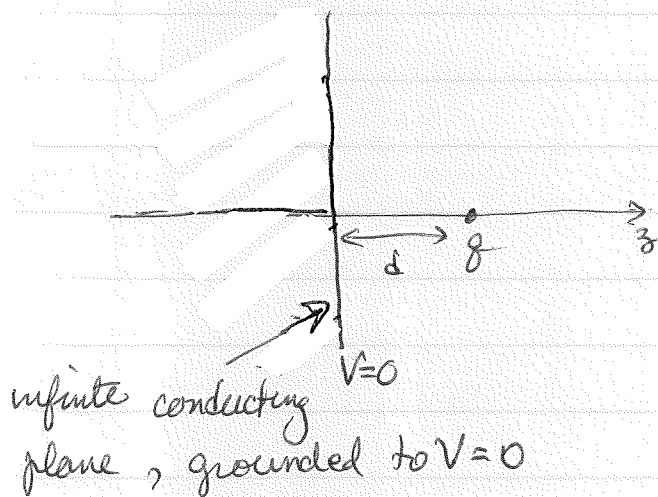


Image charge method



boundary value problem

$$\textcircled{1} \nabla^2 V = -\frac{\rho}{\epsilon_0} = -\frac{q}{\epsilon_0} \delta(\vec{r} - d\hat{z})$$

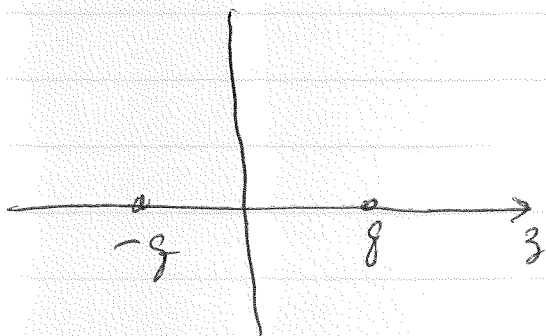
on right of plane

$$\textcircled{2} V=0 \text{ on the plane}$$

(i.e. on the boundary of the region to right of plane)

If we find a solution, it is the unique solution,

Solution: put fictitious "image charge" $-q$ at d to left of plane



$$V(\vec{r}) = \frac{q}{4\pi\epsilon_0} \frac{1}{|\vec{r} - d\hat{z}|} - \frac{q}{4\pi\epsilon_0} \frac{1}{|\vec{r} + d\hat{z}|}$$

$$\Rightarrow V(x, y, 0) = 0$$

$$= \frac{q}{4\pi\epsilon_0} \left\{ \frac{1}{\sqrt{x^2 + y^2 + (z-d)^2}} - \frac{1}{\sqrt{x^2 + y^2 + (z+d)^2}} \right\}$$

$$V(x, y, 0) = 0$$

$$\text{also } \nabla^2 V(\vec{r}) = -\frac{q}{\epsilon_0} \delta(\vec{r} - d\hat{z})$$

$$+ \frac{q}{\epsilon_0} \delta(\vec{r} + d\hat{z})$$

but $\delta(\vec{r} + d\hat{z}) = 0$ everywhere on right side plane

$$\Rightarrow \nabla^2 V = -\frac{\rho}{\epsilon_0} \text{ on right side}$$

So V is the solution

Since we have conducting plane, we can find electric field to right of plane, and use it to find surface charge induced on plane.

$$\vec{E} = -\vec{\nabla} V$$

$$\vec{E}(x, y, 0) = \frac{\sigma(x, y)}{\epsilon_0} \hat{m} = \frac{\sigma(x, y)}{\epsilon_0} \hat{z}$$

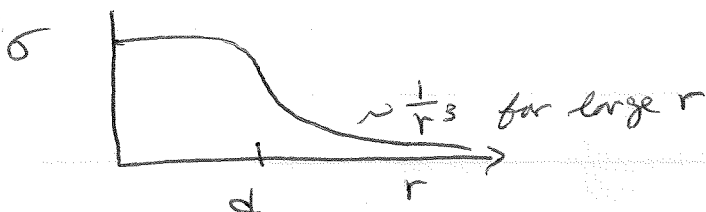
$$\Rightarrow \sigma(x, y) = \epsilon_0 E_z(x, y, 0) \\ = -\epsilon_0 \frac{\partial V(x, y, 0)}{\partial z}$$

$$E_z = -\frac{\partial V}{\partial z} = -\frac{q}{4\pi\epsilon_0} \left\{ \left(\frac{-1}{2}\right) \frac{2(z-d)}{[x^2+y^2+(z-d)^2]^{3/2}} - \left(\frac{-1}{2}\right) \frac{2(z+d)}{[x^2+y^2+(z+d)^2]^{3/2}} \right\}$$

$$E_z(x, y, 0) = +\frac{q}{4\pi\epsilon_0} \frac{-d-d}{[x^2+y^2+d^2]^{3/2}} = -\frac{qd}{2\pi\epsilon_0 [x^2+y^2+d^2]^{3/2}}$$

$$\sigma(x, y) = -\frac{qd}{2\pi} \frac{1}{[x^2+y^2+d^2]^{3/2}} = -\frac{qd}{2\pi} \frac{1}{[r^2+d^2]^{3/2}}$$

where $\sqrt{x^2+y^2} = r$ is radial dist in plane of conductor



total charge induced on the grounded plane is

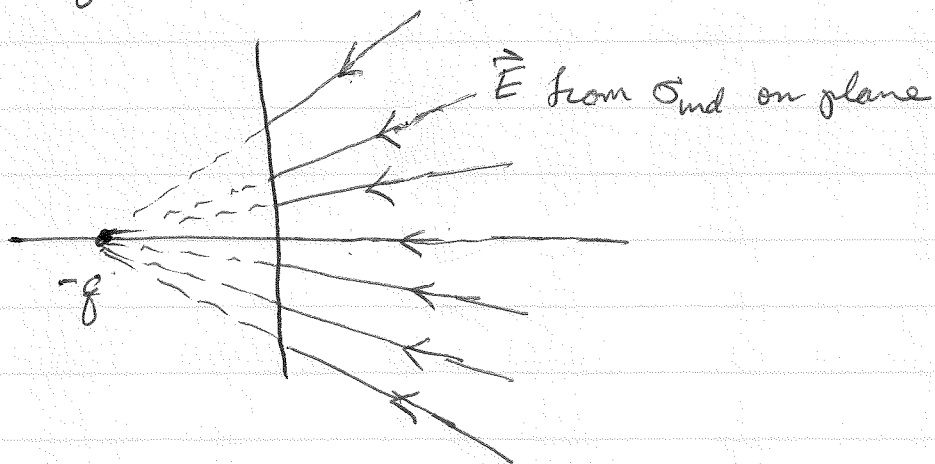
$$\begin{aligned} q_{ind} &= \int_{\text{plane}} da \sigma = \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy \sigma(x, y) = \int_0^{2\pi} d\phi \int_0^{\infty} dr r \sigma(r, \phi) \\ &= 2\pi \int_0^{\infty} dr \frac{(-gd)}{2\pi} \frac{r}{(r^2+d^2)^{3/2}} = -gd \int_0^{\infty} \frac{1}{(r^2+d^2)^{3/2}} \\ &= -gd \left[0 - \frac{-1}{d} \right] = -g \end{aligned}$$

That the plane is grounded at $V=0$ means that it is attached to a battery (charge reservoir) that can transfer or remove charge from the plane, so as to always keep its potential $V=0$, ~~is~~ with respect to pt $r \rightarrow \infty$. In present problem, battery transfers charge $-g$ to the conducting plane.

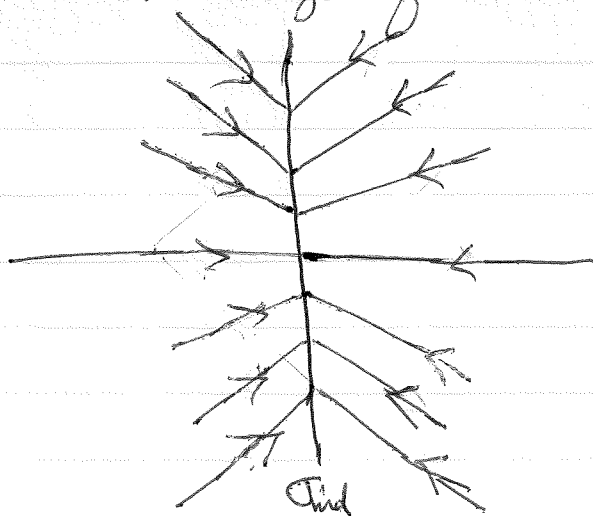
Image charge method solves for V only in region of interest, contained within the specified boundary. It does not give V on the outside of the boundary.

In the present example, we can find \vec{E} to left of plane, by using symmetry arguments

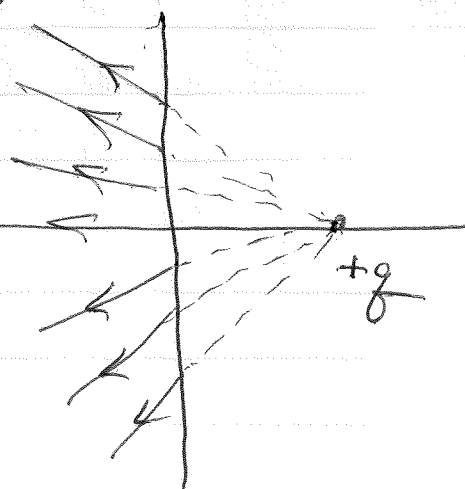
\vec{E} to right of plane, due to σ_{ind} on plane, is equal to field of image charge.



By symmetry, the \vec{E} on left of plane is just mirror image of \vec{E} on right of plane due to σ_{ind} (due to σ_{ind})

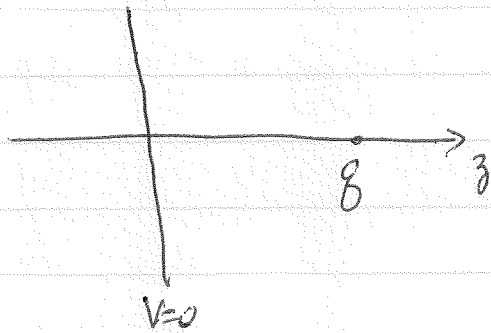


\vec{E} here looks just like field from pt charge $-q$ at $+d\hat{z}$



add to \vec{E} from real charge $+q$ at $d\hat{z}$
 \Rightarrow total \vec{E} on left of plane vanishes.

Force on charge in front of grounded conductor:



Force on q is force due to all other charges in system. In this case it is due to the σ_{ind} on the conducting plane

$$\vec{F} = q \vec{E}_{ind} \quad \text{but } \vec{E}_{ind} \text{ is same as field produced by a } -q \text{ at } -d\hat{z}$$

$$= q \left[\frac{-q}{4\pi\epsilon_0} \frac{(d\hat{z} - (-d\hat{z}))}{|d\hat{z} - (-d\hat{z})|^3} \right]$$

where $d\hat{z}$ is location of q and $-d\hat{z}$ is location of image

$$= \frac{-q^2 2d\hat{z}}{4\pi\epsilon_0 (2d)^3}$$

$$\vec{F} = \frac{-q^2}{16\pi\epsilon_0 d^2} \hat{z} \quad \text{attracts } q \text{ to the plane.}$$

Work done to move q in from infinity to the pt $\vec{r} = d\hat{z}$

$$W = - \int_{\infty}^d d\vec{l} \cdot \vec{F} = - \int_{\infty}^d dz F_z \quad \text{choose path along } z \text{ axis}$$

$$= \int_d^{\infty} dz \left(\frac{-q^2}{16\pi\epsilon_0} \right) \frac{1}{z^2} = \frac{q^2}{16\pi\epsilon_0} \left(\frac{1}{z} \right)_d^{\infty}$$

$$W = \frac{-q^2}{16\pi\epsilon_0 d}$$

$W < 0$ means energy is released as charge moves to plane - this is because electrostatic force is attractive

Note, W above is not the energy ~~of~~ that would be if the image charge was really present. The energy of a $+q$ and $-q$ separated by distance $2d$

$$\text{is } W = \frac{-q^2}{4\pi\epsilon_0 (2d)} = \frac{-q^2}{8\pi\epsilon_0 d} \text{ twice that above.}$$

we can explain the difference as follows:

$$W = \frac{q^2}{2\epsilon_0}$$

Energy of interaction