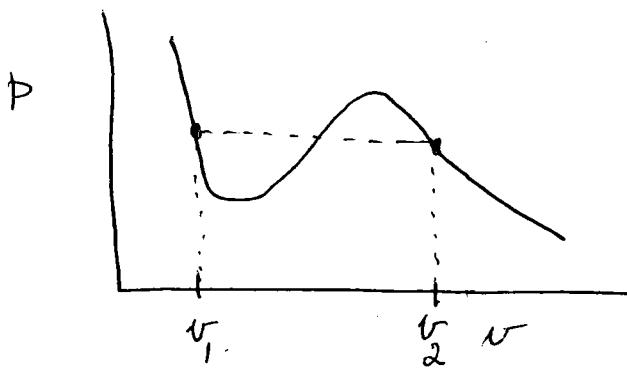


Law of corresponding states: If scale  $P, v, T$  by values at critical point to get  $\bar{P}, \bar{v}, \bar{T}$ , then  
 isotherms in terms of  $\bar{P}, \bar{v}, \bar{T}$  are independent of  $a$  and  $b$   
 - will be the same for any gas.

Now consider  $T < T_c$  we want to determine the coexisting values of  $v_e$  and  $v_g$  that the physical system will have.

Two ways to do this

### ① Maxwell Construction



$$\text{Gibbs free energy } G(T, P, N) = \mu N$$

$$dG = -SdT + Vdp + \mu dN$$

$$= \mu dN + N d\mu$$

$$\Rightarrow d\mu = -\left(\frac{S}{N}\right)dT + Vdp$$

↗ Gibbs-Duhem relation

along isotherm  $dT = 0$

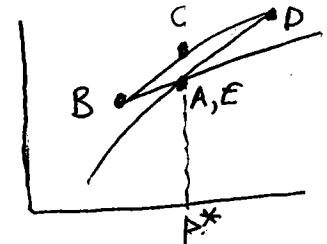
$$\Rightarrow d\mu = Vdp$$

Now at liquid-gas coexistence, the condition for chemical equilib  $\Rightarrow \mu_e = \mu_g$

$$\Rightarrow \int_1^2 d\mu = \mu_2 - \mu_1 = \int_1^2 Vdp = 0 \text{ if}$$

$v_1$  is  $v_e$  and  $v_2$  is  $v_g$  for coexisting  $v_e$  and  $v_g$

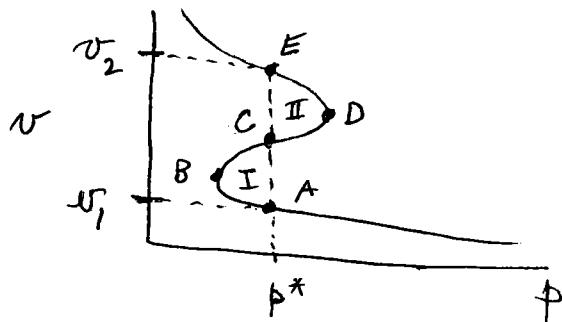
graph of  $\int v dp \rightarrow \mu = \frac{G}{N}$



$\frac{G}{N}$  has cusp at  $P^*$   
 $\frac{dG/N}{dP} = v$  is discontinuous at  $P^*$   
 physical  $G/N$  concave envelope

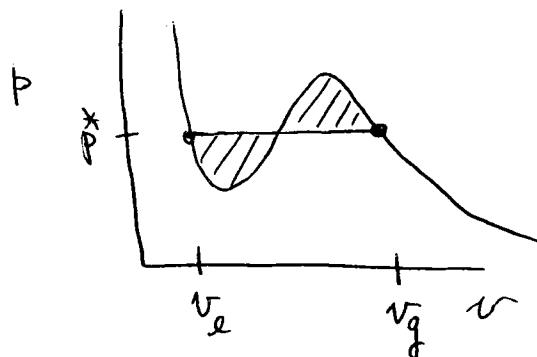
So  $\int_1^2 v dp = 0$  determines  $v_e = v_i \rightarrow v_g = v_z$

To see the geometric meaning of this



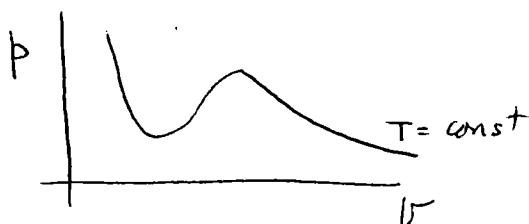
$$\begin{aligned}\int_1^2 v dp &= \int_A^B v dp + \int_B^C v dp \\ &\quad + \int_C^D v dp + \int_D^E v dp \\ &= \text{area } I - \text{area } II\end{aligned}$$

So we get condition of "equal areas"



$v_e$  and  $v_g$  determined by the cord at constant  $P^*$  such that area above cord = area below cord

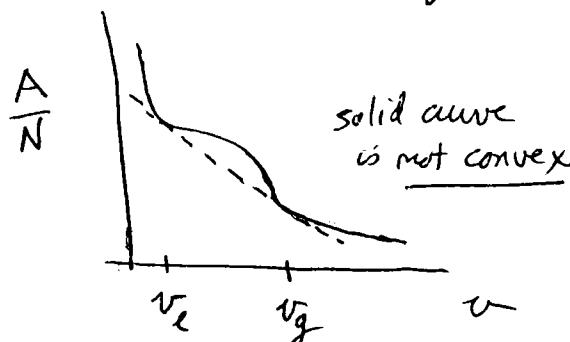
② Consider Helmholtz free energy  $A(T, V)$  fixed  $N$



$$dA = -SdT - pdV$$

$A = - \int pdV$  along isotherm

$$A/N = - \int pdv$$



$$p = -\frac{\partial A}{\partial V} = -\frac{\partial A}{N \partial V}$$

$p$  must be const from  $v_e$  to  $v_g$

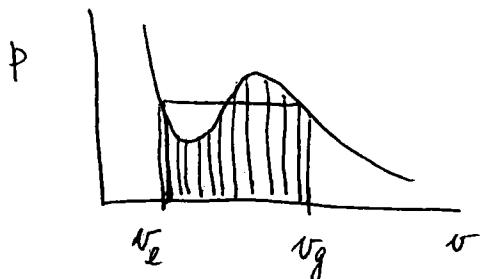
and  $A$  should be convex

$\Rightarrow$  coexistence is given by dashed line

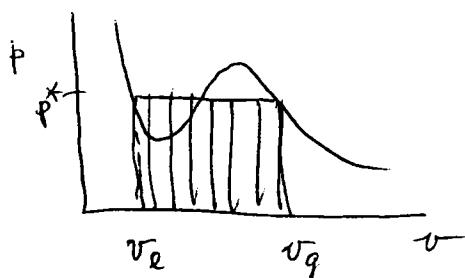
$$P^* = \frac{\left[ \left( \frac{A}{N} \right)_g - \left( \frac{A}{N} \right)_e \right]}{(v_g - v_e)}$$

coexistence pressure is slope of dashed line

$$P^*(v_g - v_e) = -\left[ \left( \frac{A}{N} \right)_g - \left( \frac{A}{N} \right)_e \right] = + \int_e^g p dv$$

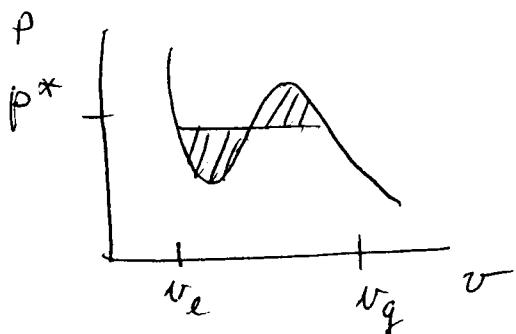


shaded area is  $\int_e^g p dv$



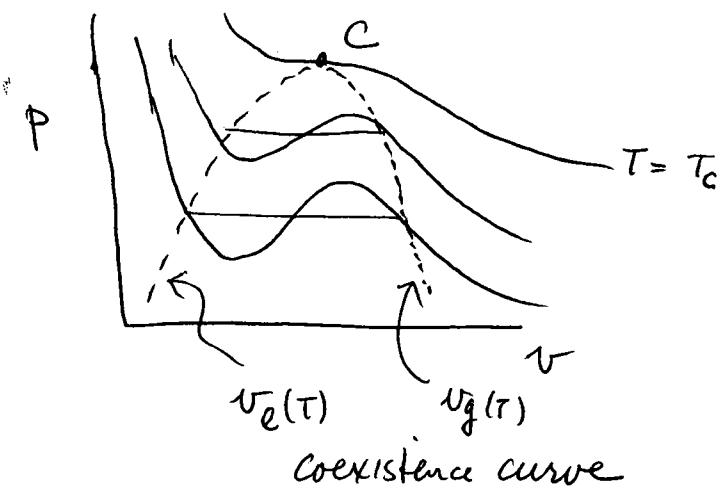
shaded area is  $P^*(v_g - v_e)$

The two areas are equal only if



the shaded area above  $P^*$   
equals shaded area below  $P^*$

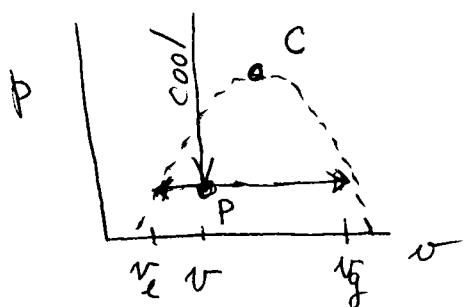
The locus of points  $v_e(T), v_g(T)$  in  $P-v$  plane  
as  $T$  varies is called the coexistence curve.



Region of  $p$ - $v$  plane under coexistence curve is the coexistence region.

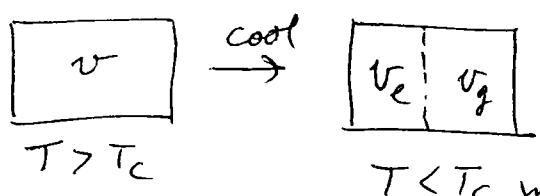
Region in which system can coexist as both liquid and gas

Coexistence region is also called the phase separation region. If sit at a fixed density  $n = 1/v$  that lies within coexistence region for some range of  $T < T_c$ , then cool system into coexistence region to point P,



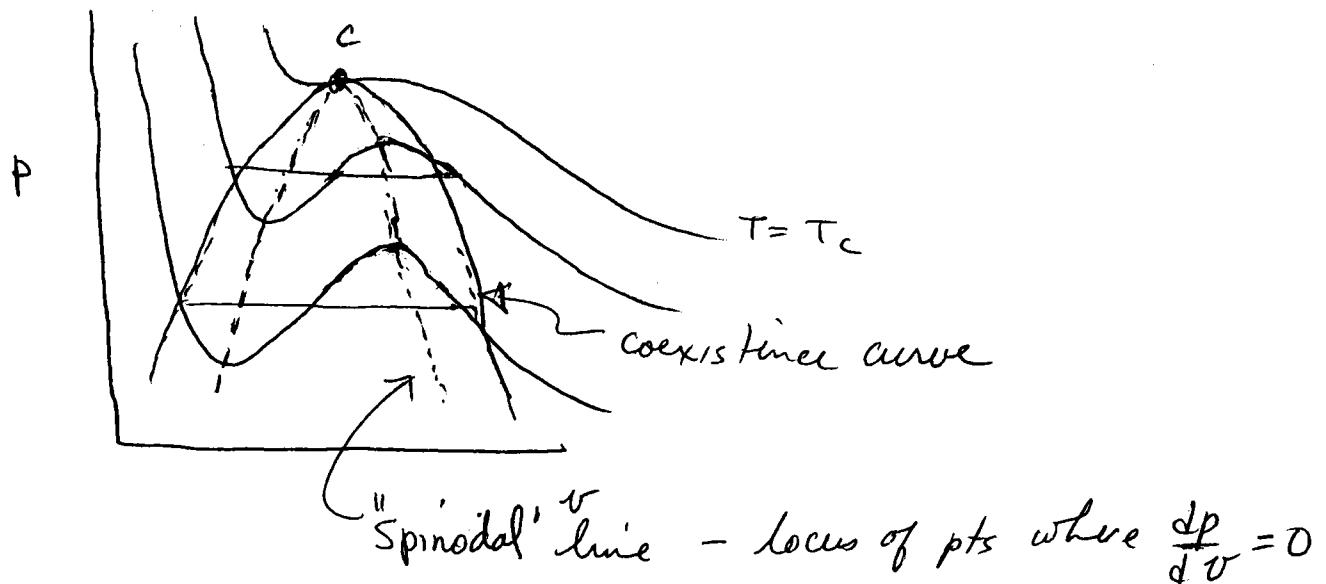
the system can no longer exist at a uniform density  $1/v$ . It will split up in a part with density  $1/v_e$  and a part with density  $1/v_g$  such that average density remains  $1/v$ .

The minimum energy state will be one in which the domain wall separates the liquid + gas to minimum, so one gets



$T < T_c$  in coexistence

the system "phase separates" into a region of liquid and a region of gas



van der waals suggested that states between coexistence curve ad spinodal line are metastable single phase regions

Law of corresponding states  $\Rightarrow$  when expressed in the scaled variables  $\tilde{\phi}$ ,  $\tilde{v}$ ,  $\tilde{T}$ , the coexistence curves of all gases should collapse to a universal curve. Expt results give a good agreement with this prediction, but near C the coexistence curve does not agree too well with the van der Waals prediction

see Fig 3.20 of Plischke + Bergersen