

# Conductors

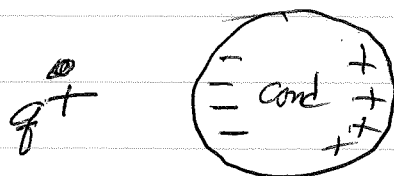
## Properties in electrostatics

- 1)  $\vec{E} = 0$  inside a conductor - if  $\vec{E} \neq 0$  then a current would flow  $\Rightarrow$  not an electrostatic situation.  
When first apply  $\vec{E}$  to a conductor, charges flow until they result in a configuration such that the charges in the conductor produce an electric field that exactly cancels the applied field so that the total field inside the conductor is zero
- 2)  $\rho = 0$  inside conductor - follows from  $\vec{\nabla} \cdot \vec{E} = \rho/\epsilon_0$  and  $\vec{E} = 0$  inside
- 3) (2)  $\Rightarrow$  any net charge on a conductor must lie on the surface of the conductor
- 4)  $V = \text{constant}$  throughout the conductor - since  $-\vec{\nabla}V = \vec{E}$  and  $\vec{E} = 0$  inside
- 5) Just outside the conductor  $\vec{E} \perp$  to surface  
- since if  $\vec{E}$  did have a component  $\parallel$  to surface, it would exert forces on charges on surface causing them to flow - ~~same for~~ so would not be electrostatics

## Consequences of above properties

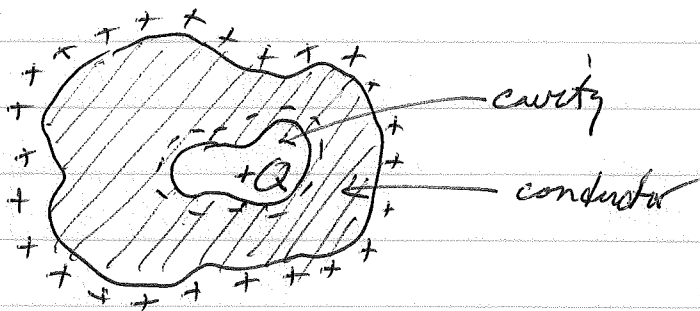
### Induced charge

point charge  
in front of  
conducting  
sphere



conductor has no net charge  
 $\Rightarrow$  induced charge on surface  
to cancel out  $\vec{E}$  from  
 $q$  outside so  $\vec{E} = 0$  inside  
sphere. But now there  
will be a net attractive  
force between  $q$  and the conductor  
even though conductor has no net charge

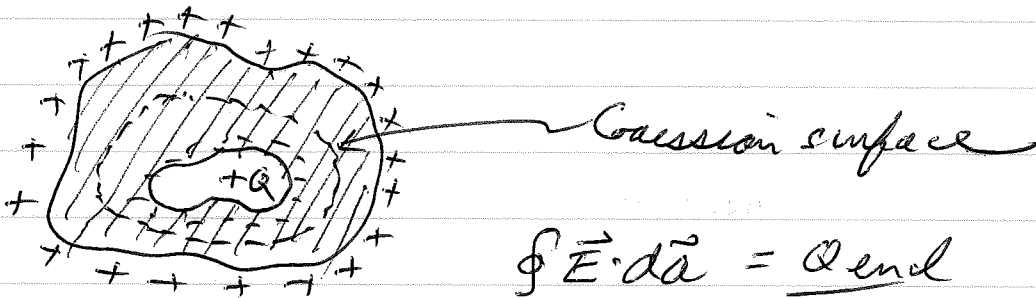
## Cavities in conductors



+Q inside cavity attracts oppositely charged (-) electrons in conductor to ~~inner~~ inner surface, so that net field inside conductor is zero. Equal amount of (+) charge builds up on outside surface of conductor to give  $\vec{E}$  field outside

Amount of induced charge on inner surface totals to -Q  
 " " " " " outer " " " +Q

to see, consider a Gaussian surface inside conductor that encloses the cavity



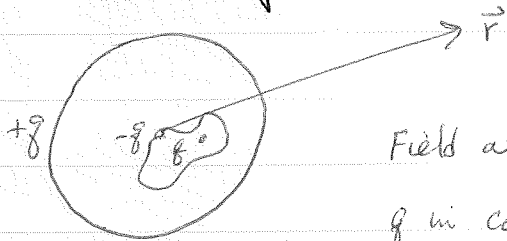
$$\oint \vec{E} \cdot d\vec{a} = \frac{Q_{\text{enc}}}{\epsilon_0}$$

but  $\vec{E} = 0$  inside conductor so

$$\oint \vec{E} \cdot d\vec{a} = 0 \Rightarrow Q_{\text{enc}} = 0$$

$\Rightarrow$  induced charge on surface cavity must total to -Q

conducting ~~shell~~ sphere with a cavity inside



Field at  $\vec{r}$ ?

$q$  in cavity induces  $-q$  on inside surface that cancels out net electric field inside conductor  $\frac{E}{2}$   
(field inside cavity ~~need~~ is not in general zero)

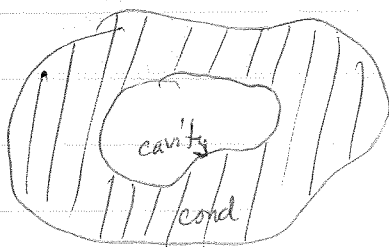
$+q$  is induced on outer surface, spreads

uniform since  $E \neq 0$  inside  $\rightarrow$  no interaction with charges inside

uniformly  $\Rightarrow \vec{E}(\vec{r}) = \frac{q}{4\pi\epsilon_0 r^2} \hat{r}$   $\leftarrow$  gives  $\vec{E} = 0$  inside sphere

true for any shape cavity at all!

Field inside charge free cavity = 0



$$-\nabla^2 V = \rho/\epsilon_0 = 0 \text{ inside cavity}$$

$$V(\vec{r}) = V_0 \text{ constant on surface of cavity}$$

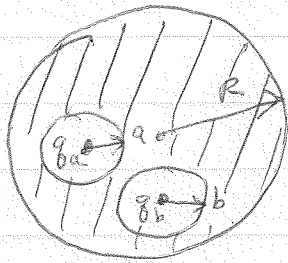
$$\Rightarrow V(\vec{r}) = V_0 \text{ inside cavity is solution}$$

$$(-\nabla^2 V_0 = 0)$$

$$\Rightarrow \vec{E} = -\vec{\nabla} V = 0 \text{ inside}$$

Faraday cage: putting apparatus inside a grounded conducting cage screens out all external electric fields

Prob 2.39



$q_a, q_b$  at centers of spherical cavities

a) surface charge  $\sigma_a, \sigma_b, \sigma_R$

uniform  $\left\{ \begin{array}{l} \sigma_a = -q_a/4\pi a^2 - \text{take gaussian surface about cavity a. } \oint \vec{E} \cdot d\vec{a} = 0 \\ \Rightarrow Q_{\text{enc}} = 0 \\ \sigma_b = -q_b/4\pi b^2 \\ \sigma_R = (q_a + q_b)/4\pi R^2 \end{array} \right.$

b)  $\vec{E}$  outside  $\hookrightarrow \vec{E} = \frac{(q_a + q_b)}{4\pi\epsilon_0 r^2} \hat{r}$   $\leftarrow \text{ie } \sigma_R \text{ is uniform}$

c)  $\vec{E}$  inside cavities

$$\vec{E}_a = \frac{q_a}{4\pi\epsilon_0 r^2} \hat{r}$$

$$\vec{r} = \vec{r} - \vec{r}_a$$

- field has spherical symmetry about center of cavity

$$\rightarrow \oint \vec{E} \cdot d\vec{a} = \frac{q_{\text{enc}}}{\epsilon_0}$$

easy to get  $\vec{E}$

$$\vec{E}_b = \frac{q_b}{4\pi\epsilon_0 r^2} \hat{r}$$

$$\vec{r} = \vec{r} - \vec{r}_b$$

d) force on  $q_a$  and  $q_b = 0$

Considers force on  $q_a$ :  $\vec{E}$  field at  $\vec{r}_a$  from  $\sigma_R = 0$  as we are inside uniformly charged shell

$\vec{E}$  field from  $q_b$  exactly cancelled by  $\vec{E}$  field from  $\sigma_b$

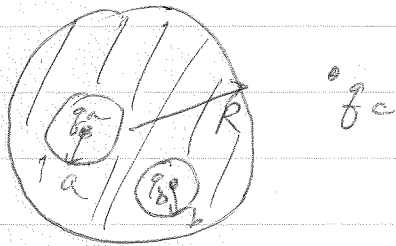
$\vec{E}$  field at  $\vec{r}_a$  from  $\sigma_a = 0$  as we are inside uniformly charged shell

$$\Rightarrow \vec{E}(\vec{r}_a) = 0. \text{ Similarly } \vec{E}(\vec{r}_b) = 0 \Rightarrow \text{force on a, b} = 0$$

$\uparrow$  field from all charges excluding  $q_a$

$\uparrow$  field from all charges excluding  $q_b$

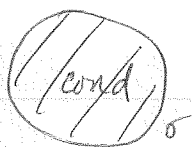
e) What changes if  $q_c$  ~~is~~ brought near outside of conductor?



- $\sigma_a, \sigma_b$  unchanged.  $\sigma_R$  still has total charge  $+q_a + q_b$  but now charge is more concentrated on side of sphere closest to  $q_c$
- $\vec{E}$  outside changes, as result of redistribution of charge  $\sigma_R$
- $\vec{E}_a \neq \vec{E}_b$  do not change
- Force on  $q_a + q_b$  do not change.  $\sigma_R$  must redistribute itself so that field from  $\sigma_R$  and  $q_c$  cancel inside the conductor + so give no force on  $q_a$  or  $q_b$ . Force on  $q_a$  from  $\sigma_a, \sigma_b$  and  $q_b$  is as before, so total force on  $q_a = 0$ . Same for  $q_b$

Conductor acts to insulate the cavities from each other and from the outside.

12.5.3) Surface charge on conductor



$\vec{E}_{above} - \vec{E}_{below} = \frac{\sigma}{\epsilon_0} \hat{n}$  as cross surface charge density  $\sigma$ .

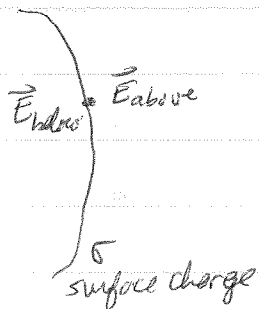
for conductor,  $\vec{E}_{below} = \vec{E}_{inside} = 0 \Rightarrow \vec{E}_{above} = \boxed{\vec{E} = \frac{\sigma}{\epsilon_0} \hat{n}}$

field at surface is in normal direction

or  $\vec{E} \cdot \hat{n} = -\vec{\nabla}V \cdot \hat{n} = \boxed{-\frac{\partial V}{\partial n} = \frac{\sigma}{\epsilon_0}}$

$\Rightarrow$  can determine  $\sigma$  on surface of conductor if know  $V$  (or  $\vec{E}$ ) outside

Electrostatic pressure at surface charge layer



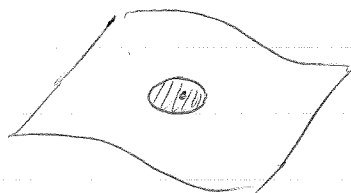
electric field  $\vec{E}$  exerts force  $\vec{F}$  on charge density  $\sigma$ .

$\vec{F} = \sigma \vec{E}$  ← but do we use  $\vec{E}_{above}$  or  $\vec{E}_{below}$ ?

↑ force/area

ans:  $\vec{E} = \frac{1}{2} (\vec{E}_{above} + \vec{E}_{below})$   
average value

proof:



write  $\vec{E} = \vec{E}_{patch} + \vec{E}_{other}$

↑ electric field from  $\sigma$  on small patch at  $\vec{r}$  on surface

$\vec{E}_{patch}$  does not exert force on the charge on the patch itself  
ie we do not want any self forces. patch feels force only from the other charges off the patch, ie  $\sigma$  outside patch + any external charges off the surface - these are the charges that determine  $\vec{E}_{other}$

So  $\vec{f}_{\text{patch}} = \sigma_{\text{patch}} \vec{E}_{\text{other}}$

Now  $\vec{E} = \vec{E}_{\text{other}} + \vec{E}_{\text{patch}}$

$\swarrow$  continuous       $\nwarrow$  discontinuous

$\Rightarrow \vec{E}_{\text{above}} = \vec{E}_{\text{other}} + \vec{E}_{\text{patch above}}$

$\Rightarrow \vec{E}_{\text{below}} = \vec{E}_{\text{other}} + \vec{E}_{\text{patch below}}$

or  $\vec{E}_{\text{above}} = \vec{E}_{\text{other}} + \frac{\sigma}{2\epsilon_0} \hat{m}$

$\vec{E}_{\text{below}} = \vec{E}_{\text{other}} - \frac{\sigma}{2\epsilon_0} \hat{m}$

$\Rightarrow \frac{1}{2} (\vec{E}_{\text{above}} + \vec{E}_{\text{below}}) = \vec{E}_{\text{other}}$

so  $\vec{f}_{\text{patch}} = \frac{1}{2} \sigma(\vec{r}) (\vec{E}_{\text{above}}(\vec{r}) + \vec{E}_{\text{below}}(\vec{r}))$

$\uparrow$  position of "patch"

where  $\vec{E}_{\text{other}}$  is continuous as move across patch, as it is field that would exist if a hole were cut out at the patch.

But  $\vec{E}_{\text{patch}}$  is discontinuous

$\vec{E}_{\text{patch above}} = \frac{\sigma}{2\epsilon_0} \hat{m}$  ← just above center of patch, it looks like infinite flat plane

$\vec{E}_{\text{patch below}} = -\frac{\sigma}{2\epsilon_0} \hat{m}$

For a conductor,  $\vec{E}_{\text{below}} = 0$ ,  $\vec{E}_{\text{above}} = \vec{E}$  on surface

$\Rightarrow \vec{f} = \frac{1}{2} \sigma \vec{E}$       but  $\vec{E} = \frac{\sigma}{\epsilon_0} \hat{m}$

$\Rightarrow \vec{f} = \frac{1}{2} \frac{\sigma^2}{\epsilon_0} \hat{m}$

or equivalently  $\vec{f} = \frac{\epsilon_0}{2} (\vec{E} \cdot \hat{m}) \vec{E} = \frac{\epsilon_0}{2} E^2 \hat{m}$

Electrostatic pressure

$p = \vec{f} \cdot \hat{m} = \frac{\epsilon_0}{2} E^2 = \frac{\sigma^2}{2\epsilon_0}$

as  $\vec{E} = E \hat{m}$

Note:  $\vec{f}$  is always <sup>outward</sup> along  $\hat{m}$  no matter what sign of  $\sigma$   
 i.e. force always draws conductor into the field.