

Radiation from a Localized Oscillating Source

$$\vec{A}(\vec{r}, t) = \frac{1}{c} \int d^3r' dt' \frac{\delta(t-t' - \frac{|\vec{r}-\vec{r}'|}{c})}{|\vec{r}-\vec{r}'|} \vec{j}(\vec{r}', t')$$

For pure harmonic oscillation in current

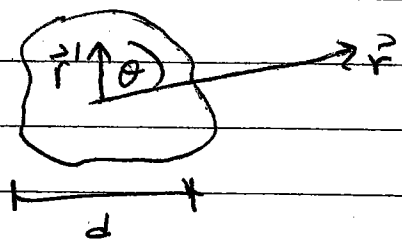
$$\vec{j}(\vec{r}', t) = \text{Re} \{ \vec{j}_\omega(\vec{r}') e^{-i\omega t} \}$$

$$\Rightarrow \vec{A}(\vec{r}, t) = \text{Re} \{ \vec{A}_\omega(\vec{r}) e^{-i\omega t} \}$$

$$\Rightarrow \vec{A}_\omega(\vec{r}) e^{-i\omega t} = \frac{1}{c} \int d^3r' \vec{j}_\omega(\vec{r}') e^{-i\omega t} \frac{e^{i\omega |\vec{r}-\vec{r}'|/c}}{|\vec{r}-\vec{r}'|}$$

$$\vec{A}_\omega(\vec{r}) = \frac{1}{c} \int d^3r' \vec{j}_\omega(\vec{r}') \frac{e^{i\omega |\vec{r}-\vec{r}'|/c}}{|\vec{r}-\vec{r}'|}$$

Assume source is localized, i.e. $\vec{j}_\omega(\vec{r}') \approx 0$ for $|\vec{r}'| > d$



Approx ①

for $r \gg d$, far from sources

$$\begin{aligned} |\vec{r}-\vec{r}'| &= \sqrt{r^2 + r'^2 - 2rr' \cos \theta} \\ &= r \sqrt{1 + \left(\frac{r'}{r}\right)^2 - \frac{2r'}{r} \cos \theta} \\ &\approx r \left(1 - \frac{r'}{r} \cos \theta\right) \\ &= r - \vec{r}' \cdot \hat{r} + o\left(\left(\frac{r'}{r}\right)^2\right) \\ &\quad \hat{r} \equiv \frac{\vec{r}}{r} \end{aligned}$$

$$\vec{A}_\omega(\vec{r}) = \frac{1}{c} \int d^3r' \frac{\vec{j}_\omega(\vec{r}') e^{ik(r-\vec{r}'\cdot\hat{r})}}{r-\vec{r}'\cdot\hat{r}} \quad \text{where } k \equiv \frac{\omega}{c}$$

$$= \frac{e^{ikr}}{cr} \int d^3r' \frac{\vec{j}_\omega(\vec{r}') e^{-ik\vec{r}'\cdot\hat{r}}}{1 - \frac{\hat{r}\cdot\vec{r}'}{r}}$$

$$\approx \frac{e^{ikr}}{cr} \int d^3r' \vec{j}_\omega(\vec{r}') e^{-ik\hat{r}\cdot\vec{r}'} \left(1 + \frac{\hat{r}\cdot\vec{r}'}{r}\right)$$

↑ when combine with the $e^{-i\omega t}$ piece, this gives outgoing spherical wave $\frac{e^{i(kr-\omega t)}}{r}$

oscillating charge radiates outgoing spherical electromagnetic waves

the $\int d^3r' \vec{j}_\omega(\vec{r}') \dots$ term will determine the angular dependence of the radiation.

Approx ② $\lambda \gg d$ long wave length approx

$$\text{or } kd \ll 1 \Rightarrow \frac{\omega}{c} d \ll 1 \text{ or } \frac{d}{\tau} \ll c$$

where τ is period of oscillation.

Since $\frac{d}{\tau}$ is max speed of the oscillating charges $\rightarrow \lambda \gg d$ is a non-relativistic approximation

$$kd \ll 1 \Rightarrow e^{-ik\hat{r} \cdot \vec{r}'} \approx 1 - ik\hat{r} \cdot \vec{r}' + \text{higher orders}$$

$$\vec{A}_\omega(\vec{r}) = \frac{e^{ikr}}{cr} \int d^3r' \vec{j}_\omega(\vec{r}') (1 - ik\hat{r} \cdot \vec{r}') \left(1 + \frac{\hat{r} \cdot \vec{r}'}{r}\right)$$

$$= \frac{e^{ikr}}{cr} \int d^3r' \vec{j}_\omega(\vec{r}') \left[1 + \hat{r} \cdot \vec{r}' \left(\frac{1}{r} - ik\right) \right]$$

+ higher order in $\frac{d}{r}$ or kd

$$\vec{A}_\omega(\vec{r}) = \frac{e^{ikr}}{r} \left[-\vec{I}_1 + \left(\frac{1}{r} - ik\right) \vec{I}_2 \right]$$

where $\vec{I}_1 \equiv \frac{1}{c} \int d^3r' \vec{j}_\omega(\vec{r}')$

$$\vec{I}_2 \equiv \frac{1}{c} \int d^3r' \hat{r} \cdot \vec{r}' \vec{j}_\omega(\vec{r}')$$

Consider first \vec{I}_1 i th component (\vec{I}_1 vanishes in statics)

$$\int d^3r j_i(\vec{r}) = - \int d^3r r_i \vec{\nabla} \cdot \vec{j} \quad \text{integration by parts}$$

$$= \int d^3r r_i \frac{\partial \rho}{\partial t} \quad \text{as } \vec{\nabla} \cdot \vec{j} + \frac{\partial \rho}{\partial t} = 0$$

$$\int d^3r j_i(\vec{r}) = -i\omega \int d^3r r_i \rho_\omega(\vec{r})$$

$$\Rightarrow \vec{I}_1 = -\frac{i\omega}{c} \int d^3r \vec{r} \rho_\omega(\vec{r}) = -\frac{i\omega}{c} \vec{P}_\omega$$

electric dipole moment

Electric dipole approximation from \vec{I}_1

$$\vec{A}_{E1}(\vec{r}) = \frac{e^{-ikr}}{r} (-i\omega \vec{p}_\omega) = -i \vec{p}_\omega \frac{k e^{-ikr}}{r} \quad \omega = ck$$

Consider \vec{I}_2

$$\vec{I}_2 = \frac{1}{c} \int d^3r' \hat{r} \cdot \vec{r}' \vec{j}_\omega(\vec{r}') = \frac{1}{c} \hat{r} \cdot \int d^3r' \vec{r}' \vec{j}_\omega(\vec{r}')$$

we saw this tensor earlier when we did the ^{tensor} magnetic dipole approx, and when we derived the macroscopic Maxwell equations

$$\begin{aligned} \int d^3r' \vec{r}' \vec{j}_\omega(\vec{r}') &= - \int d^3r' \vec{j}_\omega(\vec{r}') \vec{r}' - \int d^3r' \vec{r}' \vec{r}' (\vec{\nabla}' \cdot \vec{j}_\omega(\vec{r}')) \\ &= \frac{1}{2} \int d^3r' [\vec{r}' \vec{j}_\omega - \vec{j}_\omega \vec{r}'] - \frac{1}{2} \int d^3r' \epsilon \omega \vec{r}' \vec{r}' \rho_\omega \end{aligned}$$

using $\vec{\nabla}' \cdot \vec{j} = -\frac{\partial \rho}{\partial t}$

$$\begin{aligned} \vec{I}_2 &= \frac{1}{2c} \int d^3r' [(\hat{r} \cdot \vec{r}') \vec{j}_\omega - (\hat{r} \cdot \vec{j}_\omega) \vec{r}'] - \frac{1}{2} \frac{\epsilon \omega}{c} \hat{r} \cdot \int d^3r' (\vec{r}' \vec{r}') \rho_\omega(\vec{r}') \\ &= -\frac{1}{2c} \int d^3r' [\hat{r} \times (\vec{r}' \times \vec{j}_\omega)] - \frac{1}{2} \frac{\epsilon \omega}{c} \hat{r} \cdot \int d^3r' (\vec{r}' \vec{r}') \rho_\omega(\vec{r}') \\ &= -\hat{r} \times \vec{m}_\omega - \frac{1}{2} \frac{i\omega}{3c} \hat{r} \cdot \vec{Q}_\omega \end{aligned}$$

where $\vec{m}_\omega = \frac{1}{2c} \int d^3r' \vec{r}' \times \vec{j}_\omega(\vec{r}')$ is magnetic dipole moment

$\vec{Q}_\omega = \int d^3r' \epsilon \vec{r}' \vec{r}' \rho_\omega(\vec{r}')$ looks almost like electric quadrupole tensor

to make it look like the proper quadrupole moment

$$\vec{Q}_\omega = \int d^3r' (3\vec{r}'\vec{r}' - r'^2 \mathbf{I}) \rho_\omega(\vec{r}')$$

we can write

$$\vec{Q}'_\omega = \vec{Q}_\omega + \mathbf{I} \int d^3r' r'^2 \rho_\omega(\vec{r}')$$

↑ identity matrix $I_{ij} = \delta_{ij}$

$$\vec{I}_2 = -\hat{r} \times \vec{m}_\omega - \frac{1}{2} \frac{i\omega}{3c} \hat{r} \cdot \vec{Q}_\omega - \frac{i\omega}{6c} \hat{r} C_\omega$$

$$\text{where } C_\omega \equiv \int d^3r' r'^2 \rho_\omega(\vec{r}')$$

is a scalar

Magnetic dipole approximation from \vec{I}_2

$$\vec{A}_{M1}(\vec{r}) = \frac{e^{i\vec{k}\vec{r}}}{r} \left(\frac{1}{r} - ik \right) (-\hat{r} \times \vec{m}_\omega)$$

Electric quadrupole approximation from \vec{I}_2

$$\vec{A}_{E2}(\vec{r}) = \frac{e^{i\vec{k}\vec{r}}}{r} \left(\frac{1}{r} - ik \right) \left(-\frac{i\omega}{6c} \hat{r} \cdot \vec{Q}_\omega \right)$$

The last piece $\frac{e^{i\vec{k}\vec{r}}}{r} \left(\frac{1}{r} - ik \right) \left(-\frac{i\omega}{6c} \hat{r} C_\omega \right)$

can always be ignored - it is a radial function and so its curl always vanishes \rightarrow gives no contribution to \vec{B} . Similarly, since $-\frac{i\omega}{c} \vec{E}_\omega = c\vec{k} \times \vec{B}_\omega$ by Ampere's law, this term will give no contribution to \vec{E} .

So with these two approximations ① and ②

$$\vec{A}_\omega(\vec{r}) = \vec{A}_{E1}(\vec{r}) + \vec{A}_{M1}(\vec{r}) + \vec{A}_{E2}(\vec{r})$$

keeping higher order terms would give magnetic quadrupole, electric octopole etc.

Compare strengths of the terms above

Approx ③

Radiation Zone: far from sources,

$\frac{1}{r} \ll k$ so $(\frac{1}{r} - ik) \approx -ik$ in \vec{A}_{M1} and \vec{A}_{E2}
only keep terms of $O(\frac{1}{r})$

electric dipole

$$\vec{p}_\omega \sim qd$$

$$\vec{A}_{E1} \sim qkd$$

magnetic dipole

$$\vec{m}_\omega \sim \frac{v}{c} qd$$

$$\vec{A}_{M1} \sim qkd \left(\frac{v}{c}\right)$$

$$\text{use } v \sim \frac{d}{\tau} \sim d\omega \sim dk c \Rightarrow \vec{A}_{M1} \sim q(kd)^2$$

electric quadrupole

$$\vec{Q}_\omega \sim qd^2$$

$$\vec{A}_{E2} \sim qd^2 k \frac{\omega}{c} \sim q(kd)^2$$

Since Approx ② assumed $kd \approx \frac{v}{c} \ll 1$

above is expansion in powers of kd

leading term is electric dipole

next order are

[magnetic dipole
electric quadrupole]

$$\frac{A_{M1}}{A_{E1}} \sim \frac{A_{E2}}{A_{E1}} \sim kd$$

next order terms are smaller than A_{E1} by factor $(kd)^2$
etc.

Electric Dipole Approximation - The leading term in non-relativistic expansion

$$\vec{A}_{E1}(\vec{r}) = -ik \vec{p}_\omega \frac{e^{ikr}}{r}$$

$$\vec{\nabla} \times (\phi \vec{F}) = (\vec{\nabla} \phi) \times \vec{F} + \phi \vec{\nabla} \times \vec{F}$$

$$\vec{B}_{E1} = \vec{\nabla} \times \vec{A}_{E1} = -ik \left(\vec{\nabla} \frac{e^{ikr}}{r} \right) \times \vec{p}_\omega$$

$$= -ik \left(ik - \frac{1}{r} \right) \frac{e^{ikr}}{r} \hat{r} \times \vec{p}_\omega$$

$$= k^2 \frac{e^{ikr}}{r} \left(1 + \frac{i}{kr} \right) \hat{r} \times \vec{p}_\omega$$

In radiation zone approx, $kr \gg 1$

$$\vec{B}_{E1} \approx k^2 \frac{e^{ikr}}{r} \hat{r} \times \vec{p}_\omega$$

To get electric field, use Ampere's Law

$$\vec{\nabla} \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t} \quad (\text{away from source where } \vec{J} = 0)$$

For oscillating fields $\vec{E} = E_\omega e^{-i\omega t}$

$$\vec{\nabla} \times \vec{B}_\omega = -\frac{i\omega}{c} \vec{E}_\omega \Rightarrow E_{E1} = \frac{i}{k} \vec{\nabla} \times \vec{B}_{E1}$$

$$\vec{E}_{E1} = \frac{i}{k} \vec{\nabla} \times \left[k^2 \frac{e^{ikr}}{r} \left(1 + \frac{i}{kr} \right) \hat{r} \times \vec{p}_\omega \right]$$

Ignore in RZ approx

$$\vec{E}_{E1} = \frac{i}{k} (\vec{\nabla} e^{ikr}) \times \left[\frac{k^2}{r} \left(1 + \frac{i}{kr}\right) \hat{r} \times \vec{p}_\omega \right] + \frac{i}{k} e^{ikr} \vec{\nabla} \times \left[\frac{k^2}{r} \left(1 + \frac{i}{kr}\right) \hat{r} \times \vec{p}_\omega \right]$$

this will always be of order $1/r^2$
So can ignore it in radiation zone approx

So in radiation zone approx

$$\vec{E}_{E1} = (\vec{\nabla} e^{ikr}) \times \left[\frac{ik}{r} \hat{r} \times \vec{p}_\omega \right]$$

$$\vec{E}_{E1} = -k^2 \frac{e^{ikr}}{r} \hat{r} \times (\hat{r} \times \vec{p}_\omega)$$

if do not make radiation zone approx, one gets

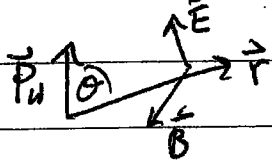
$$\vec{E}_{EI} = \frac{k^2 e^{ikr}}{r} \left[\vec{p}_\omega - \hat{r}(\vec{p}_\omega \cdot \hat{r}) - \frac{i}{kr} \left(1 + \frac{i}{kr}\right) (3\hat{r}(\vec{p}_\omega \cdot \hat{r}) - \vec{p}_\omega) \right]$$

Using radiation zone approx:

$$\vec{E}_{EI} = k^2 \frac{e^{ikr}}{r} \hat{r} \times (\vec{p}_\omega \times \hat{r}) \quad |\vec{E}_{EI}| = |\vec{B}_{EI}|$$

$$\vec{B}_{EI} = -k^2 \frac{e^{ikr}}{r} \vec{p}_\omega \times \hat{r} \quad \vec{E}_{EI} \perp \vec{B}_{EI}$$

If choose coordinates so that \vec{p}_ω is along \hat{z} axis, then



$$\vec{E}_{EI} = -k^2 p_\omega \frac{e^{ikr}}{r} \sin\theta \hat{\theta}$$

$$\vec{B}_{EI} = -k^2 p_\omega \frac{e^{ikr}}{r} \sin\theta \hat{\phi}$$

Emitted power

radiating vector $\vec{S}_{EI}(\vec{r}, t) = \frac{c}{4\pi} \text{Re}\{\vec{E}_{EI}\} \times \text{Re}\{\vec{B}_{EI}\}$

need to take real parts of complex expressions before multiplying

$$\text{Re}\{\vec{E}_{EI}(\vec{r}, t)\} = -k^2 p_\omega \frac{\cos(kr - \omega t)}{r} \sin\theta \hat{\theta}$$

$$\text{Re}\{\vec{B}_{EI}(\vec{r}, t)\} = -k^2 p_\omega \frac{\cos(kr - \omega t)}{r} \sin\theta \hat{\phi}$$

$$\vec{S}_{EI}(\vec{r}, t) = \frac{c}{4\pi} k^4 p_\omega^2 \frac{\cos^2(kr - \omega t)}{r^2} \sin^2\theta \hat{r}$$

$\vec{S}_{EI} \sim \hat{r} \Rightarrow$ energy flows radially outwards

$\vec{S}_{EI} \sim \frac{1}{r^2} \Rightarrow$ energy conserved

$$\oint da \hat{m} \cdot \vec{S}_{EI} = \text{constant for all } R$$

sphere
radius R

Question - what about the

time averaged energy current

$\frac{1}{r^n}, n > 2$, terms if we do not
make radiation zone approx?

$$\langle \vec{S}_{EI} \rangle = \frac{1}{T} \int_0^T dt \vec{S}_{EI}(\vec{r}, t)$$

$$T \text{ is period } T = \frac{2\pi}{\omega}$$

$$\langle \cos^2(\gamma) \rangle = \frac{1}{2}$$

$$= \frac{c}{8\pi} k^4 P_w^2 \frac{\sin^2 \theta}{r^2} \hat{r}$$

average energy flowing through an element
of area at spherical angles θ, ϕ is

$$dP_{EI} = \hat{r} \cdot \langle \vec{S}_{EI} \rangle \underbrace{r^2 \sin \theta d\theta d\phi}_{\text{area of surface element}}$$

area of surface element

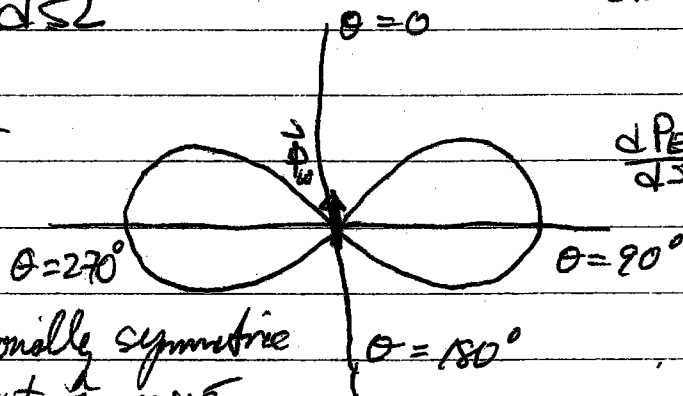
$$= r^2 d\Omega$$

Ω is solid angle

$$= \hat{r} \cdot \langle \vec{S}_{EI} \rangle r^2 d\Omega$$

$$\frac{dP_{EI}}{d\Omega} = \hat{r} \cdot \langle \vec{S}_{EI} \rangle r^2 = \frac{c}{8\pi} k^4 P_w^2 \sin^2 \theta \sim \omega^4 \sin^2 \theta$$

polar plot



rotationally symmetric
about $\frac{1}{2} \omega \times 15$

$$\frac{dP_{EI}}{d\Omega} \sim \sin^2 \theta$$

most of power is
directed outwards
into plane $\perp \vec{P}_w$,
i.e. peaked about $\theta = 90$

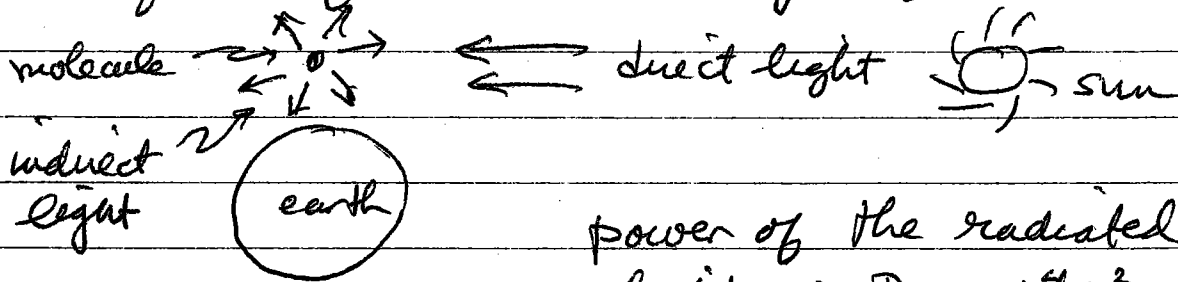
Total power radiated is

$$P_{EI} = \int \frac{dP_{EI}}{d\Omega} d\Omega = \frac{ck^4 p_0^2}{8\pi} 2\pi \int_0^\pi \sin\theta \sin^2\theta$$

$$P_{EI} = \frac{ck^4 p_0^2}{3} = \frac{p_0^2 \omega^4}{3c^3} \sim \omega^4$$

why the sky is blue - Lord Rayleigh

when look up at sky, you are seeing the indirect light of the sun, which is the light emitted by the atoms and molecules of the atmosphere as they oscillate, and so radiate, due to the electric field of the direct light from the sun



power of the radiated indirect light is $P \sim \omega^4 p_0^2$

$$\vec{p} = \alpha \vec{E} \quad \alpha \sim \frac{e^2}{m} \frac{1}{\omega_0^2 - \omega^2 - i\omega\gamma}$$

For molecules in atmosphere (N_2 etc) ω_0 is typically at a freq higher than visible spectrum. Therefore, in visible spectrum $\alpha \sim \frac{e^2}{m\omega_0^2}$ indep of ω .

\Rightarrow power radiated is $P \sim \omega^4$

$P \sim \omega^4$ largest at high freq

Since light from sun is "white light"

it has components of all freqs. Of these freqs, the higher ones are scattered the most & make up the indirect light we see.

Since blue is the largest ω in visible spectrum, the sky is blue!

When we look at sunrise or sunset, we are looking at the direct rays of the sun. Since these rays are least scattered at low $\omega \Rightarrow$ sunset and sunrise are red!

Magnetic Dipole approx - Radiation Zone for $\gg 1$

$$\vec{A}_{MI} = \frac{e^{ikr}}{r} \left(\frac{1}{r} - ik \right) (-\hat{r} \times \vec{m}_\omega)$$

$$\approx ik \hat{r} \times \vec{m}_\omega \frac{e^{ikr}}{r} \quad \text{in RZ}$$

$$\vec{B}_{MI} = \vec{\nabla} \times \vec{A}_{MI} = (\vec{\nabla} e^{ikr}) \times \left(\frac{ik \hat{r} \times \vec{m}_\omega}{r} \right)$$

$$+ e^{ikr} \vec{\nabla} \times \left(\frac{ik \hat{r} \times \vec{m}_\omega}{r} \right)$$

will give terms of $o(\frac{1}{r^2})$
so ignore in RZ approx

$$\vec{B}_{MI} = -k^2 \frac{e^{ikr}}{r} \hat{r} \times (\hat{r} \times \vec{m}_\omega)$$

From Ampere's Law

$$\vec{E}_{MI} = \frac{c}{k} \vec{\nabla} \times \vec{B}_{MI} = -ik (\vec{\nabla} e^{ikr}) \times \left(\frac{\hat{r} \times [\hat{r} \times \vec{m}_\omega]}{r} \right)$$

$$- ik e^{ikr} \vec{\nabla} \times \left(\frac{\hat{r} \times [\hat{r} \times \vec{m}_\omega]}{r} \right)$$

will give terms of $o(\frac{1}{r^2})$
so ignore in RZ approx

$$\vec{E}_{MI} = k^2 \frac{e^{ikr}}{r} \hat{r} \times (\hat{r} \times (\hat{r} \times \vec{m}_\omega))$$

triple product rule

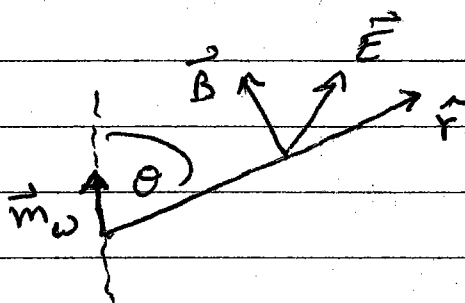
$$= k^2 \frac{e^{ikr}}{r} \left\{ \hat{r} [\hat{r} \cdot (\hat{r} \times \vec{m}_\omega)] - (\hat{r} \times \vec{m}_\omega) [\hat{r} \cdot \hat{r}] \right\}$$

$$\vec{E}_{MI} = -\frac{k^2}{r} e^{ikr} (\hat{r} \times \vec{m}_\omega)$$

If align axes so that $\vec{m}_\omega = m_\omega \hat{z}$ then

$$\vec{E}_{M1} = m_\omega \frac{k^2}{r} e^{i kr} \sin \theta \hat{\phi}$$

$$\vec{B}_{M1} = -m_\omega \frac{k^2}{r} e^{i kr} \sin \theta \hat{\theta}$$



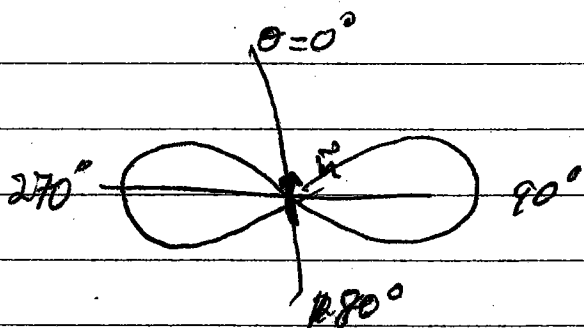
Poynting vector

$$\vec{S}_{M1} = \frac{c}{4\pi} \text{Re} \{ \vec{E}_{M1} \} \times \text{Re} \{ \vec{B}_{M1} \}$$

$$= \frac{c}{4\pi} \frac{k^4 m_\omega^2}{r^2} \cos^2(kr - \omega t) \hat{r}$$

$$\langle \vec{S}_{M1} \rangle = \frac{c}{8\pi} \frac{k^4 m_\omega^2 \sin^2 \theta}{r^2} \hat{r}$$

$$\frac{dP_{M1}}{d\Omega} = \hat{n} \cdot \langle \vec{S}_{M1} \rangle r^2 = \frac{c}{8\pi} k^4 m_\omega^2 \sin^2 \theta \sim \omega^4 \sin^2 \theta$$



rotationally symmetric
about \hat{z} axis

$$P_{M1} = \int d\Omega \frac{dP_{M1}}{d\Omega} = \frac{c k^4}{3} m_\omega^2 = \frac{m_\omega^2 \omega^4}{3c^3}$$

$$\frac{P_{M1}}{P_{E1}} = \frac{m_\omega^2}{P_\omega^2} \sim \left(\frac{v}{c} \right)^2 \quad m_\omega \sim \frac{df}{c} \sim dg \frac{v}{c}$$

$$P_\omega \sim dg$$

Electric Quadrupole radiation - radiation zone approx

$$\vec{A}_{E2} = \frac{e^{i\mathbf{k}\cdot\mathbf{r}}}{r} \left(\frac{1}{r} - ik \right) \left(\frac{-i\omega}{6c} \hat{\mathbf{r}} \cdot \vec{Q}_{\omega} \right)$$

$$= -\frac{e^{i\mathbf{k}\cdot\mathbf{r}}}{r} \frac{k^2}{6} \hat{\mathbf{r}} \cdot \vec{Q}_{\omega} \quad \text{in RZ approx}$$

$$\vec{B}_{E2} = \vec{\nabla} \times \vec{A}_{E2} = -(\vec{\nabla} e^{i\mathbf{k}\cdot\mathbf{r}}) \times \left[\frac{k^2 \hat{\mathbf{r}} \cdot \vec{Q}_{\omega}}{6r} \right]$$

$$- e^{i\mathbf{k}\cdot\mathbf{r}} \vec{\nabla} \times \left[\frac{k^2 \hat{\mathbf{r}} \cdot \vec{Q}_{\omega}}{6r} \right]$$

$\mathcal{O}(1/r^2)$ so ignore in RZ approx

$$\vec{B}_{E2} = -ik^3 \frac{e^{i\mathbf{k}\cdot\mathbf{r}}}{6r} \hat{\mathbf{r}} \times (\hat{\mathbf{r}} \cdot \vec{Q}_{\omega})$$

$$\vec{E}_{E2} = \frac{i}{k} \vec{\nabla} \times \vec{B}_{E2} = k^2 (\vec{\nabla} e^{i\mathbf{k}\cdot\mathbf{r}}) \times \left[\frac{\hat{\mathbf{r}} \times (\hat{\mathbf{r}} \cdot \vec{Q}_{\omega})}{6r} \right]$$

$$+ k^2 e^{i\mathbf{k}\cdot\mathbf{r}} \vec{\nabla} \times \left[\frac{\hat{\mathbf{r}} \times (\hat{\mathbf{r}} \cdot \vec{Q}_{\omega})}{6r} \right]$$

$\mathcal{O}(1/r^2)$ so ignore in RZ approx

$$\vec{E}_{E2} = \frac{ik^3 e^{i\mathbf{k}\cdot\mathbf{r}}}{6r} \hat{\mathbf{r}} \times [\hat{\mathbf{r}} \times (\hat{\mathbf{r}} \cdot \vec{Q}_{\omega})]$$

triple product rule

$$= \frac{ik^3 e^{i\mathbf{k}\cdot\mathbf{r}}}{6r} \left\{ \hat{\mathbf{r}} [\hat{\mathbf{r}} \cdot (\hat{\mathbf{r}} \cdot \vec{Q}_{\omega})] - (\hat{\mathbf{r}} \cdot \vec{Q}_{\omega}) [\hat{\mathbf{r}} \cdot \hat{\mathbf{r}}] \right\}$$

$$\vec{E}_{E2} = \frac{ik^3 e^{i\mathbf{k}\cdot\mathbf{r}}}{6r} \left\{ \hat{\mathbf{r}} (\hat{\mathbf{r}} \cdot \vec{Q}_{\omega} \cdot \hat{\mathbf{r}}) - (\hat{\mathbf{r}} \cdot \vec{Q}_{\omega}) \right\}$$

Poynting vector

$$\vec{S}_{E2} = \frac{-c}{4\pi} \operatorname{Re} \{ \vec{E}_{E2} \} \times \operatorname{Re} \{ \vec{B}_{E2} \}$$

$$= \frac{-k^6}{36r^2} \sin^2(kr - \omega t) \left\{ \hat{r} (\hat{r} \cdot \vec{Q}_\omega \cdot \hat{r}) - (\hat{r} \cdot \vec{Q}_\omega) \right\} \times \left[\hat{r} \times (\hat{r} \cdot \vec{Q}_\omega) \right]$$

$$= \frac{-k^6}{36r^2} \sin^2(kr - \omega t) \left\{ \hat{r} \left[\hat{r} (\hat{r} \cdot \vec{Q}_\omega \cdot \hat{r}) \cdot (\hat{r} \cdot \vec{Q}_\omega) \right] \right. \\ \left. - (\hat{r} \cdot \vec{Q}_\omega) \left[\hat{r} (\hat{r} \cdot \vec{Q}_\omega \cdot \hat{r}) \cdot \hat{r} \right] \right.$$

$$\left. - \hat{r} \left[(\hat{r} \cdot \vec{Q}_\omega) \cdot (\hat{r} \cdot \vec{Q}_\omega) \right] \right. \\ \left. + (\hat{r} \cdot \vec{Q}_\omega) \left[(\hat{r} \cdot \vec{Q}_\omega) \cdot \hat{r} \right] \right\}$$

$$= \frac{-k^6}{36r^2} \sin^2(kr - \omega t) \left\{ (\hat{r} \cdot \vec{Q}_\omega \cdot \hat{r})^2 \hat{r} - (\hat{r} \cdot \vec{Q}_\omega) (\hat{r} \cdot \vec{Q}_\omega \cdot \hat{r}) \right. \\ \left. - (\hat{r} \cdot \vec{Q}_\omega \cdot \vec{Q}_\omega \cdot \hat{r}) \hat{r} + (\hat{r} \cdot \vec{Q}_\omega) (\hat{r} \cdot \vec{Q}_\omega \cdot \hat{r}) \right\}$$

$$\vec{S}_{E2} = \frac{-k^6}{36r^2} \sin^2(kr - \omega t) \left\{ (\hat{r} \cdot \vec{Q}_\omega \cdot \hat{r})^2 - (\hat{r} \cdot \vec{Q}_\omega)^2 \right\} \hat{r}$$

$$\langle \vec{S}_{E2} \rangle = \frac{-k^6}{72r^2} \left\{ (\hat{r} \cdot \vec{Q}_\omega \cdot \hat{r})^2 - (\hat{r} \cdot \vec{Q}_\omega)^2 \right\} \hat{r}$$

$$\frac{dP_{E2}}{d\Omega} = \hat{r} \cdot \langle \vec{S}_{E2} \rangle r^2 = \frac{k^6}{72} \left\{ (\hat{r} \cdot \vec{Q}_\omega)^2 - (\hat{r} \cdot \vec{Q}_\omega \cdot \hat{r})^2 \right\}$$

angular dependence of $\frac{dP_{E2}}{d\Omega}$ depends
on specific form of the tensor \vec{Q}_ω

For example: suppose $Q_{ij} = 0$ except for Q_{zz}
 $\Rightarrow \vec{Q}_W = Q_{zz} \hat{z} \hat{z}$

$$(\hat{r} \cdot \vec{Q}_W \cdot \hat{r})^2 = (Q_{zz} \cos^2 \theta)^2$$

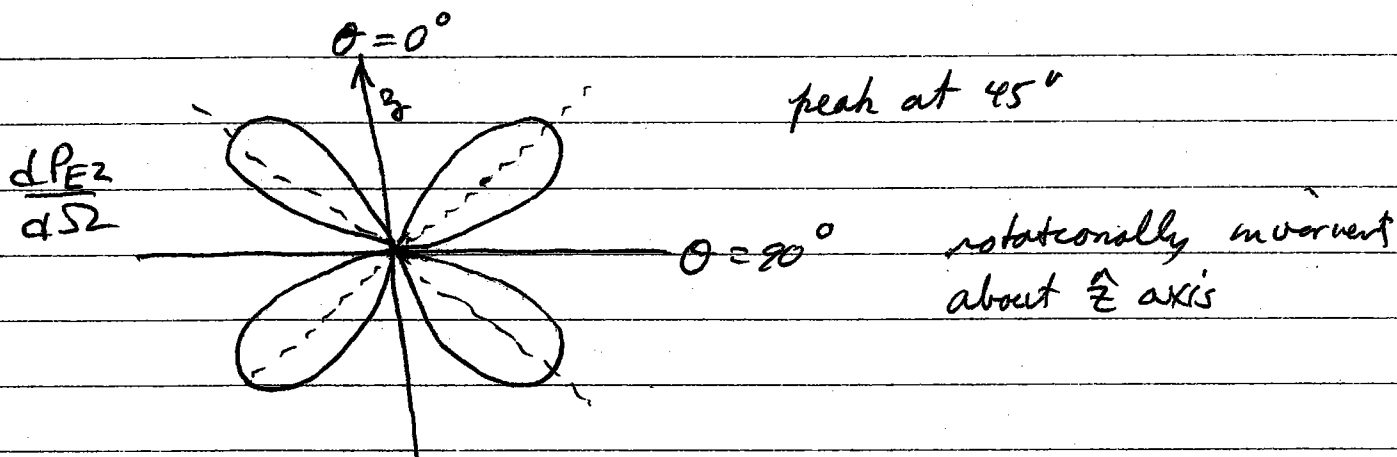
$$(\hat{r} \cdot \vec{Q}_W)^2 = Q_{zz}^2 \cos^2 \theta$$

$$\frac{dP_{E2}}{d\Omega} = \frac{k^6}{72} Q_{zz}^2 [\cos^2 \theta - \cos^4 \theta]$$

$$= \frac{k^6}{72} Q_{zz}^2 \cos^2 \theta \sin^2 \theta$$

$$\cos \theta \sin \theta = \frac{1}{2} \sin 2\theta$$

$$= \frac{k^6}{288} Q_{zz}^2 \sin^2 2\theta$$



$$\frac{P_{E2}}{P_{E1}} \sim \frac{k^6 Q^2}{k^4 p^2} \sim \frac{k^2 (q d^2)^2}{(q d)^2} \sim k^2 d^2 \sim \left(\frac{v}{c}\right)^2$$

$$P_{E2} \sim P_{M1}$$

For more general case, choose axes so that \vec{Q}_0 is diagonal - can always do this since \vec{Q}_0 is symmetric

$$(\hat{r} \cdot \vec{Q}_0 \cdot \hat{r}) = \hat{r} \cdot \begin{pmatrix} Q_{xx} & 0 & 0 \\ 0 & Q_{yy} & 0 \\ 0 & 0 & Q_{zz} \end{pmatrix} \cdot \hat{r}$$

$$= \hat{r} \cdot \begin{pmatrix} Q_{xx} \sin^2 \theta \cos^2 \varphi \\ Q_{yy} \sin^2 \theta \sin^2 \varphi \\ Q_{zz} \cos^2 \theta \end{pmatrix} = Q_{xx} \sin^2 \theta \cos^2 \varphi + Q_{yy} \sin^2 \theta \sin^2 \varphi + Q_{zz} \cos^2 \theta$$

$$(\hat{r} \cdot \vec{Q}_0)^2 = Q_{xx}^2 \sin^2 \theta \cos^2 \varphi + Q_{yy}^2 \sin^2 \theta \sin^2 \varphi + Q_{zz}^2 \cos^2 \theta$$

$$\frac{dP_{E2}}{d\Omega} = \frac{k^6}{72} \left\{ Q_{zz}^2 (\cos^2 \theta - \cos^4 \theta) + Q_{xx}^2 (\sin^2 \theta \cos^2 \varphi - \sin^4 \theta \cos^4 \varphi) + Q_{yy}^2 (\sin^2 \theta \sin^2 \varphi - \sin^4 \theta \sin^4 \varphi) \right\}$$

$$\frac{dP_{E2}}{d\Omega} = \frac{k^6}{72} \left\{ Q_{zz}^2 \cos^2 \theta \sin^2 \theta + Q_{xx}^2 \sin^2 \theta \cos^2 \varphi (1 - \sin^2 \theta \cos^2 \varphi) + Q_{yy}^2 \sin^2 \theta \sin^2 \varphi (1 - \sin^2 \theta \sin^2 \varphi) \right\}$$

no special symmetries - varies with θ and φ