## Unit 5-6: The Kramers-Kronig Relation

For want of a better place to put this topic, I will put it here!

In our discussion of the frequency dependent atomic polarizability  $\alpha(\omega)$  we saw that the response had to be *causal*, i.e. there is no response before the driving force begins. More mathematically,

$$\mathbf{p}(t) = \int_{-\infty}^{\infty} dt' \,\tilde{\alpha}(t-t') \,\mathbf{E}(t') \tag{5.6.1}$$

and the causality meant that we must have  $\tilde{\alpha}(t) = 0$  for t < 0, so that **p** responds only to **E** at *earlier times*.

In terms of the Fourier transform  $\alpha(\omega)$  we found that causality implied that  $\alpha(\omega)$  could have no poles in the upper half of the complex  $\omega$  plane.

This is a general feature of any causal response function, and it implies that there is a relation between the real and the imaginary parts of the response function. This is the Kramers-Kronig relation.

If  $\alpha(\omega)$  is a causal response function, having no poles in the upper half of the complex  $\omega$  plane (UHP), then for any complex valued  $\bar{\omega}$  in the (UHP), we can write  $\uparrow \operatorname{Im} \omega$ 

where the contour C goes down the real axis, and then closes up with an infinite semicircle in the UHP.

The above result follows because, due to causality,  $\alpha(\omega)$  has no poles in the UHP, and so the only pole of the integrand in the UHP is the pole at  $\omega' = \bar{\omega}$ .

Now the closing semicircular path of the contour in the UHP at infinity should give no contribution to the integral assuming that  $\alpha(\omega)$  decays sufficiently quickly as  $|\omega| \to \infty$ . We thus have,

$$\alpha(\bar{\omega}) = \frac{1}{2\pi i} \int_{-\infty}^{\infty} d\omega' \, \frac{\alpha(\omega')}{\omega' - \bar{\omega}}$$
(5.6.3)

Now consider  $\bar{\omega} = \omega + i\delta$  where  $\omega$  and  $\delta$  are real valued, and  $\delta \to 0$ . Then

$$\alpha(\omega) = \lim_{\delta \to 0} \frac{1}{2\pi i} \int_{-\infty}^{\infty} d\omega' \frac{\alpha(\omega')}{\omega' - \omega - i\delta}$$
(5.6.4)

Now

$$\frac{1}{\omega' - \omega - i\delta} = \mathbb{P}\left(\frac{1}{\omega' - \omega}\right) + i\pi\delta(\omega' - \omega) \qquad \text{where } \mathbb{P} \text{ stands for the principle part of the integral}$$
(5.6.5)

The principle part is defined by,

$$\mathbb{P}\int_{-\infty}^{\infty} d\omega' \frac{\alpha(\omega')}{\omega' - \omega} \equiv \lim_{\epsilon \to 0} \left[ \int_{-\infty}^{\omega - \epsilon} d\omega' \frac{\alpha(\omega')}{\omega' - \omega} + \int_{\omega + \epsilon}^{\infty} d\omega' \frac{\alpha(\omega')}{\omega' - \omega} \right]$$
(5.6.6)

Substituting Eq. (5.6.5) into (5.6.4) then gives

$$\alpha(\omega) = \frac{1}{2\pi i} \mathbb{P} \int_{-\infty}^{\infty} d\omega' \, \frac{\alpha(\omega')}{\omega' - \omega} + \frac{1}{2\pi i} \int_{-\infty}^{\infty} d\omega' \, i\pi \delta(\omega' - \omega) \alpha(\omega) = \frac{1}{2\pi i} \mathbb{P} \int_{-\infty}^{\infty} d\omega' \, \frac{\alpha(\omega')}{\omega' - \omega} + \frac{1}{2} \alpha(\omega) \tag{5.6.7}$$

This then gives

$$\alpha(\omega) = \frac{1}{\pi i} \mathbb{P} \int_{-\infty}^{\infty} d\omega' \, \frac{\alpha(\omega')}{\omega' - \omega}$$
(5.6.8)

Equating the real and the imaginary parts on both sides of the above then gives the Kramers-Kronig relations,

$$\operatorname{Re}\left[\alpha(\omega)\right] = \frac{1}{\pi} \mathbb{P} \int_{-\infty}^{\infty} d\omega' \frac{\operatorname{Im}\left[\alpha(\omega')\right]}{\omega' - \omega} \quad \text{and} \quad \operatorname{Im}\left[\alpha(\omega)\right] = -\frac{1}{\pi} \mathbb{P} \int_{-\infty}^{\infty} d\omega' \frac{\operatorname{Re}\left[\alpha(\omega')\right]}{\omega' - \omega}$$
(5.6.9)

So if a response function is causal, then if one knows the real part  $\operatorname{Re}[\alpha]$ , then one can reconstruct the imaginary part  $\operatorname{Im}[\alpha]$ , and vice-versa.