

## From Coulomb to Maxwell

Electrodynamics is concerned with one particular attribute of matter - charge

Experimentally it was observed that certain bodies exert long range forces on each other that are certainly not gravitational - they are not proportional to the mass and they can be repulsive as well as attractive. The source of this new force was defined to be the "charge" of the object

### Electrostatics

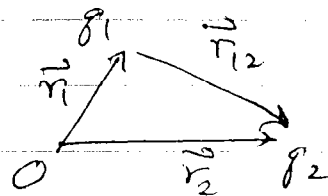
Coulomb's Law - for charge  $q_1$  at  $\vec{r}_1$  and charge  $q_2$  at  $\vec{r}_2$ , if separation  $|\vec{r}_2 - \vec{r}_1|$  is much greater than the size of either charge, then

$$\vec{F}_{12} = k_1 q_1 q_2 \frac{\hat{r}_{12}}{r_{12}^2}$$

↑  
force on 2 due to 1

$$\begin{aligned} \vec{r}_{12} &= \vec{r}_2 - \vec{r}_1 \\ \hat{r}_{12} &= \frac{\vec{r}_{12}}{|\vec{r}_{12}|} \end{aligned}$$

central force - points from 1 to 2  
inverse square law



(2)

$k_1$  is a universal constant of nature that determines the strength of the force when  $q$  is expressed in terms of some arbitrary reference charge.

Since we only know about charge by measuring the Coulomb forces, we are in principle free to choose  $k_1$  to be anything we like - our choice then determines the units that charge is measured in.

In MKS system of units (same as SI system) charge is measured in the historical unit, the "coulomb." Then  $k_1$  has the value  $k_1 = \frac{1}{4\pi\epsilon_0} = 10^{-7} \text{ C}^2$ , where  $c$  is speed of light in a vacuum. The units of  $k_1$  are  $\text{Nt} \cdot \text{m}^2 / \text{coul}^2$

In CGS system of units (abt called esu - electrostatic units) one fixes  $k_1 = 1$  and charge is measured in "statcoulombs"  $k_1$  is taken dimensionless, so  $\text{statcoulomb} = (\text{Nt} \cdot \text{m}^2)^{1/2}$

Another reasonable modern choice would be to measure charge in integer multiples of the electron charge. This would yield a different value for  $k_1$ .

In this class we will be using CGS units. But we keep  $k_1$  general for now.

### Superposition

For charges  $q_i$  at positions  $\vec{r}_i$ , the force on charge  $Q$  at position  $\vec{r}$  is

$$\vec{F} = k_1 Q \sum_i q_i \frac{\vec{r} - \vec{r}_i}{|\vec{r} - \vec{r}_i|^3} \quad \text{forces add linearly}$$

Conservation of charge - charge is neither created nor destroyed

$$\frac{d}{dt} \sum_i q_i = 0 \quad \text{where sum is over all charges in system}$$

### Continuum charge density

for charges  $q_i$  at positions  $\vec{r}_i$ , define,

$$\rho(\vec{r}) = \sum_i q_i \delta(\vec{r} - \vec{r}_i)$$

$\delta(\vec{r} - \vec{r}_i)$  is Dirac  $\delta$ -function with properties:

$$\int_V d^3r \delta(\vec{r} - \vec{r}_i) = \begin{cases} 1 & \text{if } \vec{r}_i \in V \\ 0 & \text{otherwise} \end{cases}$$

$$\int_V d^3r f(\vec{r}) \delta(\vec{r} - \vec{r}_i) = \begin{cases} f(\vec{r}_i) & \text{if } \vec{r}_i \in V \\ 0 & \text{otherwise} \end{cases}$$

for any scalar function  $f(\vec{r})$

$$\int_V d^3r \rho(\vec{r}) = \sum_i q_i \int d^3r \delta(\vec{r} - \vec{r}_i)$$

=  $Q_{\text{enc}}$  total charge enclosed by volume  $V$

$\Rightarrow \rho$  has units of charge per volume

$\Rightarrow \delta(\vec{r})$  has units of  $1/\text{vol}$

Coulomb

$$\vec{E} = k_1 Q \int d^3r' \rho(\vec{r}') \frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3}$$

We will often forget that  $\rho(\vec{r})$  is in principle made up of a distribution of point charges, and take it to be a smooth continuous function.

charge conservation -  $\frac{d}{dt} \int_V d^3r \rho(\vec{r}) = 0$

assuming  $V$  is so big that it contains all the charge, and no charge flows through the surface of  $V$

# Electric Field

$\vec{E}(\vec{r})$  is the force per unit charge that would be felt by an infinitesimal test charge  $sq$  at position  $\vec{r}$ .

$$\vec{E}(\vec{r}) = \frac{1}{sq} \vec{F} = k_1 \int d^3r' \rho(\vec{r}') \frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3} \quad (*)$$

In principle, the above is solution to all electrostatic problems. In practise, we may not always know  $\rho(\vec{r})$ , but may need to solve for it self consistently with  $\vec{E}$ . It will help to have another formulation of the above in terms of differential equations. We get these by taking the divergence and curl of eq (\*) above to get (see proof later)

$$\Rightarrow \left\{ \begin{array}{l} \vec{\nabla} \cdot \vec{E} = 4\pi k_1 \rho \quad (1) \quad \text{Gauss' Law} \\ \vec{\nabla} \times \vec{E} = 0 \quad (2) \quad \leftarrow \text{true only for statics!} \end{array} \right.$$

the proof of the above will follow on next page.

we also can recast (1) and (2) in integral form as follows

by Gauss' Theorem  $\int_V d^3r \vec{\nabla} \cdot \vec{E} = \oint_S da \hat{n} \cdot \vec{E} = 4\pi k_1 \int_V d^3r \rho$

total charge enclosed in V

by Stokes Theorem  $\int_S da \hat{n} \cdot \vec{\nabla} \times \vec{E} = \oint_C d\vec{l} \cdot \vec{E} = 0$

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in above,  $\hat{n}$  is outward pointing normal to surface  $S$   
 $d\vec{l}$  is differential tangent to curve  $C$  bounding surface  $S$

Proof that (1) and (2) follow from (\*)

First Note that 
$$\frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3} = -\vec{\nabla} \left( \frac{1}{|\vec{r}-\vec{r}'|} \right)$$

to see this, let  $\vec{r} \equiv \vec{r}-\vec{r}'$ , and do the calculation in spherical coords centered at  $\vec{r}=0$ .

$$\frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3} = \frac{\hat{r}}{r^2}$$

$$-\vec{\nabla} \left( \frac{1}{|\vec{r}-\vec{r}'|} \right) = -\vec{\nabla} \frac{1}{r}$$

since  $\hat{r} = \vec{r}-\vec{r}'$ , we have  $\vec{\nabla}$  (which differentiates with respect to  $\vec{r}$ ) is the same as  $\vec{\nabla}_r$  (which differentiates with respect to  $r$ )

$$\begin{aligned} \text{So } -\vec{\nabla} \left( \frac{1}{r} \right) &= \hat{r} \frac{d}{dr} \left( \frac{1}{r} \right) \quad \text{using } \vec{\nabla} \text{ in spherical coords} \\ &= -\frac{\hat{r}}{r^2} \end{aligned}$$

So from Coulomb we have

$$\vec{E}(\vec{r}) = k_1 \int d^3r' \rho(\vec{r}') \frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3}$$

$$= -\vec{\nabla} \left( k_1 \int d^3r' \frac{\rho(\vec{r}')}{|\vec{r}-\vec{r}'|} \right) \quad \vec{E} \text{ is gradient of scalar function}$$

$$\Rightarrow \boxed{\vec{\nabla} \times \vec{E} = 0}$$

since the curl of a gradient always vanishes  
 $\vec{\nabla} \times \vec{\nabla} \phi = 0$  for any scalar function  $\phi$

(7)

$$\vec{\nabla} \cdot \vec{E} = -\nabla^2 \left( k_1 \int d^3r' \frac{\rho(r')}{|\vec{r} - \vec{r}'|} \right) \quad \text{where } \nabla^2 = \vec{\nabla} \cdot \vec{\nabla}$$

Consider

$$\nabla^2 \left( \frac{1}{|\vec{r} - \vec{r}'|} \right)$$

as before, define  $\vec{r} = \vec{r} - \vec{r}'$ , so  $\vec{\nabla}_r = \vec{\nabla}_r$ , and go to spherical coords centered at  $\vec{r} = 0$ .

$$\nabla^2 \left( \frac{1}{|\vec{r} - \vec{r}'|} \right) = \nabla_r^2 \left( \frac{1}{r} \right) \quad \text{use expression for } \nabla^2 \text{ in spherical coords}$$

$$= \frac{1}{r} \frac{d^2}{dr^2} r \left( \frac{1}{r} \right)$$

$$= \begin{cases} 0! & \text{for } r \neq 0 \\ \text{singular} & \text{at } r = 0 \end{cases}$$

So

$\nabla^2 \left( \frac{1}{r} \right)$  vanishes everywhere except at  $r = 0$

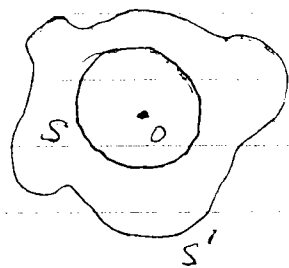
to see what happens at  $r = 0$ , consider integrating over a sphere  $S$  of radius  $R$  centered at the origin

$$\int_V d^3r \nabla^2 \left( \frac{1}{r} \right) = \int_V d^3r \vec{\nabla} \cdot \vec{\nabla} \left( \frac{1}{r} \right) = \oint_S da \hat{n} \cdot \vec{\nabla} \left( \frac{1}{r} \right) \quad \text{using Gauss' Theorem}$$

integrate  $\hat{n} \cdot \vec{\nabla} \left( \frac{1}{r} \right) = \frac{d}{dr} \left( \frac{1}{r} \right) = -\frac{1}{r^2}$  is constant on surface  $S$  so

$$\oint_S da \hat{n} \cdot \vec{\nabla} \left( \frac{1}{r} \right) = 4\pi R^2 \left( \frac{-1}{R^2} \right) = -4\pi$$

Above was integrating over a sphere, but we would get same result if integrated over any volume containing  $\vec{r}=0$ .



S is sphere of radius R

S' is any surface

let  $V'$  be volume between S and S'

Then by Gauss theorem

$$\int_{V'} d^3r \nabla \cdot \vec{v}(\frac{1}{r}) = \oint_{S'} da \hat{n} \cdot \vec{v}(\frac{1}{r}) - \oint_S da \hat{n} \cdot \vec{v}(\frac{1}{r})$$

$$= 0 \text{ since } \nabla^2(\frac{1}{r}) = 0 \text{ everywhere in } V'$$

$$\Rightarrow \oint_{S'} da \hat{n} \cdot \vec{v}(\frac{1}{r}) = \oint_S da \hat{n} \cdot \vec{v}(\frac{1}{r})$$

$$\Rightarrow \int_{V'} d^3r \nabla^2(\frac{1}{r}) = \int_V d^3r \nabla^2(\frac{1}{r})$$

$V'$  bounded by  $S'$        $V$  bounded by  $S$

So we conclude:  $\int_V d^3r \nabla^2(\frac{1}{r}) = \begin{cases} -4\pi & \text{if } \vec{r}=0 \text{ in } V \\ 0 & \text{if } \vec{r}=0 \text{ not in } V \end{cases}$

$\Rightarrow \nabla^2(\frac{1}{r}) = -4\pi \delta(\vec{r})$  Dirac delta function

$$\nabla^2\left(\frac{1}{|\vec{r}-\vec{r}'|}\right) = -4\pi \delta(|\vec{r}-\vec{r}'|)$$



So now

$$\vec{\nabla} \cdot \vec{E} = -k_1 \int d^3r' \rho(r') \nabla^2 \left( \frac{1}{|\vec{r} - \vec{r}'|} \right)$$

acts only on  $\vec{r}$

$$= -k_1 \int d^3r' \rho(r') (-4\pi) \delta(\vec{r} - \vec{r}')$$

$$= 4\pi k_1 \rho(\vec{r}) \quad \text{by property of } \delta \text{-function}$$

proof is done!

we have shown that

$$\vec{E}(\vec{r}) = k_1 \int d^3r' \rho(\vec{r}') \frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3} \Rightarrow \begin{cases} \vec{\nabla} \cdot \vec{E} = 4\pi k_1 \rho \\ \vec{\nabla} \times \vec{E} = 0 \end{cases}$$

is the reverse true? i.e. is the formulation in terms of partial differential equations completely equivalent to Coulomb's law? yes! because of Helmholtz's Theorem.

Helmholtz Theorem of vector calculus — if one specifies the divergence and curl of a vector function, and boundary conditions (here  $E \rightarrow 0$  as  $r \rightarrow \infty$  and one is away from all charges), then vector function is uniquely determined.