

# Separation of variables & Sturm-Liouville theory

In physics, we often have problems that depend on several variables. Consider for example the case of a wave in a medium, in 1D, with vanishing boundary conditions:

$$\hat{D} U(x,t) - \frac{1}{c^2} \frac{\partial^2 U(x,t)}{\partial t^2} = 0$$

where  $c$  is the wave's speed, and  $\hat{D}$  is a (e.g. second order) differential operator  $\hat{D} = \frac{d^2}{dx^2} + p(x) \frac{d}{dx} + q(x)$ .

with  $U(x_1, t) = U(x_2, t) = 0$  BC

$$U(x, 0) = U_0(x), \quad \frac{\partial U}{\partial t}(x_0, 0) = U_0'(x). \quad \text{IC.}$$

The standard way of solving this problem is through separation of variables. That is, assume:

$$U(x,t) = y(x)T(t).$$

Substituting this in the D.E. gives

$$T \hat{D} y - \frac{1}{c^2} y T'' = 0$$

or 
$$\frac{\hat{D} y(x)}{y(x)} = \frac{T''(t)}{c^2 T(t)}$$

~~The~~ The LHS depends only on  $x$ , the RHS on  $t$ . Yet, they must be equal  $\forall x, t$ .

Therefore, both sides must be constant

Let this constant be called  $\lambda$ . Then

$$\hat{L}y = \lambda y \quad (1)$$

$$T''(t) = c^2 \lambda T \quad (2)$$

(2) Has a simple solution

$$T(t) = A_1 e^{c\sqrt{\lambda}t} + A_2 e^{-c\sqrt{\lambda}t},$$

while (1) gives the so-called "Sturm-Liouville" problem

$\hat{L}y - \lambda y = 0$	D.E.
$x_1 \leq x \leq x_2$	R.O.I.
$y(x_1) = y(x_2) = 0$	B.C.

As discussed in the notes about selfadjoint operators, we can make the operator selfadjoint by writing  $y(x) = \alpha(x) u(x)$ , where  $\alpha(x)$  is the Wronskian. Then

$$\hat{L}y = \hat{L}u = (\alpha u')' + \beta u, \quad \alpha = w_0 e^{-p(x)}, \quad \beta = (q-p')\alpha$$

The equation for  $u$  is then

$\hat{L}u = \lambda w u$	D. E.
$x_1 \leq x \leq x_2$	R. O. I.
$u(x_1) = u(x_2) = 0$	

where the "weight function"  $w(x)$  is the Wronskian  $\alpha(x)$ .

This problem is analogous to the eigenvalue problems for Matrices:

$$[M] \vec{u} = \lambda \vec{u}$$

which have a discrete set of solutions (eigenvectors)  $\vec{u}_i$  with corresponding eigenvalues  $\lambda$ .  
The same is true for the functional problem:

$$\hat{L} u_i = \lambda_i w u_i$$

i.e. there are a discrete set of solutions  $u_i(x)$  (eigenfunction) with eigenvalues  $\lambda_i$ .

Note that, if we define the inner product:

$$\langle u, v \rangle = \int_{x_1}^{x_2} u^*(x) w(x) v(x) dx$$

↑ weight function, which happens to be equal to the ~~Jacobian~~,  
Wronskian

Then

$$\int_{x_1}^{x_2} u_i^* \hat{L} u_j dx = \lambda_j \int_{x_1}^{x_2} u_i^* w u_j dx = \lambda_j \langle u_i, u_j \rangle$$

but, because  $\hat{L}$  is self adjoint and  $u_i(x_1) = u_i(x_2) = 0$ ,

$$\int_{x_1}^{x_2} u_i^* \hat{L} u_j dx = \left[ \int_{x_1}^{x_2} u_j^* \hat{L} u_i dx \right]^* \quad \text{so}$$

$$\lambda_j \langle u_i, u_j \rangle = \lambda_i^* \underbrace{\langle u_j, u_i \rangle^*}_{\langle u_i, u_j \rangle}$$

or

$$\frac{(\lambda_j - \lambda_i^*) \langle u_i, u_j \rangle = 0}{/}$$

• if  $i=j$ , this means that

$$\lambda_i^* = \lambda_i \Rightarrow \lambda_i = \text{real}.$$

The eigenvalues of self-adjoint operators are always real.

• If  $i \neq j$  and  $\lambda_i \neq \lambda_j$ , then

$$\langle u_i, u_j \rangle = 0$$

The eigenfunctions (with different eigenvalues) of a self-adjoint operator are orthogonal.

• If  $i \neq j$  and  $\lambda_i = \lambda_j$  then  $u_i$  &  $u_j$  are said to be degenerate.  $u_i$  &  $u_j$  are not necessarily orthogonal, but can be recombined into orthogonal eigenfunctions using a Gram-Schmidt procedure.

Notice that the Sturm-Liouville equation can be written as

$$\begin{aligned} \hat{L}u - \lambda w u &= (\alpha u')' + (\beta - \lambda w)u = 0 \\ &= (\alpha u')' + \bar{B}u, \quad \text{where } \bar{B} = \beta - \lambda w \\ &= \beta - \lambda \alpha \end{aligned}$$

Also, recall that  $u(x_1) = u(x_2) = 0$ .

Therefore, we are looking for MODES of the modified self-adjoint operator

$$\hat{L}u = (\alpha u')' + \bar{B}u$$

We know (from the notes on self-adjoint operators) that there are no modes if

$$\alpha \bar{B} \leq 0 \quad \text{for all } x \in [x_1, x_2], \text{ i.e.}$$

$$\alpha \beta - \lambda \alpha^2 \leq 0, \quad \text{that is:}$$

No modes with  $\lambda \geq \beta/\alpha \quad \forall x \in [x_1, x_2]$

That is, there is an upper bound on the possible value of the eigenvalue

The ~~eigenvalue~~ <sup>eigenfunction</sup> with the largest eigenvalue is called the "ground state". (In some applications we let  $\lambda = -\kappa^2$ , where  $\kappa$  then has a lower bound.)