

## 1. MHD eq's including radiation

$$\frac{\partial \vec{p}}{\partial t} = -\nabla \cdot (\rho \vec{v}) \quad (1)$$

$$\frac{\partial \vec{p}}{\partial t} = -\nabla \cdot (\vec{v} \otimes \vec{p} - \tau) - \nabla P + \vec{j} \times \vec{B} + \rho \vec{g} - \overleftrightarrow{\nabla P}_{\text{rad}} \quad (2)$$

$$\frac{\partial e}{\partial t} = -\nabla \cdot (e \vec{v}) - P \nabla \cdot \vec{v} + Q + Q_{\text{rad}} \quad (3)$$

$\vec{p}$  - the momentum density,  $\tau$  - the stress tensor

$P$  - the gass pressure,  $\vec{j}$  - the current density

$\vec{B}$  - the magnetic field vector,  $\vec{g}$  - the acceleration due to gravity

$\overleftrightarrow{\nabla P}_{\text{rad}}$  - the radiation pressure tensor,  $e$  - the internal energy

$Q_{\text{rad}}$  - heating/cooling due to rad.,  $Q$  - other energy exchange

## 2. Energy density of radiation and matter

$$T_{\text{rad}} = 5777 \text{ K}$$

$$E_{\text{rad}} = \frac{4\sigma}{c} T_{\text{rad}}^4 = 0.84 \text{ J m}^{-3} \quad (4)$$

assume B·B:

$$F_{\text{rad}} = \sigma T_{\text{rad}}^4 = 6.3 \times 10^7 \text{ W m}^{-2} \quad (5)$$

## 3. Radiation pressure and force

Isotropic B·B rad.

$$\frac{4\sigma}{3c} T^4 = 0.27 \text{ Pa} \quad @ T_{\text{pl}} = 5700 \text{ K} \quad (6)$$

$$P_{\text{gas}} \sim 10^4 \text{ Pa}$$

$$a_{\text{rad}} \approx \frac{\kappa F}{c} = 3 \times 10^{-3} \text{ ms}^{-2} \quad (7)$$

$\kappa$  - the Rosseland ~~mean~~ opacity,  $F$  - the freq.-integrated rad. flux

$$g_{\theta, \text{surf}} \sim 274 \text{ ms}^{-2}$$

4. Energy exchange between radiation and matter  
energy gaining rate per volume, by a medium from rad. field

$$Q_{\text{rad}} = - \int_0^\infty \nabla \cdot \vec{F}_\nu d\nu = - \nabla \cdot \vec{F} \quad (8)$$

$\vec{F}$  - the total radiative flux

5. Explicit expression of the

$$\vec{F}_\nu = \oint \hat{n} I_\nu d\Omega \quad (9)$$

Source function:  $S_\nu = j_\nu / (\kappa_\nu \rho)$

$$\begin{aligned} \nabla \cdot \vec{F} &= \int_0^\infty \oint (j_\nu - \kappa_\nu \rho I_\nu) d\Omega d\nu \\ &= \int_0^\infty \oint \kappa_\nu \rho (S_\nu - I_\nu) d\Omega d\nu \end{aligned} \quad (10)$$

If both  $S_\nu$  and  $\kappa_\nu$  do not depend on direction:

$$\nabla \cdot \vec{F} = \int_0^\infty 4\pi \kappa_\nu \rho (S_\nu - J_\nu) d\nu \quad (11)$$

the angle-averaged intensity:

$$J_\nu = \frac{1}{4\pi} \oint I_\nu d\Omega \quad (12)$$

## 6. Light travel time and hydrodynamical timescales

$$v_{\text{bulk, solar atm}} \quad l \sim > 300 \text{ km s}^{-1}$$

$$v_{\text{Alfvén}} \quad \text{up to } 2.2 \times 10^3 \text{ km s}^{-1}$$

$$t_{\text{hydro}} \quad \begin{array}{ll} \sim 5 \text{ min} & \text{in photosphere} \\ \sim 1 \text{ min} & \text{in chromosphere} \end{array}$$

few min in corona

$$t_{\text{light cross}} \quad \sim 1 \text{ sec}$$

$t$ -dependent transfer eq.

$$\frac{1}{c} \frac{\partial I_\nu}{\partial t} + \hat{n} \cdot \nabla I_\nu = j_\nu - \alpha_\nu I_\nu \quad (13)$$

Simplifies to  $\rightarrow$

$$\hat{n} \cdot \nabla I_\nu = j_\nu - \alpha_\nu I_\nu \quad (14)$$

$j_\nu$  - emissivity ,  $\alpha_\nu = \kappa_\nu \rho$  extinction coeff. per volume

## 7. Diffusion approximation

At very large optical depth

$$J_\nu \approx S_\nu + \frac{1}{3} \frac{d^2 S_\nu}{dT_\nu^2} \quad (15)$$

$$\vec{F}_\nu \approx -\frac{4}{3\kappa_\nu P} \nabla S_\nu \quad (16)$$

for all freq.  $\Rightarrow$

assume  $S_\nu = B_\nu$

$$\vec{F} \approx -\frac{16}{3} \frac{\sigma T^3}{\kappa_R P} \nabla T \quad (17)$$

Rosseland opacity

$$\frac{1}{\kappa_R} = \left( \int_0^\infty \frac{1}{\kappa_\nu} \frac{dB_\nu}{dT} d\nu \right) / \left( \int_0^\infty \frac{dB_\nu}{dT} d\nu \right) \quad (18)$$