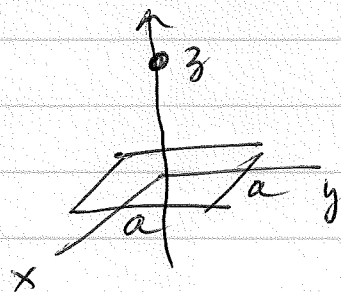
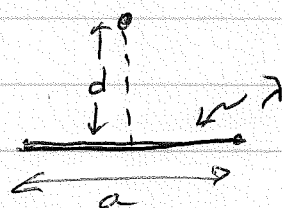


1) Field along z axis above square loop of side a with uniform line charge λ



compare to Ex 2.2
with $L \rightarrow \frac{a}{2}$
 $z \rightarrow d$

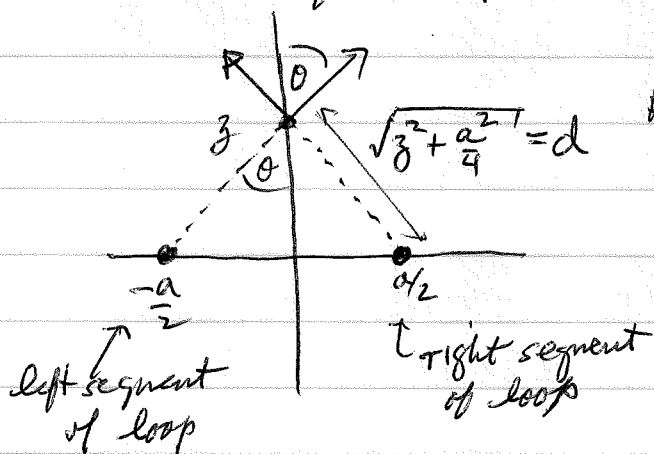
From lecture we saw that the \vec{E} field, at a point directly above the midpoint of a line of length a with uniform λ is



$$E = \frac{1}{4\pi\epsilon_0} \frac{\lambda a}{d\sqrt{d^2 + (\frac{a}{2})^2}}$$

\vec{E} points in direction from midpoint of line to observer

For the square loop, consider the y - z plane



Field from each of the two segments point as shown. We see the components $\perp z$ will cancel, while the z components add. So the net field above the square loop is in \hat{z} direction, and equal to 4 times the z -component of field from one side of the square

$$\cos\theta = \frac{z}{\sqrt{z^2 + \frac{a^2}{4}}}$$

The z component of \vec{E} from one side of square has magnitude

$$\frac{1}{4\pi\epsilon_0} \frac{\lambda a}{d\sqrt{d^2 + \frac{a^2}{4}}} \cos\theta$$

↑ project on \hat{z} axis

here the distance from wire to observer is

$$d = \sqrt{z^2 + \frac{a^2}{4}}$$

$$\text{and } \cos\theta = \frac{z}{\sqrt{z^2 + \frac{a^2}{4}}}$$

So

$$\begin{aligned} \vec{E}(z, z) &= \int 4 \frac{1}{4\pi\epsilon_0} \frac{\lambda a}{\sqrt{z^2 + \frac{a^2}{4}} \cdot \sqrt{z^2 + \frac{a^2}{4} + \frac{a^2}{4}}} \cdot \frac{z}{\sqrt{z^2 + \frac{a^2}{4}}} \\ &= \frac{1}{z} \frac{\lambda a}{\pi\epsilon_0} \frac{z}{\left(\frac{z^2 + \frac{a^2}{4}}{4}\right) \sqrt{z^2 + \frac{a^2}{4}}} \end{aligned}$$

Note: $E(z) = -E(-z)$ is antisymmetric as one would expect

Check: at $z=0$, $\vec{E}=0$

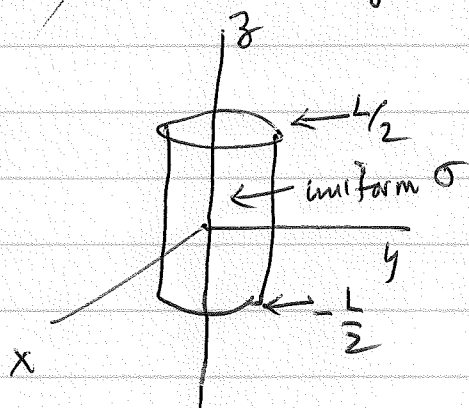
$$\text{For } z \gg a, \quad z^2 + \frac{a^2}{4} \approx z^2 + \frac{a^2}{2} \approx z^2$$

$$\vec{E} \approx \frac{1}{z} \frac{\lambda a}{\pi\epsilon_0} \frac{z}{z^2 \cdot z} = \frac{1}{z} \frac{4\lambda a}{4\pi\epsilon_0} \frac{1}{z^2}$$

Total charge is $Q = 4a\lambda$ so

$$\vec{E} = \frac{Q}{4\pi\epsilon_0} \frac{1}{z^2} \hat{z} \quad \text{as expected.}$$

2) Field along axis of cylindrical shell of charge of radius R , length L



$$\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_{\text{surface}} d\vec{a}' \sigma \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

$$= \frac{\sigma}{4\pi\epsilon_0} \int_{-\frac{L}{2}}^{\frac{L}{2}} dz' \int_0^{2\pi} d\phi R \frac{(z-z')\hat{z} - R\cos\phi\hat{x} - R\sin\phi\hat{y}}{[(z-z')^2 + R^2]^{3/2}}$$

$$\vec{r} = z\hat{z}$$

$$\vec{r}' = z'\hat{z} + R(\cos\phi\hat{x} + \sin\phi\hat{y})$$

when integrate over ϕ , $\int_0^{2\pi} d\phi \cos\phi = \int_0^{2\pi} d\phi \sin\phi = 0$

so only \hat{z} component does not vanish

$$\vec{E}(\hat{z}) = \hat{z} \frac{2\pi\sigma R}{4\pi\epsilon_0} \int_{-\frac{L}{2}}^{\frac{L}{2}} dz' \frac{(z-z')}{[(z-z')^2 + R^2]^{3/2}}$$

change integration variable to $\bar{z} = z - z'$ $d\bar{z} = -dz'$

$z' = (\frac{L}{2}, \frac{L}{2}) \rightarrow \bar{z} = (z + \frac{L}{2}, z - \frac{L}{2})$ flip limits of integration gives $(-)$

$$\vec{E}(\hat{z}) = \hat{z} \frac{2\pi\sigma R}{4\pi\epsilon_0} \int_{z-\frac{L}{2}}^{z+\frac{L}{2}} d\bar{z} \frac{\bar{z}}{(\bar{z}^2 + R^2)^{3/2}}$$

$$= \hat{z} \frac{2\pi\sigma R}{4\pi\epsilon_0} \left(\frac{-1}{(\bar{z}^2 + R^2)^{1/2}} \right)_{z-\frac{L}{2}}^{z+\frac{L}{2}}$$

Note: $E(z) = -E(-z)$ is antisymmetric as one would expect

$$\vec{E}(\hat{z}) = \hat{z} \frac{2\pi\sigma R}{4\pi\epsilon_0} \left[\frac{1}{\sqrt{(z-\frac{L}{2})^2 + R^2}} - \frac{1}{\sqrt{(z+\frac{L}{2})^2 + R^2}} \right]$$

Check limits: at $z=0$, $\vec{E}=0$ as expected by symmetry

What if $z \gg L$, $z \gg R$? Find \vec{E} to lowest non zero order in $\frac{L}{z}$ and $\frac{R}{z}$.

$$\frac{1}{\sqrt{(z-\frac{L}{2})^2 + R^2}} = \frac{1}{z} \frac{1}{\sqrt{(1-\frac{L}{2z})^2 + \frac{R^2}{z^2}}} = \frac{1}{z} \sqrt{1 - \frac{L}{z} + \frac{L^2}{4z^2} + \frac{R^2}{z^2}}$$

to lowest order in small quantities $\frac{L}{z}$, $\frac{R}{z}$

$$\approx \frac{1}{z} \frac{1}{(1 - \frac{L}{2z})} \approx \frac{1}{z} \left(1 + \frac{L}{2z}\right) \quad \text{ignore } \frac{L^2}{z^2} \text{ and } \frac{R^2}{z^2} \text{ terms}$$

similarly

$$\frac{1}{\sqrt{(z+\frac{L}{2})^2 + R^2}} \approx \frac{1}{z} \left(1 - \frac{L}{2z}\right)$$

$$\text{So } \vec{E}(z\hat{z}) \approx \frac{\hat{z}}{z} \frac{2\pi\sigma R}{4\pi\epsilon_0} \frac{1}{z} \left[\left(1 + \frac{L}{2z}\right) - \left(1 - \frac{L}{2z}\right) \right]$$

$$\approx \frac{\hat{z}}{z} \frac{2\pi\sigma R}{4\pi\epsilon_0} \frac{1}{z} \left[\frac{L}{z} \right] = \frac{\hat{z}}{z} \frac{2\pi R L \sigma}{4\pi\epsilon_0 z^2}$$

Total charge on cylindrical shell is $\sigma(\text{area}) = 2\pi R L \sigma = Q$

$$\vec{E}(z\hat{z}) = \frac{Q}{4\pi\epsilon_0 z^2} \hat{z} \quad \text{just like point charge } Q \text{ when } \vec{r} \text{ is far away!}$$