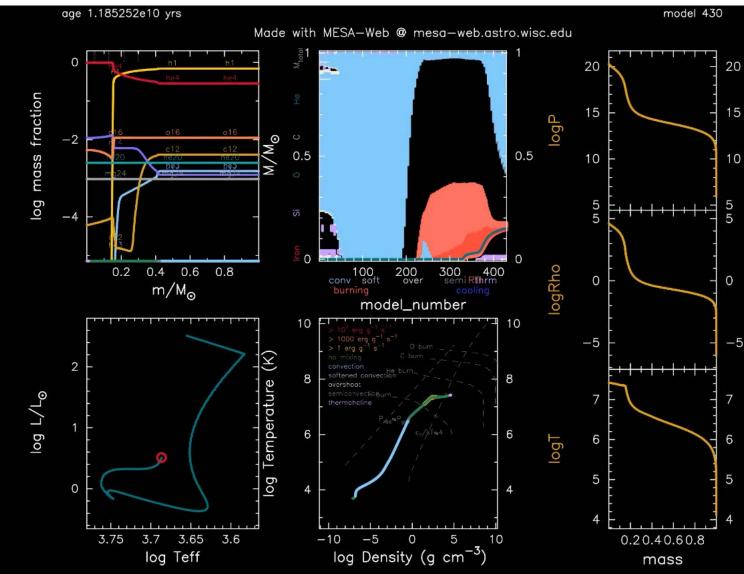
Today in Astronomy 241: stellar evolution on, and slightly off, the main sequence

- Changes in the stellar core composition
- Buildup of an isothermal core
- The maximum in weight which can be supported by an isothermal core
- Reading: C&O chapter 13, pp. 446-457

MESA model of the Sun, through age 12 Gyr.



Stellar evolution on the main sequence

- As hydrogen burns in the core, the mean molecular weight slowly increases. At a given temperature the ideal gas law says this would result in a lower pressure. Therefore
 - the core slowly contracts and heats up.
 - the radius and effective temperature of the star's **envelope** slowly increase, in response to the new internal temperature and density distribution.
 - the luminosity slowly increases, in response to the increase in radius and effective temperature.
- This all takes place much faster for higher-mass stars.

Evolution on the main sequence (continued)

• Most of these trends are evident in the results of the simplified stellar models you have recently worked out. Consider the result from classes on 26 and 28 February:

$$L = \frac{\pi^3}{64^2} \left(\frac{7}{4}\right)^{6.5} \left(\frac{3}{4}\right)^{1.5} \frac{\sigma}{5\kappa_0} \left(\frac{G\mu m_H}{k}\right)^{15/2} M^{11/2} R^{-1/2}$$
$$T_0 = \left(\left(\frac{\pi}{2}\right)^2 \left(\frac{7}{2}\right)^{6.5} \left(\frac{3}{2}\right)^{1.5} \frac{1}{2}\right)^{1/4} \left(\frac{G\mu m_H}{2}\right)^{15/8} M^{11/2} R^{-1/2}$$

$$T_e = \left(\left(\frac{\pi}{64}\right)^2 \left(\frac{7}{4}\right)^{6.5} \left(\frac{3}{4}\right)^{1.5} \frac{1}{20\kappa_0} \right)^{1/4} \left(\frac{G\mu m_H}{k}\right)^{15/8} M^{11/8} R^{-5/8}$$

$$t \propto \frac{M}{L} \propto M^{-9/2}$$

Shell hydrogen burning and the subgiant phase

- Eventually hydrogen is exhausted in the very center, and the temperature is insufficient to ignite helium burning, but is high enough just outside the center for a **shell** of hydrogen fusion to keep the star hydrostatic. Thus
 - **T** is nearly constant in the core (isothermal helium core, or *ic*), which keeps increasing in mass owing to hydrogen depletion;
 - luminosity increases and the envelope expands further;
 - effective temperature decreases further;
 - This is called the **subgiant phase**: the star moves off the main sequence, upwards and to the right in the HR diagram.
- Pressure exerted by core on envelope, C&O equation 13.7:

$$P_{ic} = \frac{3}{4\pi R_{ic}^3} \left(\frac{M_{ic} kT_{ic}}{\mu_{ic} m_{\rm H}} - \frac{1}{5} \frac{GM_{ic}^2}{R_{ic}} \right)$$

Critical point for the isothermal core

• This pressure has a maximum, given by C&O equation 13.9:

$$P_{ic, \max} = \frac{375}{64\pi G^3 M_{ic}^2} \left(\frac{kT_{ic}}{\mu_{ic}m_{\rm H}}\right)^4$$

and thus there is a maximum to the mass of the core for hydrostatic stars (Schoenberg and Chandrasekhar 1942):

$$\frac{M_{ic}}{M} \approx 0.54 \left(\frac{\mu_{env}}{\mu_{ic}}\right)^2 = 0.54 \left(\frac{0.62}{1.34}\right)^2 = 0.12 \text{ for the Sun}$$

• The subgiant phase ends when this mass exceeded, and core pressure can no longer support the envelope's weight; big changes begin to happen much faster thereafter.

Electron degeneracy in the isothermal core

- Degeneracy pressure from electrons can increase the support of the core against its weight and that of the envelope; this leads to a modest increase (to about 0.13) over the Schoenberg-Chandra limit.
- Degenerate electron gases are polytropes, of order n = 1.5 for nonrelativistic electrons and n = 3 for ultrarelativistic ones.

Today's in-class problems

- 1. Run MESA via MESA-Web.
 - a. Go to <u>MESA-Web</u> and run a model of the Sun, with the default parameters except for a maximum age of 12 Gyr (under Custom Stopping Condition), and a profile output frequency of every 5 models (under General).

While waiting:

- 2. Suppose that the pressure P in a stellar interior consists of a fraction β of radiation pressure P_{rad} and a remaining fraction 1β of gas pressure P_{gas} . Suppose further that β is uniform in the star's interior. Show that the resulting equation of state is a polytrope of order n = 3, i.e. $\gamma = C_P / C_V = 4/3$.
- 3. C&O 13.3.
- b. In the description of <u>MESA-Web's output</u>, familiarize yourself with profile output. Find where MESA lists the important stellar parameters – age, total mass, outer radius, radius, density, pressure, temperature, mean molecular weight, adiabatic gradient – can be found. You will be using MESA regularly from now on.

Today's in-class problems (continued)

Hints for the previous set of in-class problems:

- 1. I presume everybody has a smoothly running Lane-Emden solver now.
- 2. I get $\rho_C = 8.44 \text{ gm cm}^{-3}$, $P_C = 8.66 \times 10^{15} \text{ dyne cm}^{-2}$, and $T_C = 7.72 \times 10^6 \text{ K}$.
- 3. (C&O 10.19) Put the n = 5 solution for D_n into the expression for total mass of the star:

$$\begin{split} M &= -4\pi\alpha^{3}\rho_{C}\xi_{1}^{2}D_{n}'(\xi_{1}) = -4\pi\alpha^{3}\rho_{C}\xi_{1}^{2}\left[\frac{d}{d\xi}\left(1+\frac{\xi^{2}}{3}\right)^{-1/2}\right]_{\xi_{1}} \\ &= -4\pi\alpha^{3}\rho_{C}\xi_{1}^{2}\left(-\frac{1}{2}\right)\left(1+\frac{\xi_{1}^{2}}{3}\right)^{-3/2}\left(\frac{2\xi_{1}}{3}\right) = \frac{4\pi}{3}\alpha^{3}\rho_{C}\left(\frac{1}{\xi_{1}^{2}}+\frac{1}{3}\right)^{-3/2} \xrightarrow{\xi_{1}\to\infty} 3^{3/2}\frac{4\pi}{3}\alpha^{3}\rho_{C} \quad . \end{split}$$