Solving ablative RT Problem in AstroBEAR

1.Quasi-equilibrium

The quasi-equilibrium in RT problem is defined by the balance of the thermal pressure, Ram pressure and the gravity force. In equilibrium, the Euler equations are:

$$\frac{\partial \rho v}{\partial y} = 0$$

$$\rho v \frac{\partial v}{\partial y} = -\frac{\partial p}{\partial y} - \rho g$$

$$\rho v C_v \frac{\partial T}{\partial y} = -p \frac{\partial v}{\partial y} + \frac{\partial}{\partial y} (\kappa T^n \frac{\partial T}{\partial y})$$

from the first equation we know that $\rho v = \rho_0 v_0$, which is just the outflow at the bottom. If we define the heat flux as:

$$q = \kappa T^n \frac{\partial T}{\partial y}$$

we can then write the equations as:

$$\frac{\frac{\partial \rho}{\partial y}}{\frac{\partial r}{\partial y}} = \frac{\frac{\rho R q}{\kappa T^{n}} + \rho g}{\frac{\rho_{0}^{2} v_{0}^{2}}{\rho^{2}} - RT}$$

$$\frac{\frac{\partial T}{\partial y}}{\frac{\partial q}{\partial y}} = \frac{q}{\kappa T^{n}}$$

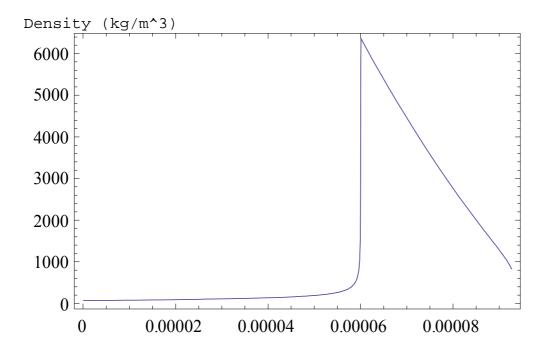
$$\frac{\frac{\partial q}{\partial y}}{\frac{\rho_{0}^{2} v_{0}^{2} - RT}{\kappa T^{n}} - \rho_{0} v_{0} RT \frac{\frac{R q}{\kappa T^{n}} + g}{\frac{\rho_{0}^{2} v_{0}^{2}}{\rho^{2}} - RT}$$

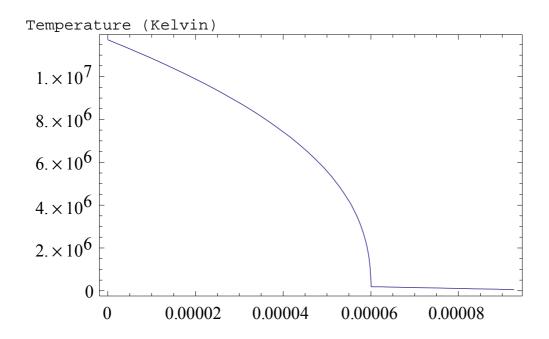
with ρ , v, q known on the two boundaries.

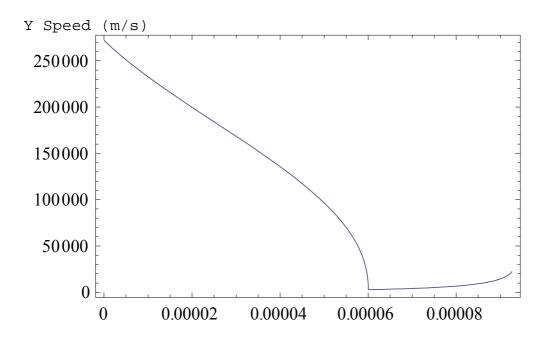
 $v_0 = -272144.604867564$

Using parameters (in SI units, Temperature is measured in Joule): R = 4.7904e+26 $C_v = 7.186e+26$ $\kappa = 3.734e+69$ g = 1.0e+14 q = -5.876e+18 $\rho_0 = 68.1622919147237$

The solutions to this equation set are posted below:







2. Euler Fluid Equations with diffusion in c.g.s.

The Euler equations can be written as:

$$\rho_t + (\rho u)_x + (\rho v)_y = 0$$

$$(\rho u)_t + (\rho u^2 + p)_x + (\rho u v)_y = 0$$

$$(\rho v)_t + (\rho u v)_x + (\rho v^2 + p)_y = -\rho g$$

$$(\frac{\rho v_t^2}{2} + \frac{p}{\gamma - 1})_t + \nabla \bullet [v_t (\frac{\rho v_t^2}{2} + \frac{\gamma p}{\gamma - 1})] = \nabla \bullet (\kappa T^n \nabla T)$$

From the last equation, we have:

$$\left(\frac{\rho v_t^2}{2}\right)_t + \left(\frac{p}{\gamma - 1}\right)_t + \left(\frac{\rho v_t^2}{2} + \frac{\gamma p}{\gamma - 1}\right) \bigtriangledown \bullet v_t + v_t \bullet \bigtriangledown \left(\frac{\rho v_t^2}{2} + \frac{\gamma p}{\gamma - 1}\right) = \bigtriangledown \bullet \left(\kappa T^n \bigtriangledown T\right)$$

The second term can be written as:

$$\frac{p}{\gamma-1} = \frac{(1+Z)\rho k_B T}{(\gamma-1)m_i}$$

If we write $c_v = \frac{(1+2)}{(\gamma-1)m_i}$, and measure T in terms of energy, then we have the EOS:

$$\frac{p}{\gamma-1} = c_{\nu} \rho T$$

then the energy equation can be written as:

$$\left(\frac{\rho v_t^2}{2}\right)_t + \left(c_v \rho T\right)_t + \left(\frac{\rho v_t^2}{2} + \gamma c_v \rho T\right) \nabla \bullet v_t + v_t \bullet \nabla \left(\frac{\rho v_t^2}{2} + \gamma c_v \rho T\right) = \nabla \bullet \left(\kappa T^n \nabla T\right)$$

Simplify this equation by using the momentum and mass conservations, we get the energy equation in the form of:

 $\rho c_{v}(\partial_{t}T + v \bullet \nabla T) = -p \nabla \bullet v + \nabla \bullet (\kappa T^{n} \nabla T)$

Notice that in this form, all quantities are in SI with T measured in Joule. So it's actually: $\rho^{s} c_{v}^{s} (\partial_{t} k^{s} T^{s} + v^{s} \bullet \nabla k^{s} T^{s}) = -p^{s} \nabla^{s} \bullet v^{s} + \nabla^{s} \bullet (\kappa^{s} (k^{s})^{n} (T^{s})^{n} \nabla^{s} k^{s} T^{s})$ To convert to the c.g.s units, we have the following conversion factors:

$\rho^{S} = 1000 \rho^{G}$	$p^{S} = 1/10 p^{G}$
$v^{S} = 1/100 v^{G}$	$\nabla^{s} = 100 \nabla^{G}$
$c_{v}^{S} = 1000 c_{v}^{G}$	$k^{S} = 10^{-7} k^{G}$

Substitute, we find that:

$$\rho^{G} c_{v}^{G} (\partial_{t} k^{G} T^{G} + v^{G} \bullet \nabla k^{G} T^{G}) = -p^{G} \nabla^{G} \bullet v^{G} + \nabla^{G} \bullet (10^{(-7n-2)} \kappa^{S} (k^{G})^{n} (T^{G})^{n} \nabla^{G} k^{G} T^{G})$$

If we define $\kappa^{G} = 10^{(-7n-2)} \kappa^{S}$, then:

$$\rho^{G} c_{v}^{G} (\partial_{t} k^{G} T^{G} + v^{G} \bullet \nabla k^{G} T^{G}) = -p^{G} \nabla^{G} \bullet v^{G} + \nabla^{G} \bullet (\kappa^{G} (k^{G})^{n} (T^{G})^{n} \nabla^{G} k^{G} T^{G})$$

Using operator splitting method, we need to solve:

$$\rho c_{v}^{G}(\partial_{t} k^{G} T) = \nabla^{G} \bullet (\kappa^{G} (k^{G})^{n} T^{n} \nabla k^{G} T)$$

Divide by k and cv in c.g.s., we get:

$$\rho(\partial_t T) = \nabla \bullet \left(\frac{\kappa^G(k^G)^n}{c_v^G} T^n \nabla T \right)$$

define:

$$\kappa_{0} = \frac{\kappa^{G}(k^{G})^{n}}{c_{v}^{G}} = \frac{10^{(-7n+1)}(k^{G})^{n}\kappa^{S}}{c_{v}^{S}}$$

then we have the equation ready to be scaled:

$$\rho(\partial_t T) = \nabla \bullet (\kappa_0 T^n \nabla T)$$

For the flux equation, we have:

$$q^{S} = \kappa^{S} (k^{S})^{n} T^{n} \nabla^{S} T$$

we need a form like:

$$q_0 = \kappa_0 T^n \nabla^G T$$

we can divide these two equations and use the relations between SI units and Gauss units. The result is:

$$q_0 = \frac{q^s}{10 c_v^s k^s}$$

The scales can be obtained by writing equation $\rho(\partial_t T) = \nabla \bullet (\kappa_0 T^n \nabla T)$ in the forms of scalings:

$$\frac{rscale\,Tscale}{tscale} = \frac{\kappa\,scale\,Tscale^{(n+1)}}{lscale^2}$$

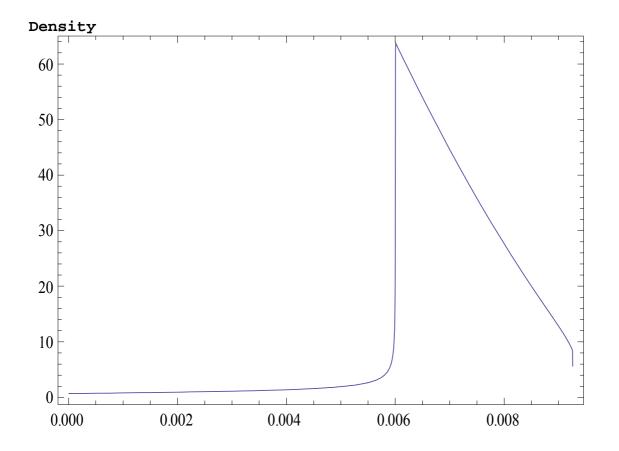
simplify this form, we get:

$$\kappa scale = \frac{lscale \, rscale^{1/2} \, pscale^{1/2}}{Tscale^n}$$

also, we have:

$$qscale = \frac{\kappa \, scale \, Tscale^{n+1}}{lscale} = rscale^{1/2} \, pscale^{1/2} \, Tscale$$

Using the rescaled parameters, the Euler equation is solved and the equilibrium is plotted as follows:



Temperature

