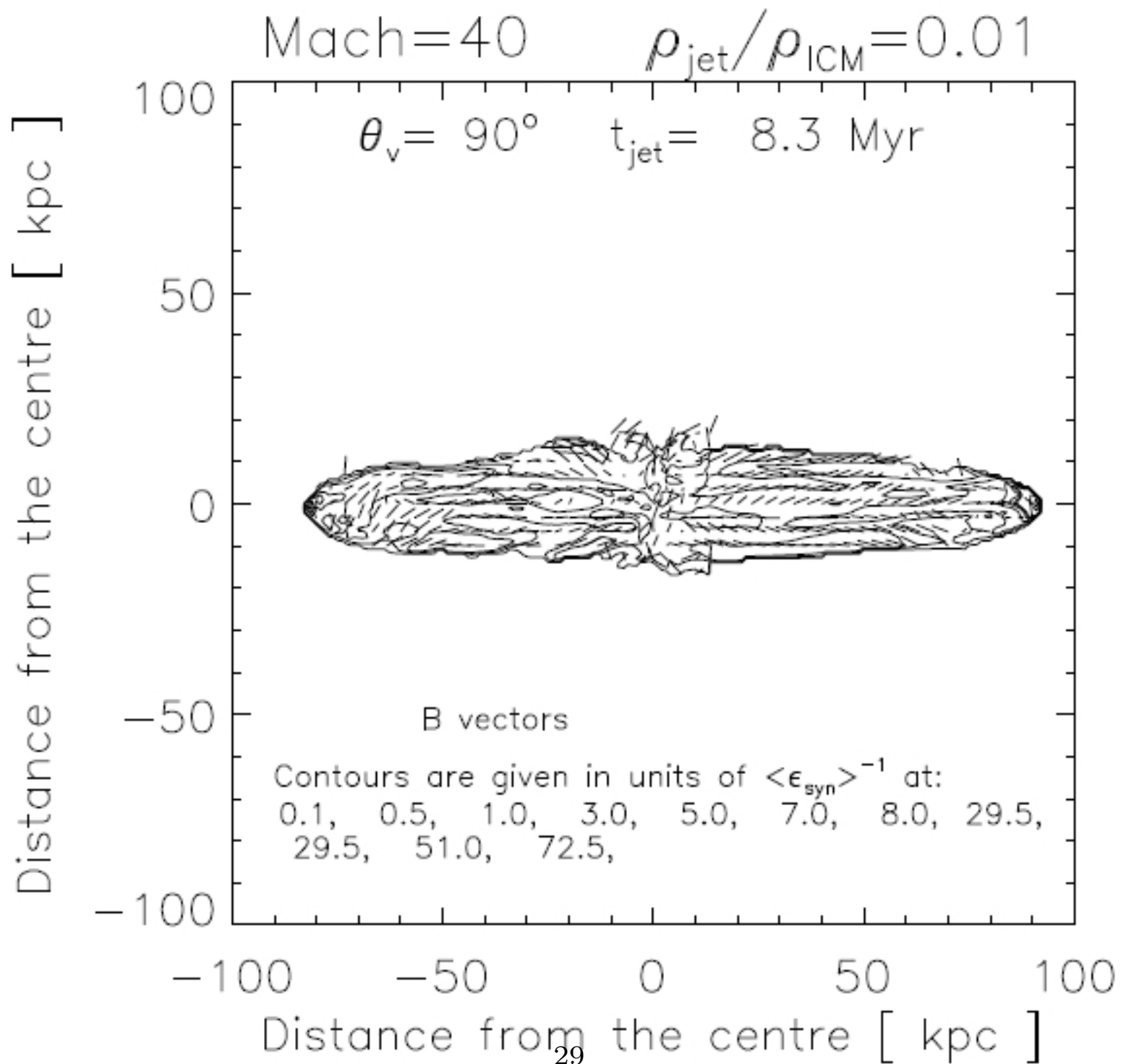


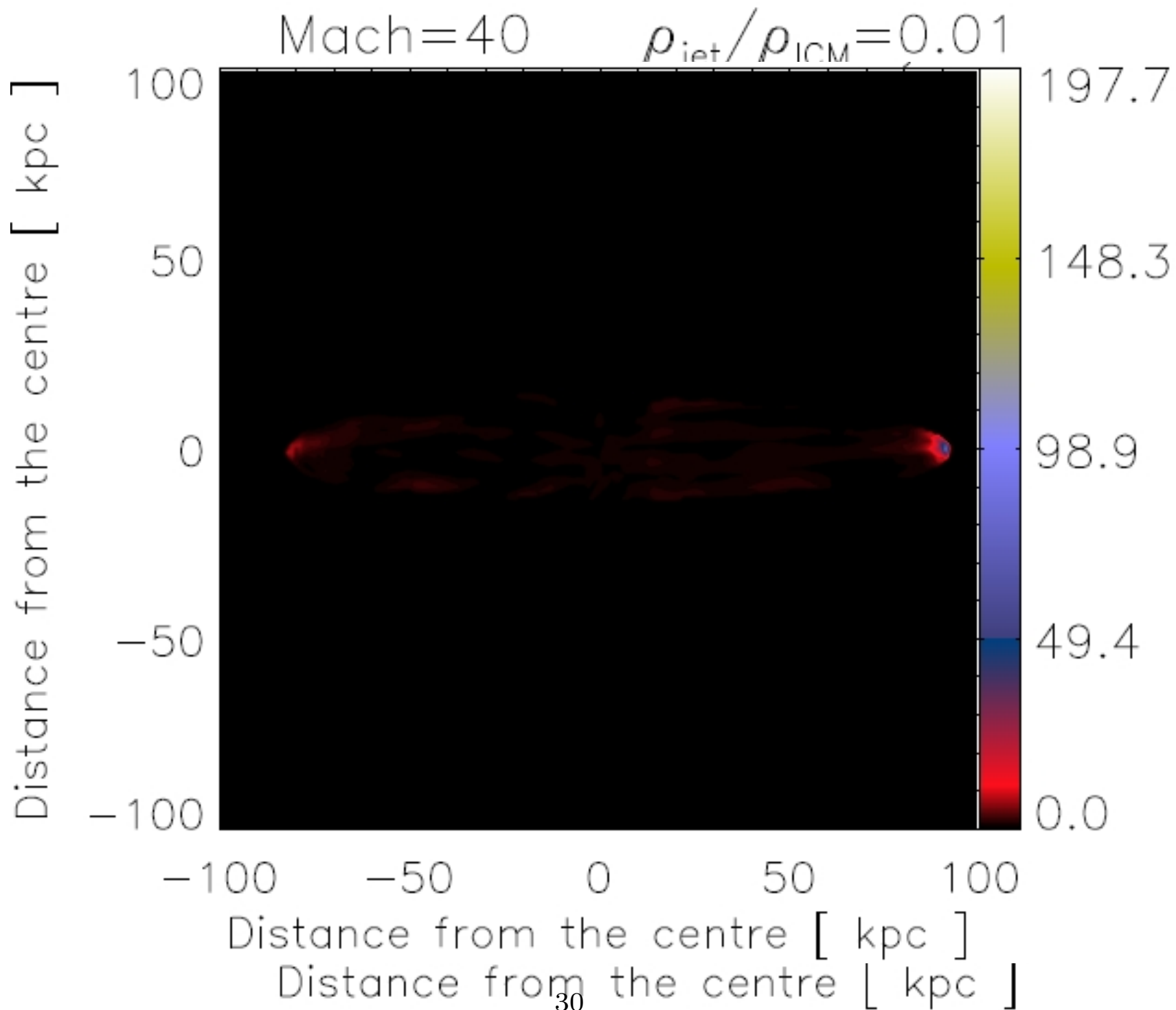


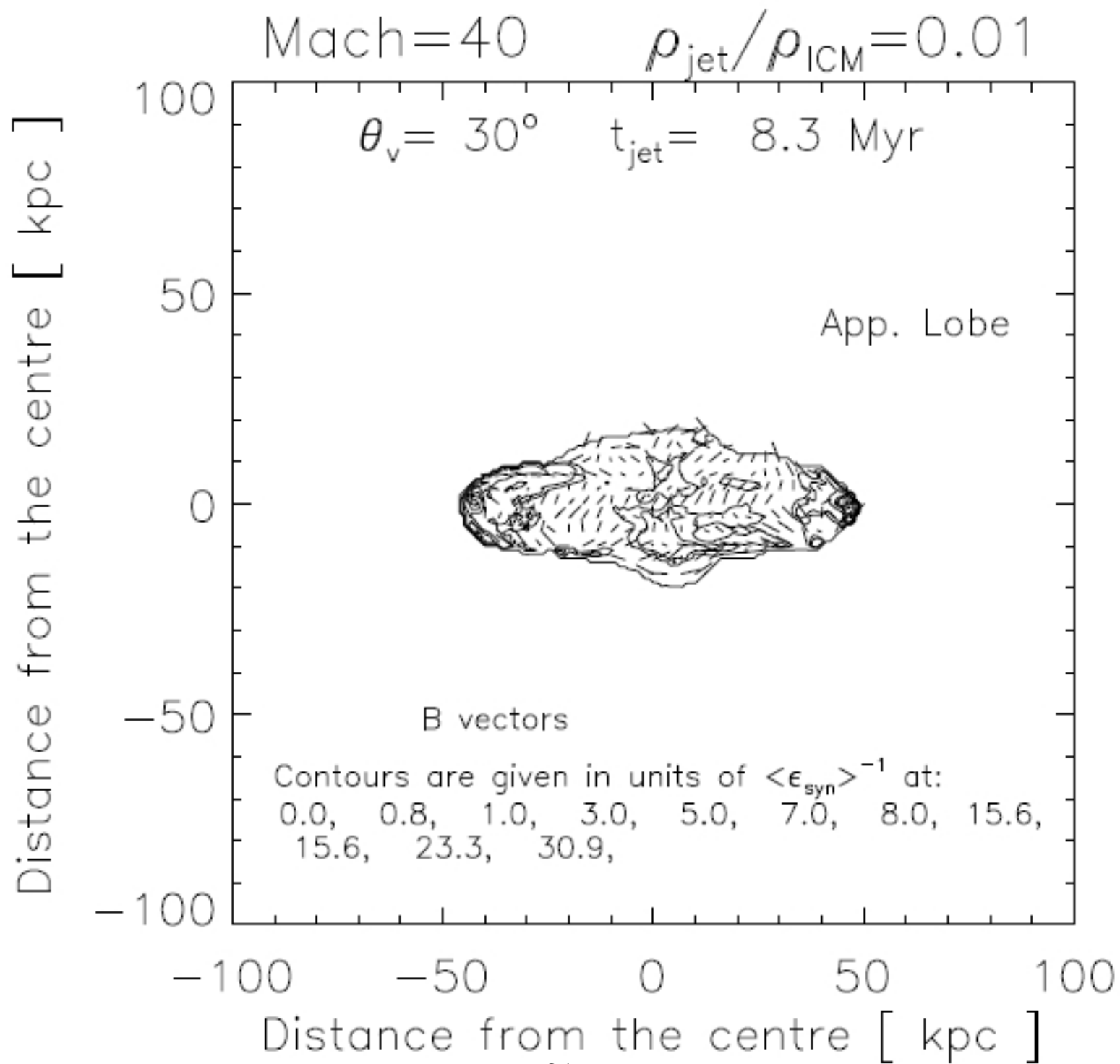
Mach=40

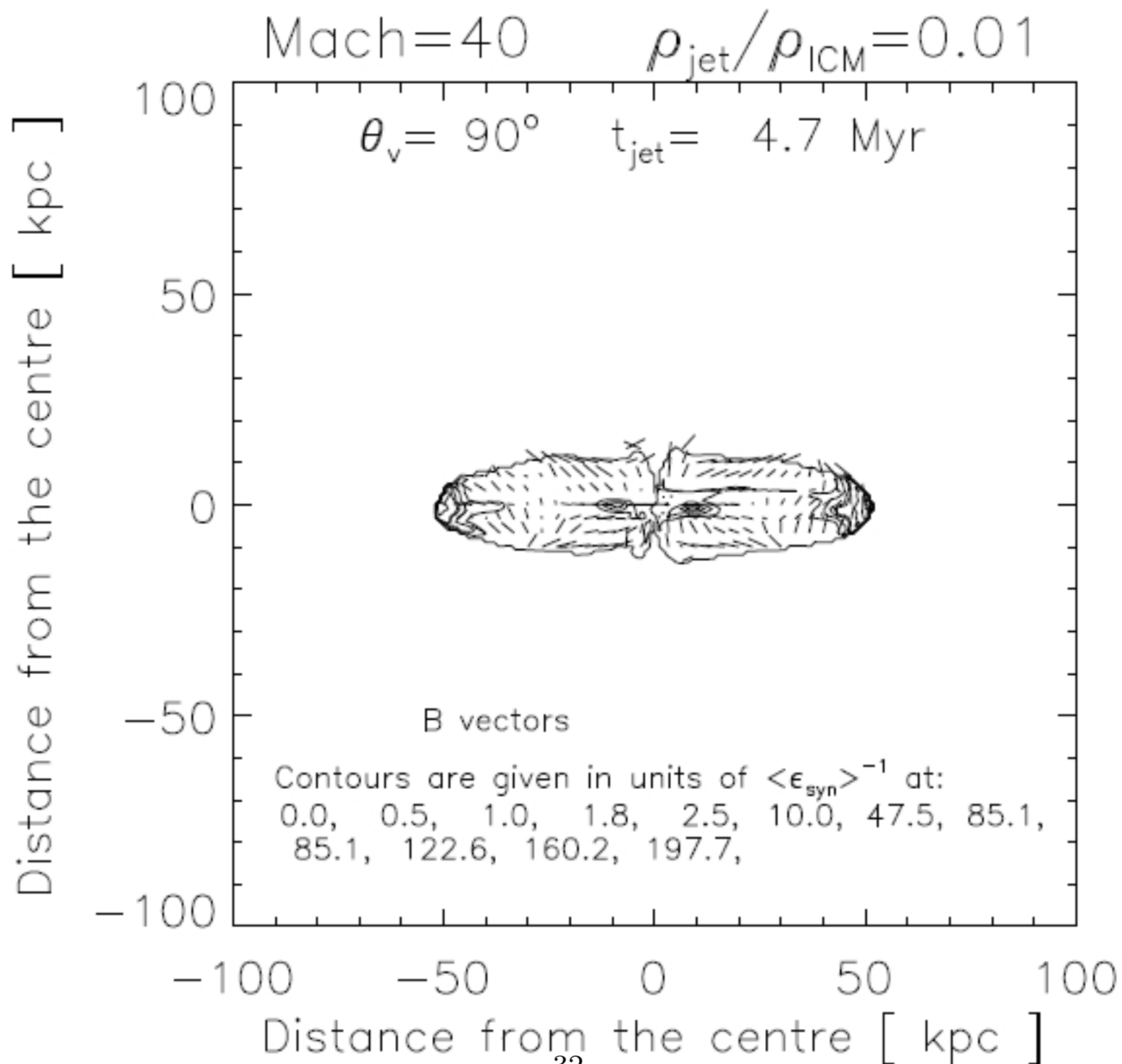
$\rho_{\text{jet}}/\rho_{\text{ICM}}=0.001$

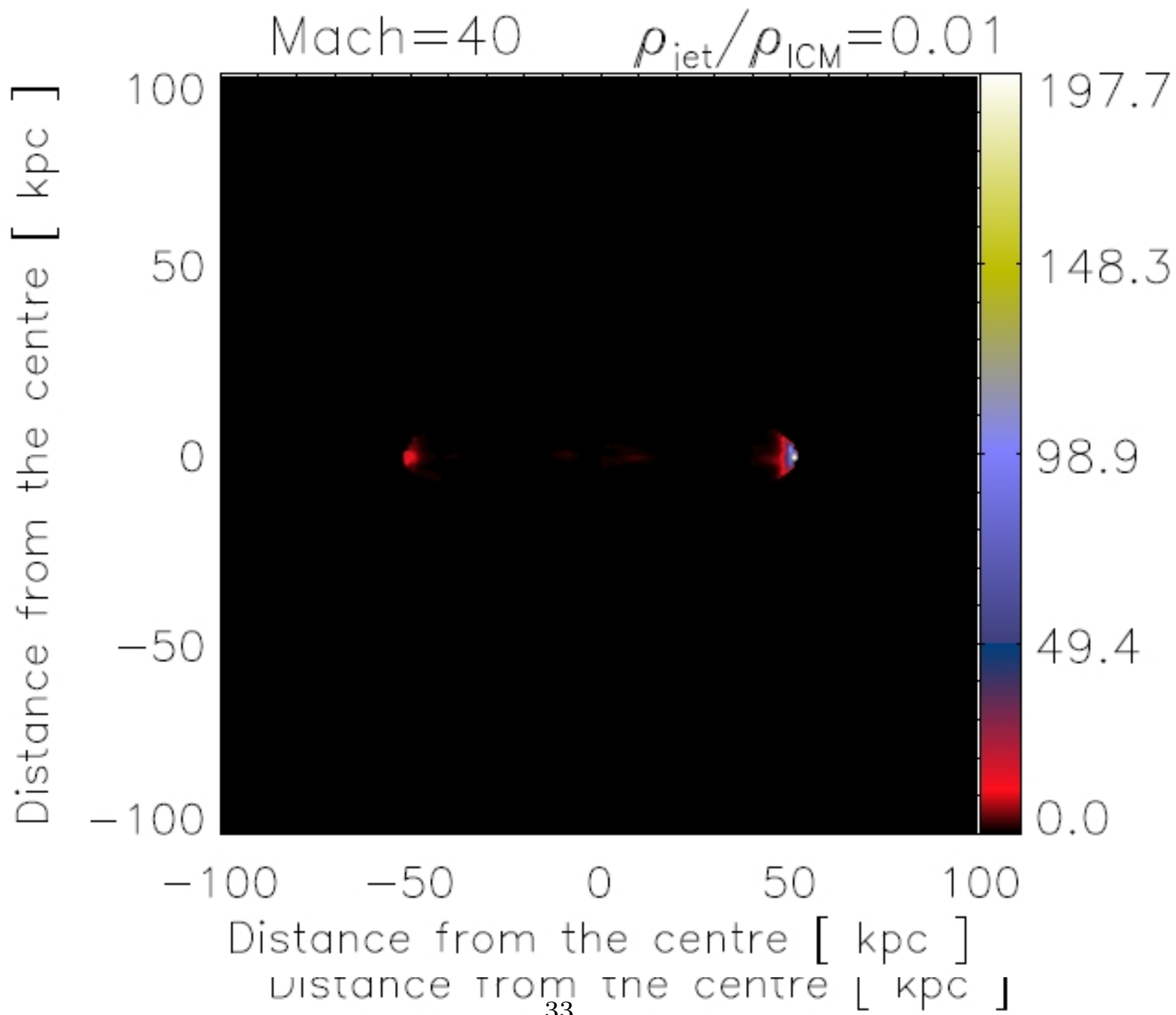












Summary and discussion

- ★ $\chi_B \sim 90^\circ$ along the jets' projection, but it is typically higher elsewhere.
Consistent with topology maps and magnetic flux freezing.
Little dependence with time.
- ★ χ_B seems bipolar when jets have $\eta = 0.01$, but not clear trends for $\eta = 0.001$.
Correlation with sources morphology and energy distribution.
- ★ The outermost B-vectors are generally tangent to the dimmest emissivity contours.
 \sim uniform B-vectors distribution close to shocks but not elsewhere.
- ★ At 90° and 60° the polarisation maps are quite similar but not for 30° .
Though to observe.
- ★ Lots of work to do: couple with RM maps; PD diagrams (luminosity vs size); depolarization; composition of AGN plasma; resolution study; etc.

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Movies @ <http://www.mrao.cam.ac.uk/~mh475/>.

Talk @ <http://www.pas.rochester.edu/martinhe/>.