

Mathematical Methods study guide

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1 Einstein Notation

- Dot Product: $\{\vec{A} \cdot \vec{B}\} = A_i B_i$
- Inner Product: $\langle \vec{A}, \vec{B} \rangle = A_i \mu_{ij} B_j$
- Cross Product: $\{\vec{A} \times \vec{B}\}_i = \epsilon_{ijk} A_j B_k$
- Tensorial Product: $\{\vec{A} \otimes \vec{B}\} = A_i B_j$
- Kronecker's delta: $\delta_{ij} = \delta_{ji}$, $\delta_{ij} A_j = A_i$, $\delta_{ii} = 3$ (# dim)
- Levi Civita: $\epsilon_{ijk} = \epsilon_{jki} = \epsilon_{kij}$, $\epsilon_{ijk} = -\epsilon_{jik}$, $\epsilon_{iik} = 0$, $\epsilon_{ijk} \epsilon_{ilm} = \delta_{jl} \delta_{km} - \delta_{jm} \delta_{kl}$
- Derivatives: $\{\vec{r}\}_i = x_i$, $\{\nabla \phi\}_i = \frac{\partial \phi}{\partial x_i} = \partial_i \phi$

2 Derivatives and Integral Theorems

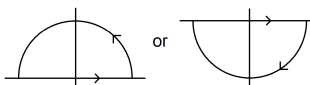
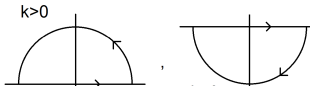
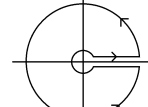
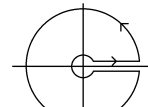
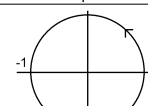
- Gradient: $\{\nabla \phi\}_i = \partial_i \phi$ it gives the rate of change and it points in the direction of the maximum increase.
- Divergence: $\nabla \cdot \vec{A} = \lim_{\Delta \tau \rightarrow 0} \frac{flux\{\vec{A}, \Delta \tau\}}{\Delta \tau}$ where we define $flux\{\vec{A}, \tau\} = \oint_{S \text{ of } \tau} \vec{A} \cdot d\vec{\sigma}$, it can be interpreted as the net flow out of the volume $d\tau$.
- Curl: $\lim_{\Delta \vec{\sigma} \rightarrow \vec{0}} \Delta \vec{\sigma} \cdot (\nabla \times \vec{A}) = \lim_{\Delta \vec{\sigma} \rightarrow \vec{0}} circ\{\vec{A}, \Delta \vec{\sigma}\}$ where we define $circ\{\vec{A}, \vec{\sigma}\} = \oint_{C \text{ of } \vec{\sigma}} \vec{A} \cdot d\vec{l}$, and it can be interpreted as the circulation per unit of area.
- Fundamental theorem of gradients: if $\vec{F}(\vec{r}) = \nabla \phi(\vec{r})$ then $\int_{C \text{ } \vec{r}_a}^{\vec{r}_b} \vec{F}(\vec{r}) \cdot d\vec{l} = \phi(\vec{r}_b) - \phi(\vec{r}_a)$
- Gauss Theorem: if $F(\vec{r}) = \nabla \cdot \vec{A}(\vec{r})$ then $\int_{\tau} \nabla \cdot \vec{A} d\tau = \int_S \vec{A} \cdot d\vec{\sigma}$
- Stokes Theorem: if $\vec{F}(\vec{r}) = \nabla \times \vec{A}(\vec{r})$ then $\int_S \nabla \times \vec{A} \cdot d\vec{\sigma} = \oint_C \vec{A} \cdot d\vec{l}$

3 Complex Variables and Functions

- Cauchy Riemann Conditions: $f(z) = u(x, y) + iv(x, y)$ is analytic if $\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}$ and $\frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x}$, a singularity is a point where they are not satisfied.
- Cauchy's integral theorem: if f is analytic inside and along C then $\oint_C f(z) dz = 0$

- Cauchy's integral formula: $f(z) = \frac{1}{2\pi i} \oint_C \frac{f(z')}{z'-z} dz'$
- Taylor series: $f(z) = \sum_{n=0}^{\infty} f^{(n)}(z_o) \frac{(z-z_o)^n}{n!}$ distance between z and z_o must be smaller than the distance between z_o and the closest singularity.
- Laurent series: $f(z) = \sum_{n=-\infty}^{\infty} a_n(z-z_o)^n$ valid within an annular region centered at z_o , the singularities where we can expand the function as a Laurent series are called poles and if the series starts at the $-n$ power we call this singularity an n^{th} order pole.
- Branch points, branch cuts & Riemann surfaces: because of a branch point the function can take several values and we have # Riemann sheets = # of values of $f(z)$. A branch cut is an artificial line separating different Riemann sheets starting from the branch point. For the special case $f(z) = (z-z_o)^\alpha$ if $\alpha = \frac{M}{N}$ with no common prime factors then there are N Riemann sheets, for the extension $f(z) = (z-z_1)^\alpha (z-z_2)^\beta \dots$ we can choose multiple branch cuts, if the numbers α, β, \dots are rationals the number of Riemann sheets is the lowest common denominator and if the powers of a subset of branch points adds up to an integer then we can join them by a finite branch cut. If instead of a product we have a sum the number of Riemann sheets is the product of the number for each term.
- Residue theorem: $\oint_C f(z) dz = 2\pi i \sum_{k \text{ inside } C} \text{Res}\{f(z); z_k\}$
- Calculating residues: for a n^{th} order pole $\text{Res}\{f(z); z_o\} = \lim_{z \rightarrow z_o} \frac{1}{(n-1)!} \frac{d^{n-1}}{dz^{n-1}} [(z-z_o)^n f(z)]$ and for simple poles $\text{Res}\{\frac{g(z)}{h(z)}; z_o\} = \frac{g(z)}{h'(z)}$

Table 1: Solving definite integrals

Type	Use	Contour	Conditions
$I_1 = \int_{-\infty}^{\infty} f(x) dx$	$\oint_C f(z) dz$		$f(z) \rightarrow 0$ faster than $\frac{1}{z}$ as $ z \rightarrow \infty$
$I_2 = \int_{-\infty}^{\infty} f(x) e^{ikx} dx$	$\oint_C f(z) e^{ikz} dz$		$f(z) \rightarrow 0$ as $ z \rightarrow \infty$
$I_3 = \int_0^{\infty} x^\alpha f(x) dx$	$\oint_C z^\alpha f(z) dz$	 Sheet $[0, 2\pi)$	$z^\alpha f(z) \rightarrow 0$ faster than $\frac{1}{z}$ as $ z \rightarrow \infty$
$I_4 = \int_0^{\infty} f(x) dx$	$\oint_C \ln(z) f(z) dz$	 Sheet $[0, 2\pi)$	$f(z) \rightarrow 0$ faster than $\frac{1}{z^{1+\delta}}$ with $\delta > 0$ as $ z \rightarrow \infty$ and $f(0)$ is finite
$I_5 = \int_0^{2\pi} F(\cos \theta, \sin \theta) d\theta$	$\oint_{u.c} f\left(\frac{z+1/z}{2}, \frac{z-1/z}{2i}\right) \frac{dz}{iz}$	 unit circle	none

- Infinite sums: $\sum_{n=-\infty}^{\infty} f(n) = - \sum_{k, \text{poles of } f} \text{Res}\{\gamma(z) f(z); z_k\}$ where $\gamma(z) = \pi \cot(\pi z)$.

- Principal values: when we have simple poles along the contour,

$$\mathcal{P} \oint_C f(z) dz = 2\pi i \sum_{k \text{ inside } C} \text{Res}\{f(z); z_k\} + \pi i \sum_{k \text{ along } C} \text{Res}\{f(z); z_k\}$$

- Gaussian integrals: $\int_{-\infty}^{\infty} x^n e^{-ax^2/2} dx = \begin{cases} \frac{(n-1)!!}{a^{n/2}} \sqrt{\frac{2\pi}{a}} & , n=\text{even} \\ 0 & , n=\text{odd} \end{cases}$
and $\int_{-\infty}^{\infty} x^n e^{ibx^2/2} dx = \begin{cases} \frac{(n-1)!!}{(-ib)^{n/2}} \sqrt{\frac{2\pi}{|b|}} e^{i\frac{\pi}{4} \text{sgn}(b)} & , n=\text{even} \\ 0 & , n=\text{odd} \end{cases}$

- Stationary phase: x_o is a stationary point if $\phi'(x_o) = 0$

$$\int_a^b A(x) e^{ik\phi(x)} dx = \sum_{j, \text{stat pts a,b}} A(x_j) e^{ik\phi(x_j)} \sqrt{\frac{2\pi}{k|\phi''(x_j)|}} e^{i\frac{\pi}{4} \text{sgn}(\phi''(x_j))} + \frac{A(x)}{ik\phi'(x)} e^{ik\phi(x)} \Big|_a^b$$

4 Integral transforms

- Functional analysis:

- Cauchy-Schwartz inequality: $\|f\|^2 \|g\|^2 \geq |\langle f, g \rangle|^2$ if $w(t) > 0$ and $\langle f, g \rangle = \int f^*(t) w(t) g(t) dt$. The equality is satisfied when f is proportional to g .
- Kernel (\sim matrix): $\tilde{f}(\tau) = \int \mathcal{M}(\tau, t) f(t) dt$ equivalent to $A' = MA$ for matrices and vectors.
- Unitary transformations: $\langle \tilde{f}, \tilde{g} \rangle = \langle f, g \rangle \Leftrightarrow \int \mathcal{U}^*(\tau, t_1) \mathcal{U}(\tau, t_2) d\tau = \delta(t_2 - t_1)$

- Fourier transform: $\mathcal{U}_F(p, x) = \sqrt{\frac{|\eta|}{2\pi}} e^{-i\eta xp}$

$$\tilde{f}(p) = \hat{\mathcal{F}}_{x \rightarrow p} f(x) = \sqrt{\frac{|\eta|}{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\eta xp} dx$$

$$f(x) = \hat{\mathcal{F}}_{x \rightarrow p}^{-1} \tilde{f}(p) = \sqrt{\frac{|\eta|}{2\pi}} \int_{-\infty}^{\infty} \tilde{f}(p) e^{i\eta xp} dx$$

- Properties:

- Shift and phase: $\hat{\mathcal{F}}_{x \rightarrow p}[f(x - a)] = \tilde{f}(p) e^{-i\eta ap}$ and $\hat{\mathcal{F}}_{x \rightarrow p}[f(x) e^{i\eta ap}] = \tilde{f}(p - a)$.
- Scale: $\hat{\mathcal{F}}_{x \rightarrow p}[f(ax)] = \frac{1}{|a|} \tilde{f}\left(\frac{p}{a}\right)$
- Convolution and product: we define $f(x) * g(x) = \int f(x') g(x - x') dx'$ so we have that $\hat{\mathcal{F}}_{x \rightarrow p}[f(x) g(x)] = \sqrt{\frac{|\eta|}{2\pi}} (\tilde{f} * \tilde{g})(p)$ and $\hat{\mathcal{F}}_{x \rightarrow p}[(f * g)(x)] = \sqrt{\frac{2\pi}{|\eta|}} \tilde{f}(p) \tilde{g}(p)$
- Derivative: $\hat{\mathcal{F}}_{x \rightarrow p}[f^{(n)}(x)] = (i\eta p)^n \tilde{f}(p)$
- Parseval's theorem: $\int_{-\infty}^{\infty} f^*(x) g(x) dx = \int_{-\infty}^{\infty} \tilde{f}^*(p) \tilde{g}(p) dp$.
- Uncertainty relation: we define the mean $\bar{x} = \frac{\int x |f(x)|^2 dx}{\int |f(x)|^2 dx}$ and the variance $\Delta x^2 = \frac{\int (x - \bar{x})^2 |f(x)|^2 dx}{\int |f(x)|^2 dx}$, we have $\Delta x \Delta p \geq \frac{1}{2|\eta|}$.

- Laplace transform:

$$\bar{f}(s) = \hat{\mathcal{L}}_{t \rightarrow s} f(t) = \int_0^{\infty} f(t) e^{-st} dt \quad t > 0, f \text{ can diverge as } t \rightarrow \infty$$

$$f(t) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \bar{f}(s) e^{st} ds \quad c \text{ large enough so that poles lie to the left of } c$$

- Properties with derivatives: $\hat{\mathcal{L}}_{t \rightarrow s} f^{(n)}(t) = s^n \bar{f}(s) - s^{n-1} f(0) - s^{n-2} f'(0) - \dots - f^{(n-1)}(0)$ used to solve ODE with constant coefficients.
- Hilbert transform: derived from $f(x) = -\frac{i}{\pi} \mathcal{P} \int_{-\infty}^{\infty} \frac{f(x')}{x' - x} dx'$ taking $f(x) = u(x) + iv(x)$

$$v(x) = -\frac{1}{\pi} \mathcal{P} \int_{-\infty}^{\infty} \frac{u(x')}{x' - x} dx'$$

$$u(x) = \frac{1}{\pi} \mathcal{P} \int_{-\infty}^{\infty} \frac{v(x')}{x' - x} dx'$$

- Titchmarsh theorem: the following statements are equivalent,

1. $\tilde{f}(p) = 0 \forall p < 0$ (assuming $\eta > 0$)
2. $f(z)$ is analytic along and above the real axis and goes to zero as $|z| \rightarrow \infty$ within the upper half plane
3. the real and imaginary parts of f are related via the Hilbert transform

5 Ordinary differential equations

- Second order: $y'' + py' + qy = -\rho$, we divide in the homogeneous and particular solution $y = \phi_h + \phi_p$, the homogeneous has two independent solutions, $\phi_h = A_1 \phi_1 + A_2 \phi_2$, where the constants are chosen so that it satisfies the IC/BC and the particular is zero at the BC/IC.
- Wronskian: $W(x) = \phi_1 \phi_2' - \phi_2 \phi_1' = e^{-\int^x p(x') dx'}$
- Variational method: if ϕ_1 is known we suppose $\phi_2 = u \phi_1$ and so we get that $\phi_2 = \phi_1 \int^x \frac{W(x')}{\phi_1^2(x')} dx'$
- Frobenius method: works if p has up to a simple pole and q has up to a second order pole at x_o

1. set $\phi(x) = \sum_{n=0}^{\infty} a_n (x - x_o)^{n+\gamma}$ where $a_0 = 1$

2. expand p and q around x_o

3. find equations in terms of a_n and γ

4. determine $\gamma_{1,2}$ from the indicial equation (the first one, usually $n = 0$) and solve for a_n

Note: second solution won't work if $\gamma_1 - \gamma_2 = 0$ or $N > 0$ integer, near x_o $\phi_1 \propto (x - x_o)^{\gamma_1}$ (same for ϕ_2 if the difference of the γ 's is not an integer) and $\phi_2(x) \propto \begin{cases} (x - x_o)^{\gamma_2} & \text{if } \gamma_1 - \gamma_2 = N > 0 \\ (x - x_o)^{\gamma_1} \ln(x - x_o) & \text{if } \gamma_1 - \gamma_2 = 0 \end{cases}$

- Green function: $G(x, x') = \begin{cases} a\phi_1(x) + b\phi_2(x) & \text{if } x < x' \\ c\phi_1(x) + d\phi_2(x) & \text{if } x > x' \end{cases}$ G vanishes at BC/IC and satisfies the matching conditions $\begin{matrix} a\phi_1(x') + b\phi_2(x') & = & c\phi_1(x') + d\phi_2(x') \\ a\phi_1'(x') + b\phi_2'(x') & = & c\phi_1'(x') + d\phi_2'(x') + 1 \end{matrix}$ for IC $a = b = 0$.
- Particular solution: $\phi_p(x) = \int_{\text{R.o.I}} \rho(x')G(x, x')dx'$
- Modes: solution to homogeneous equation that vanishes at the endpoints x_0 and x_1 .
- Self-adjoint operators: $\langle v, \hat{\mathcal{L}}u \rangle = \langle u, \hat{\mathcal{L}}v \rangle^* + [\]_{x_0}^{x_1}$ if $\hat{\mathcal{L}} = p_0(x)\frac{d^2}{dx^2} + p_1(x)\frac{d}{dx} + p_2(x)$ then we must have $p_0' = p_1$ and we can write $\hat{\mathcal{L}}u = (\alpha u')' + \beta u$, if instead we have $\hat{\mathcal{D}}y = y'' + py' + qy$ which is not in self-adjoint form we let $y = \alpha u$ and so $\alpha = W$ and $\beta = (q - p')W$ ($\hat{\mathcal{D}}y = \hat{\mathcal{L}}u = 0$) and there can be no modes if $\alpha\beta \leq 0 \forall x \in [x_0, x_1]$.

6 Other useful stuff

- Geometric series: $\sum_{n=0}^{\infty} w^n = \frac{1}{1-w}$ for $|w| < 1$
- Binomial expansion: $(1 + w)^\alpha = \sum_{n=0}^{\infty} \frac{\prod_{p=0}^{n-1} (\alpha - p)}{n!} w^n$
- Partial fractions: when quadratic there are two terms.