

In CGS, charges are measured in "statcoulombs"
 current measured in "statamperes"
 magnetic field measured in "gauss"

We will use the CGS or Gaussian units

$$1) \vec{\nabla} \cdot \vec{E} = 4\pi\rho$$

$$3) \vec{\nabla} \cdot \vec{B} = 0$$

$$2) \vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$4) \vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

Maxwell's Eqs !!

$$\vec{F} = q \left(\vec{E} + \frac{\vec{v}}{c} \times \vec{B} \right)$$

Lorentz force

Physical content of Maxwell's eqs.

1) Gauss Law for electric field - charge is source of \vec{E} field. Field lines can begin and end at point charges

2) Faraday's Law of induction - time varying magnetic flux produces circulating \vec{E} -field

3) Gauss Law for magnetic fields - no magnetic monopoles. Magnetic field lines are continuous, they either close upon themselves or go off to infinity, they cannot begin nor end at any point.

4) Ampere's Law + Maxwell's correction - electric current is a source for circulating \vec{B} -field; so is a time varying \vec{E} -field. Maxwell's correction is necessary to have charge conservation and to give electromagnetic waves.

Note:

$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{E}) = 0$ as div of curl always vanishes
 then (2) $\Rightarrow -\frac{1}{c} \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{B}) = 0$

So if $\vec{\nabla} \cdot \vec{B} = 0$ at $t=0$, Equ (2) requires that $\vec{\nabla} \cdot \vec{B}$ remains zero for all time

Similarly

$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B}) = 0$
 then (4) $\Rightarrow \frac{4\pi}{c} \vec{\nabla} \cdot \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{E}) = 0$

use charge continuity $\vec{\nabla} \cdot \vec{j} = -\frac{\partial \rho}{\partial t}$ to get

$$-\frac{4\pi}{c} \frac{\partial \rho}{\partial t} + \frac{1}{c} \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{E}) = 0$$

$$\Rightarrow \frac{1}{c} \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{E} - 4\pi\rho) = 0$$

So if $\vec{\nabla} \cdot \vec{E} = 4\pi\rho$ at $t=0$, Equ(4) requires that $\vec{\nabla} \cdot \vec{E} = 4\pi\rho$ for all time

Thus Eqs (1) and (3) can be viewed as "initial conditions". If they are true for any particular moment in time, then Eqs (2) and (4) ensure that they remain true for all time.

Eqs (2) and (3) are also referred to as the homogeneous Maxwell Eqs - they involve only the fields \vec{E} and \vec{B} and not the sources ρ and \vec{j} . Eqs (1) and (4) are referred to as the inhomogeneous Maxwell's eqs - they involve the sources ρ and \vec{j} .

Electromagnetic Potentials & Gauge Invariance

$$\begin{aligned} \vec{\nabla} \cdot \vec{E} &= 4\pi\rho & \vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} &= -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} & \vec{\nabla} \times \vec{B} &= \frac{4\pi\vec{j}}{c} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t} \end{aligned}$$

Consider first statics.

Electrostatics

$$\vec{\nabla} \times \vec{E} = 0 \quad \text{since } \frac{\partial \vec{B}}{\partial t} = 0 \text{ for statics}$$

From vector calculus we know that if the curl of a vector is ~~also~~ everywhere zero, then we can always write that vector field as the gradient of some scalar function ϕ

$$\vec{E} = -\vec{\nabla}\phi \Rightarrow \vec{\nabla} \times \vec{E} = -\vec{\nabla} \times (\vec{\nabla}\phi) = 0$$

ϕ is electrostatic potential

Gauss Law becomes

$$\vec{\nabla} \cdot \vec{E} = -\vec{\nabla} \cdot (\vec{\nabla}\phi) = -\nabla^2\phi = 4\pi\rho$$

$$\boxed{\nabla^2\phi = -4\pi\rho} \quad \text{Poisson's Equation}$$

In regions where $\rho = 0$, we have

$$\nabla^2\phi = 0 \quad \text{Laplace's Equation}$$

In our discussion of Coulomb's Law we saw that the electric field from a distribution of localized charges was

$$\vec{E}(\vec{r}) = \int d^3r' \rho(\vec{r}') \frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3}$$

$$= -\vec{\nabla} \left[\int d^3r' \frac{\rho(\vec{r}')}{|\vec{r}-\vec{r}'|} \right] = -\vec{\nabla}\phi$$

We therefore see that the solution to Poisson's equ for a localized charge distribution ρ (with $\vec{E} = 0$ as $\vec{r} \rightarrow \infty$) is

$$\phi(\vec{r}) = \int d^3r' \frac{\rho(\vec{r}')}{|\vec{r}-\vec{r}'|}$$

We will soon spend a fair amount of time learning new ways to solve $\nabla^2\phi = \rho$, both for arbitrary ρ where we want an approx to the above integral (multipole expansion), and for cases where ϕ or $\vec{\nabla}\phi$ are pre determined on the surfaces of specified regions of space, such as conducting surfaces (boundary value problems).

Magnetostatics

$$\vec{\nabla} \cdot \vec{B} = 0$$

From vector calculus we know that if the divergence of a vector field vanishes everywhere, then it can always be written as the curl of another vector field \vec{A}

$$\vec{B} = \vec{\nabla} \times \vec{A} \Rightarrow \vec{\nabla} \cdot \vec{B} = \vec{\nabla} \cdot (\vec{\nabla} \times \vec{A}) = 0$$

\vec{A} is the magnetic vector potential

This remains true in general - not just in magnetostatics

Ampere's law becomes

$$\vec{\nabla} \times \vec{B} = \frac{4\pi \vec{j}}{c} \quad (\text{in magnetostatics } \frac{\partial \vec{E}}{\partial t} = 0)$$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{A}) = \frac{4\pi \vec{j}}{c}$$

$$\Rightarrow \vec{\nabla} (\vec{\nabla} \cdot \vec{A}) - \nabla^2 \vec{A} = \frac{4\pi \vec{j}}{c}$$

magnetostatic gauge invariance

There are many possible vector potentials \vec{A} that result in the same \vec{B} . If \vec{A} is such that $\vec{\nabla} \times \vec{A} = \vec{B}$, then $\vec{A}' = \vec{A} + \vec{\nabla} \chi$ also gives $\vec{\nabla} \times \vec{A}' = \vec{B}$, since $\vec{\nabla} \times \vec{\nabla} \chi = 0$ for any scalar function $\chi(r)$.

Therefore we can always choose to represent \vec{B} by a vector potential \vec{A} such that $\vec{\nabla} \cdot \vec{A} = 0$.

proof: Suppose we had $\vec{B} = \vec{\nabla} \times \vec{A}$ for some \vec{A} with $\vec{\nabla} \cdot \vec{A} = D(\vec{r}) \neq 0$. Construct an $\vec{A}' = \vec{A} + \vec{\nabla} \chi$ with χ chosen as follows:

$$\vec{\nabla} \cdot \vec{A}' = \vec{\nabla} \cdot \vec{A} + \nabla^2 \chi = 0 \Rightarrow \nabla^2 \chi = -\vec{\nabla} \cdot \vec{A} = D$$

Solve for χ , for example

$$\chi(\vec{r}) = \int d^3 r' \frac{D(\vec{r}')}{4\pi |\vec{r} - \vec{r}'|}$$

We thus have constructed an \vec{A}' such that $\vec{\nabla} \times \vec{A}' = \vec{B}$ and $\vec{\nabla} \cdot \vec{A}' = 0$.

This freedom to choose various \vec{A} 's that give the same \vec{B} is called gauge invariance. Imposing a particular additional constraint on \vec{A} that removes this freedom is called fixing the gauge. The choice $\vec{\nabla} \cdot \vec{A} = 0$ is usually known as the Coulomb gauge (or sometimes the Landau gauge). Going from \vec{A} to $\vec{A}' = \vec{A} + \vec{\nabla} \chi$ is called making a gauge transformation.

"Working in the Coulomb gauge" with $\vec{\nabla} \cdot \vec{A} = 0$, Ampere's Law becomes

$$\boxed{\nabla^2 \vec{A} = -\frac{4\pi}{c} \vec{j}} \quad \text{Poisson's Equ.}$$

For a localized current density

$$\vec{A}(\vec{r}) = \frac{1}{c} \int d^3 r' \frac{\vec{j}(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

Back to dynamics

$$\vec{\nabla} \cdot \vec{B} = 0 \Rightarrow \boxed{\vec{B} = \vec{\nabla} \times \vec{A}} \text{ remains true}$$

But now instead of $\vec{\nabla} \times \vec{E} = 0$ we have

$$\vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0$$

$$\Rightarrow \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial (\vec{\nabla} \times \vec{A})}{\partial t} = 0$$

$$\Rightarrow \vec{\nabla} \times \left(\vec{E} + \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right) = 0$$

\Rightarrow there exists a scalar potential ϕ such that

$$\vec{E} + \frac{1}{c} \frac{\partial \vec{A}}{\partial t} = -\vec{\nabla} \phi \quad \text{or} \quad \boxed{\vec{E} = -\vec{\nabla} \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}}$$

Gauss's law for electric field now becomes

$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho = -\nabla^2 \phi - \frac{1}{c} \frac{\partial (\vec{\nabla} \cdot \vec{A})}{\partial t} = 4\pi\rho$$

$$\boxed{\nabla^2 \phi + \frac{1}{c} \frac{\partial (\vec{\nabla} \cdot \vec{A})}{\partial t} = -4\pi\rho} \quad \text{Gauss law in terms of electromagnetic potentials}$$

Ampere's law becomes

$$\vec{\nabla} \times \vec{B} = \frac{4\pi\vec{j}}{c} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{A}) = \frac{4\pi\vec{j}}{c} + \frac{1}{c} \frac{\partial \vec{\nabla} \phi}{\partial t} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2}$$

$$-\nabla^2 \vec{A} + \vec{\nabla}(\vec{\nabla} \cdot \vec{A}) = \frac{4\pi}{c} \vec{j} - \frac{1}{c} \frac{\partial}{\partial t} \left(\vec{\nabla} \phi + \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right)$$

$$\text{or } -\nabla^2 \vec{A} + \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = \frac{4\pi}{c} \vec{j} - \vec{\nabla} \left(\vec{\nabla} \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

Gauge invariance

As before, we can always construct $\vec{A}' = \vec{A} + \vec{\nabla} \chi$, for any scalar function χ , that gives the same \vec{B} . But since \vec{A} now also enters expression for \vec{E} , we need to make sure that if we change \vec{A} to \vec{A}' , we must make some corresponding change ϕ to ϕ' so that \vec{E} does not change.

$$\left[\begin{array}{l} \vec{A}' = \vec{A} + \vec{\nabla} \chi \\ \phi' = \phi - \frac{1}{c} \frac{\partial \chi}{\partial t} \end{array} \right] \text{ gauge transformation}$$

For any scalar χ , the above \vec{A}' and ϕ' give the same values of \vec{E} and \vec{B} as \vec{A} and ϕ .

Proof:

$$\vec{\nabla} \times \vec{A}' = \vec{\nabla} \times \vec{A} + \vec{\nabla} \times \vec{\nabla} \chi = \vec{\nabla} \times \vec{A} = \vec{B}$$

$$\begin{aligned} \left(-\vec{\nabla} \phi' - \frac{1}{c} \frac{\partial \vec{A}'}{\partial t} \right) &= -\vec{\nabla} \phi + \frac{1}{c} \vec{\nabla} \frac{\partial \chi}{\partial t} - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} - \frac{1}{c} \frac{\partial}{\partial t} \vec{\nabla} \chi \\ &= \left(-\vec{\nabla} \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right) = \vec{E} \end{aligned}$$

As before, we can fix the gauge by imposing some additional constraint on \vec{A} and ϕ . There are two popular choices:

1) Lorentz Gauge

gauge constraint: require $\frac{1}{c} \frac{\partial \phi}{\partial t} + \vec{\nabla} \cdot \vec{A} = 0$

Then Gauss' Law becomes

$$\nabla^2 \phi + \frac{1}{c} \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{A}) = -4\pi \rho$$

$$\Rightarrow \nabla^2 \phi - \frac{1}{c} \frac{\partial}{\partial t} \left(\frac{1}{c} \frac{\partial \phi}{\partial t} \right) = -4\pi \rho$$

$$\boxed{\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = -4\pi \rho}$$

Ampere's Law becomes

$$-\nabla^2 \vec{A} + \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = \frac{4\pi}{c} \vec{j} - \vec{\nabla} \left(\underbrace{\vec{\nabla} \cdot \vec{A}}_0 + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

$$\boxed{\nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\frac{4\pi}{c} \vec{j}}$$

The combination $-\nabla^2 + \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \equiv \square^2$ is the wave equation operator.

In Lorentz gauge, \vec{A} and ϕ satisfy the inhomogeneous wave equations:

$$\boxed{\begin{aligned} \square^2 \vec{A} &= \frac{4\pi}{c} \vec{j} \\ \square^2 \phi &= 4\pi \rho \end{aligned}}$$

when $\vec{j}=0, \rho=0$
electromagnetic waves
are solution!