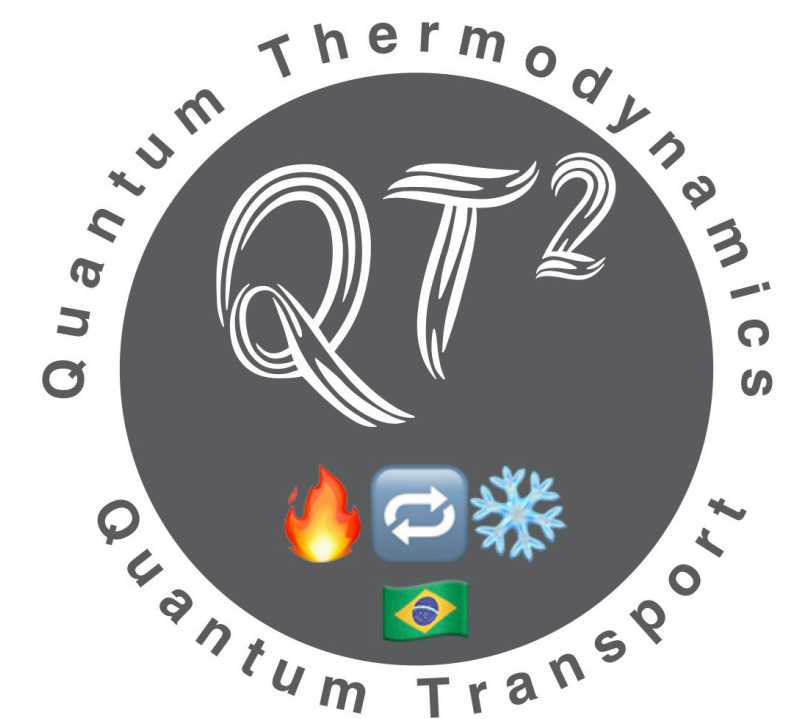


# Information & fluctuations in continuously measured systems

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September 2022 - Quantum Intelligence - Birr, Ireland



[www.fmt.if.usp.br/~gtlandi](http://www.fmt.if.usp.br/~gtlandi)

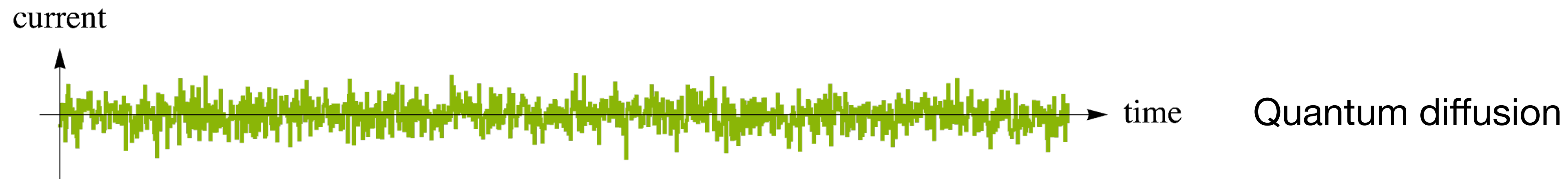
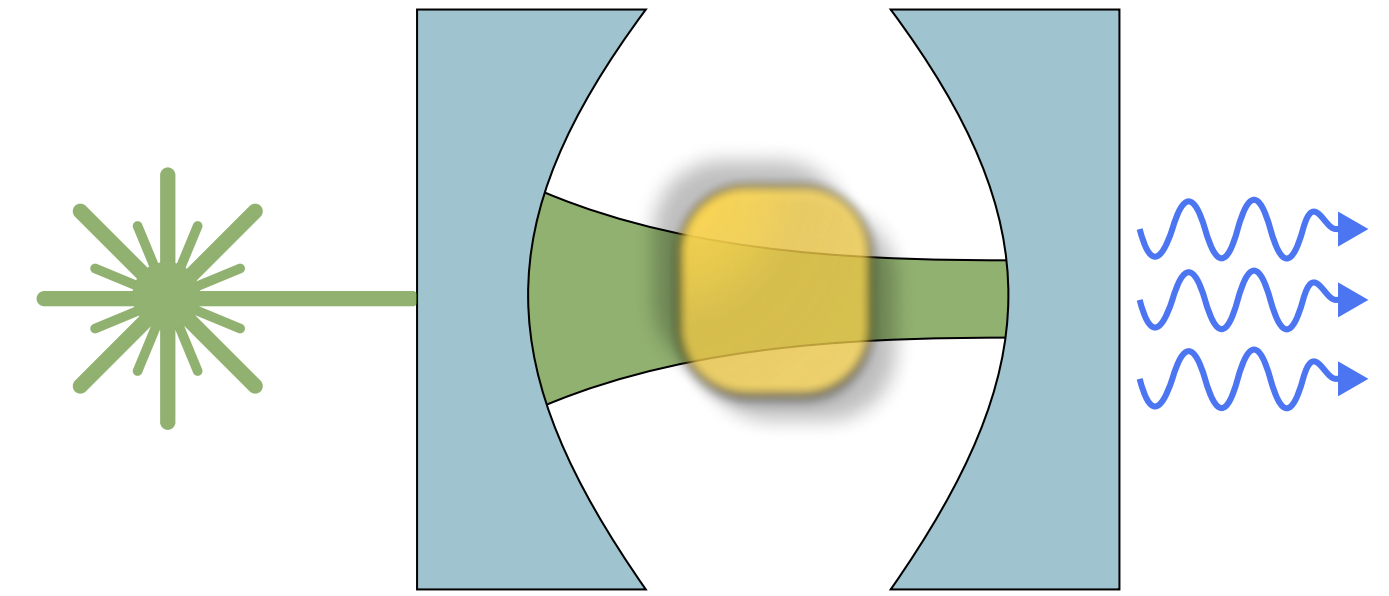
# Overview

- Open Quantum Systems + continuous weak measurement.

➔ ***Classical (stochastic) currents (time-series)***

- **Motivation:**

1. Current fluctuations & two-time correlations.
2. What information the current conveys about the system.





# Current fluctuations & two-time correlations

M. Kewming, M. Mitchison & GTL,  
“**Diverging current fluctuations in critical Kerr resonators**”,  
*arXiv 2205.02622. Accepted in PRA.*

GTL, M. Kewming, M. Mitchison, P. Potts,  
“**Current fluctuations in open quantum systems: Bridging the gap between quantum continuous measurement and full counting statistics.**”  
*In preparation. Tutorial for PRX Quantum.*



Patrick Potts



Michael Kewming



Mark Mitchison



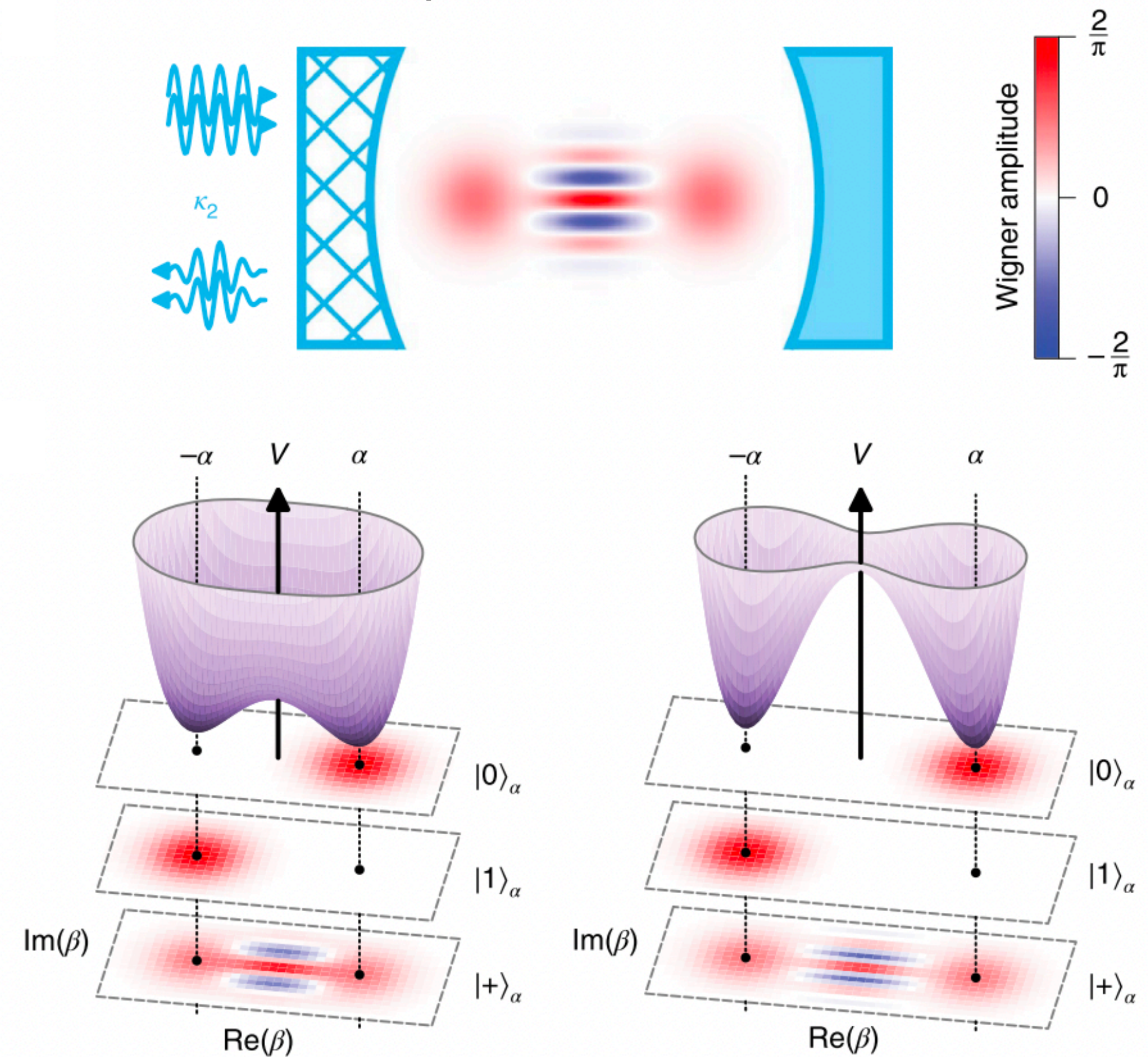
# Example: Kerr model

- Quantum Master equation:

$$\frac{d\rho}{dt} = \mathcal{L}(\rho) = -i[H(t), \rho] + \kappa \left[ a\rho a^\dagger - \frac{1}{2}\{a^\dagger a, \rho\} \right]$$

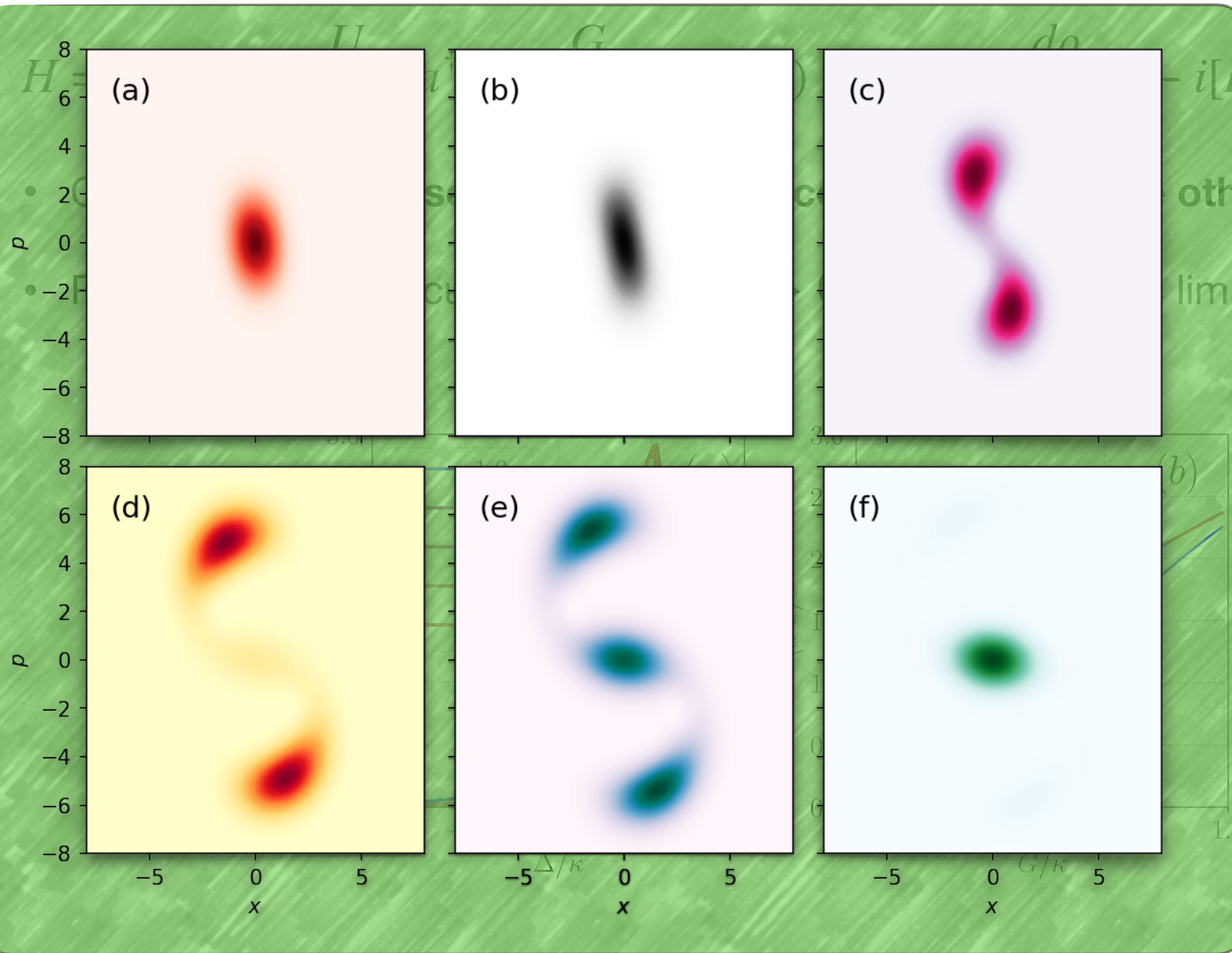
$$H = -\Delta a^\dagger a + \frac{U}{2} a^\dagger a^\dagger a a + \frac{G}{2} (a^{\dagger 2} + a^2)$$

Cat qubits:  
useful for quantum error correction



Lescanne, *et. al.*, Nature, **16**, 509-513 (2020)

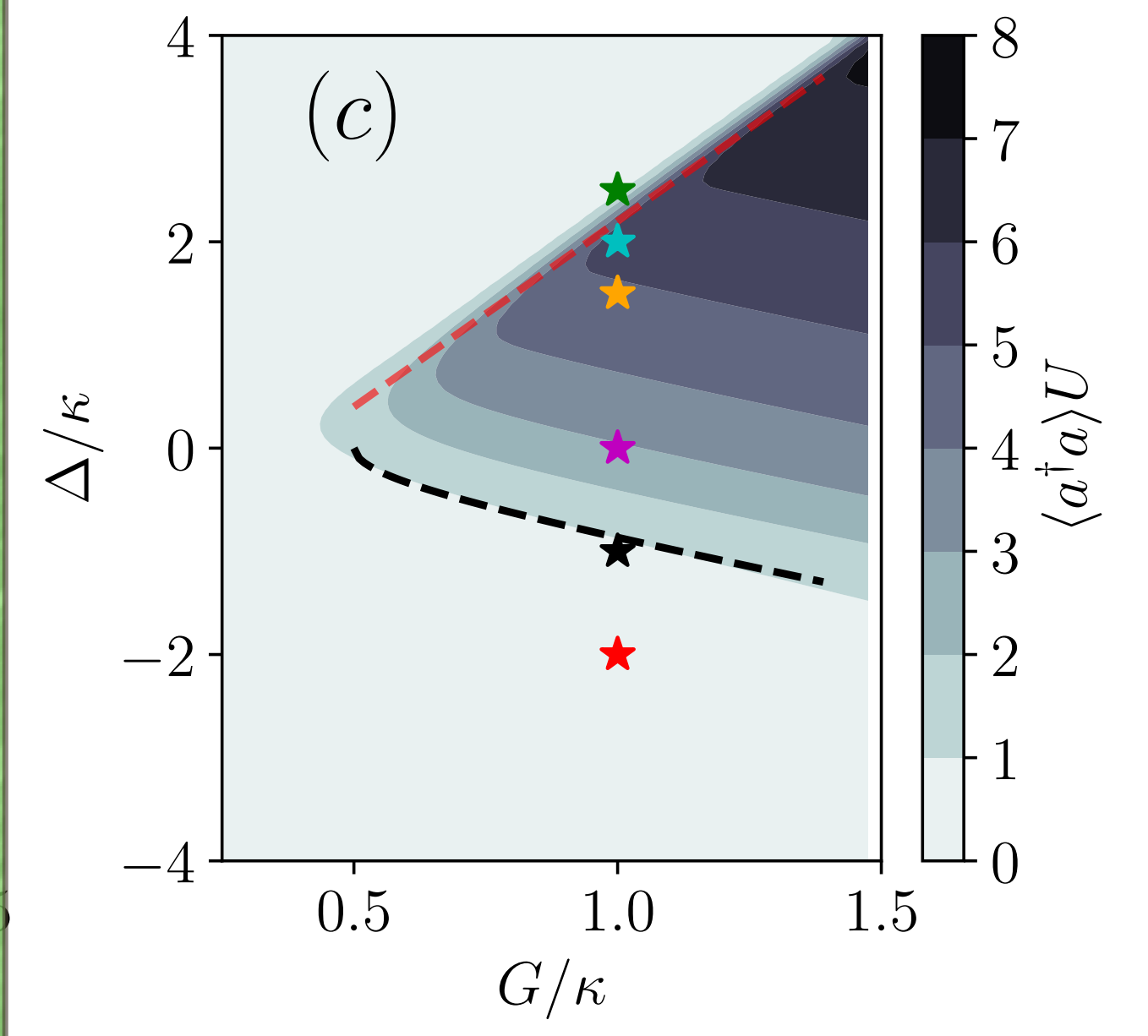




$$i[H, \rho] + \kappa D[a]$$

other discontinuous:

limit")



# Photo-detection current

- Quantum Jump unravelling:

With prob.  $p = dt \operatorname{tr}(a\rho a^\dagger)$  the system jumps:  $\rho \rightarrow \frac{a\rho a^\dagger}{\operatorname{tr}(a\rho a^\dagger)}$

With prob.  $1 - p$  it does not jump:  $\rho \rightarrow \frac{M_0\rho M_0^\dagger}{\operatorname{tr}(M_0\rho M_0^\dagger)}$

$$M_0 = 1 - dt\left(iH + \frac{\kappa}{2}a^\dagger a\right)$$

- Jump:  $dN = 1$

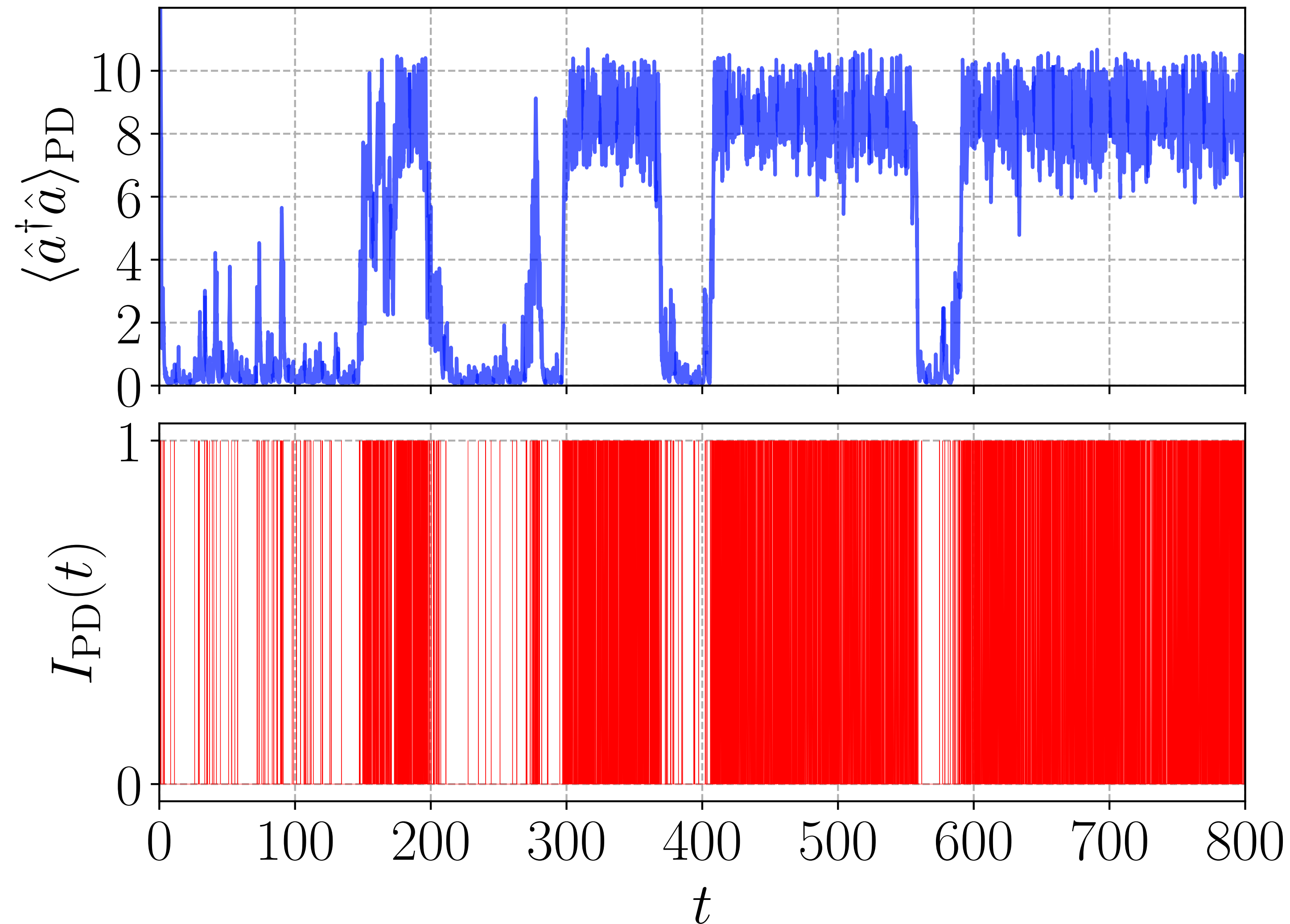
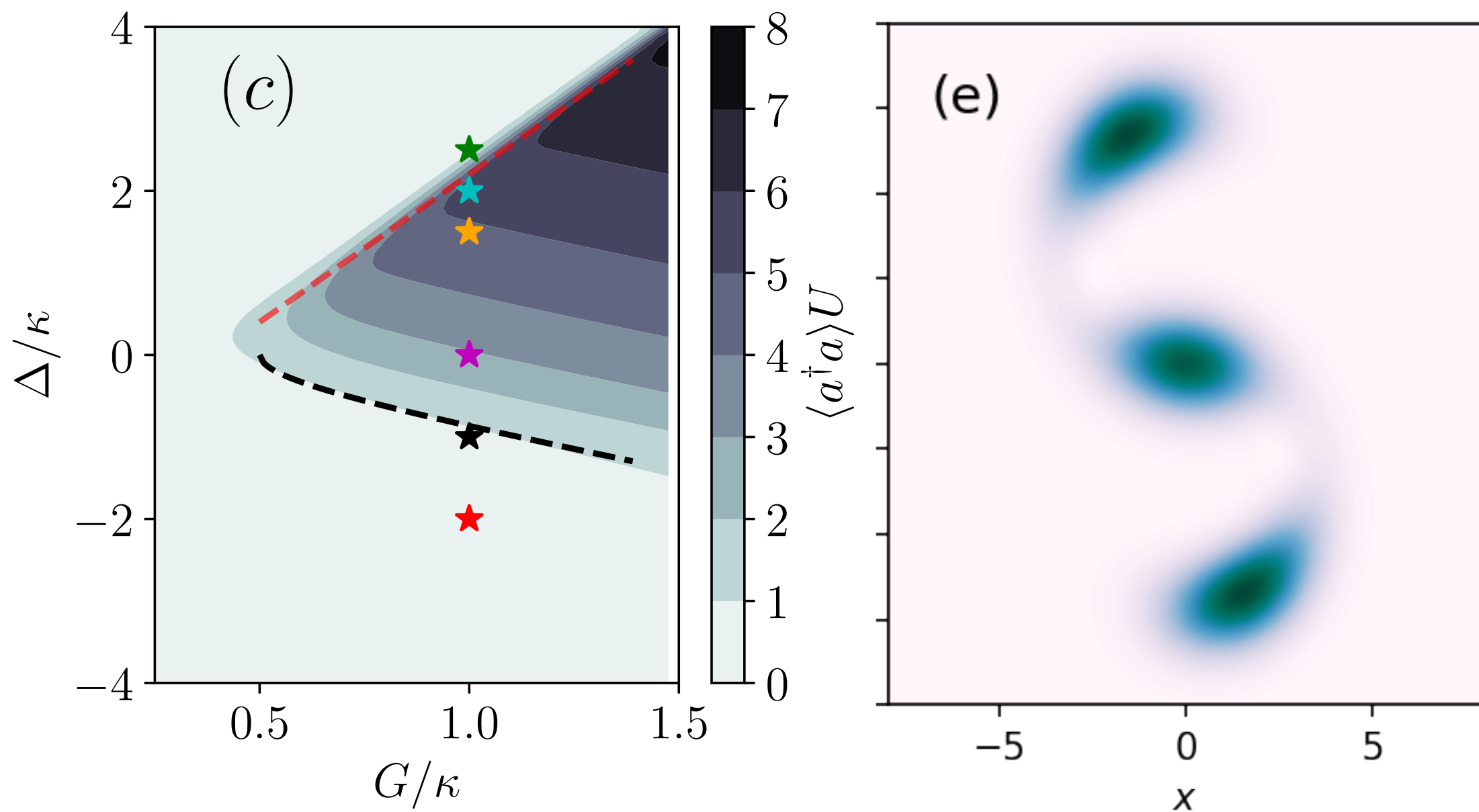
- (stochastic) current:  $I(t) = \frac{dN}{dt}$

- Integrated current (net charge):  $N(t) = \int_0^t dt' I(t')$



# Photo-detection current

- @ discontinuous transition: on/off (telegraph) behavior of the current.
- Photo-detection cannot resolve upper vs. lower blobs.



# Full Counting Statistics

- Average photon current:  $J = \text{tr}(\mathcal{J}\rho) = \kappa\langle a^\dagger a \rangle$  where

$$\mathcal{J}\rho = \kappa\rho a^\dagger$$

- Two-time correlation function:

$$F(\tau) := \text{Cov}(I(t)I(t+\tau)) = K\delta(\tau) + \text{tr}\{\mathcal{J}e^{\mathcal{L}|\tau|}\mathcal{J}\rho\} - J^2$$

$K = J$  (in this case) is the **dynamical activity**: *jumps/second*.

- Power spectrum:

$$S(\omega) = \int_{-\infty}^{\infty} e^{-i\omega\tau} F(\tau) d\tau$$

- Zero-frequency component of the power spectrum := “noise”:

$$D = S(0) = \lim_{t \rightarrow \infty} \frac{d}{dt} \text{Var}(N(t))$$

$D$  is the quantity in TURs  
and KURs

$$\text{TUR: } \frac{D}{J^2} \geq \frac{2}{\dot{\sigma}}$$

$$\text{KUR: } \frac{D}{J^2} \geq \frac{1}{K}$$

Both can be violated in the  
quantum regime.



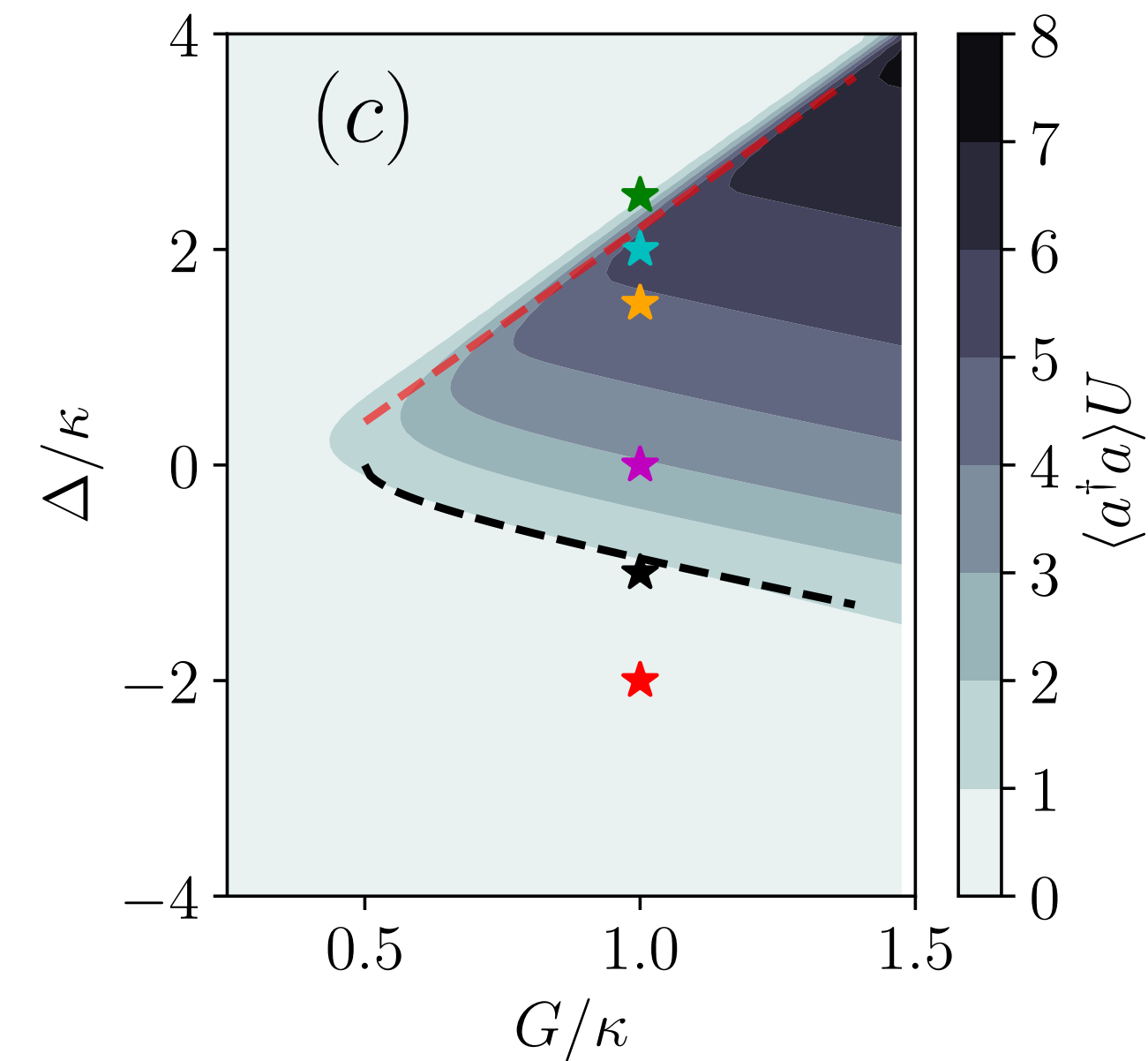
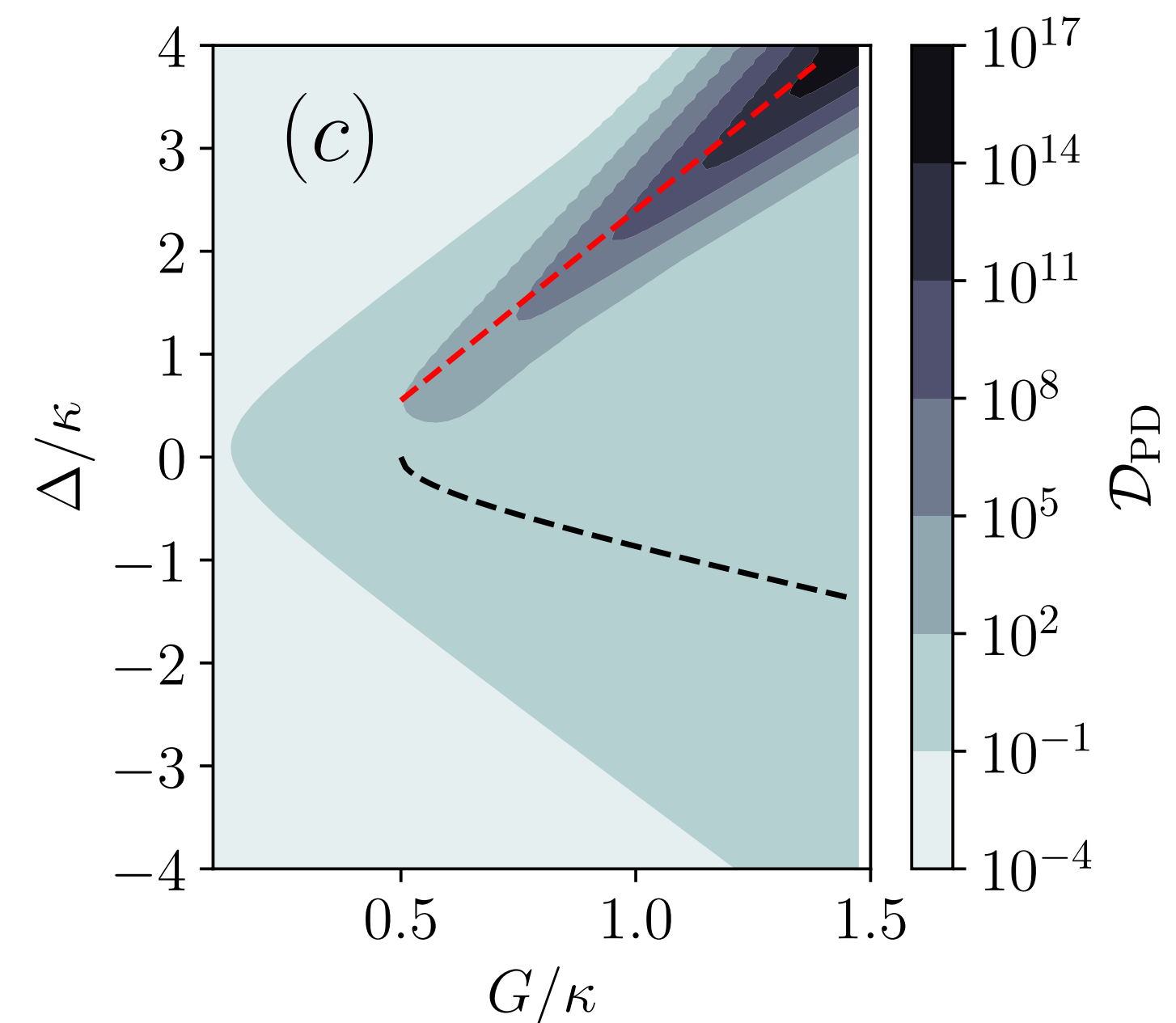
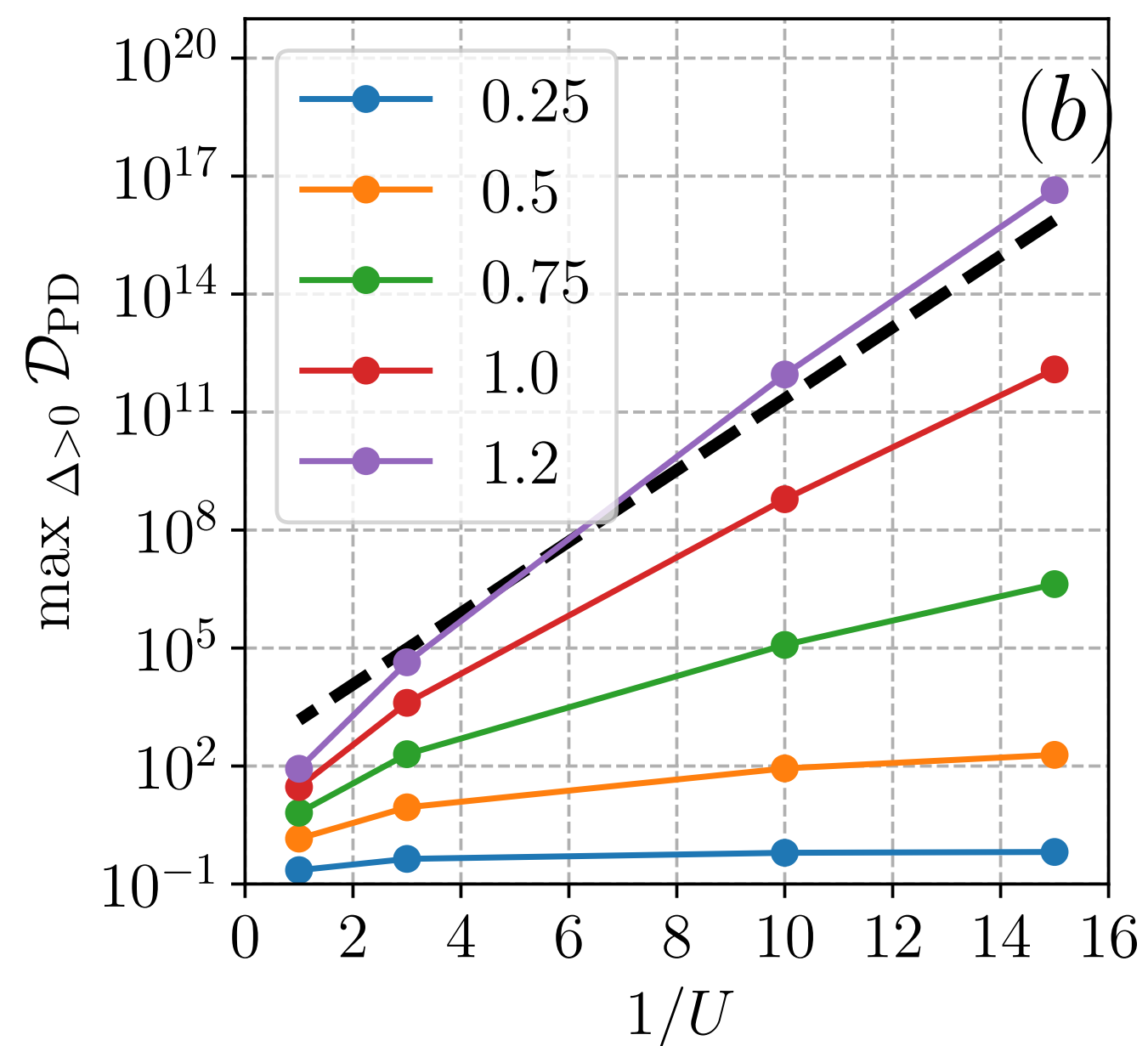
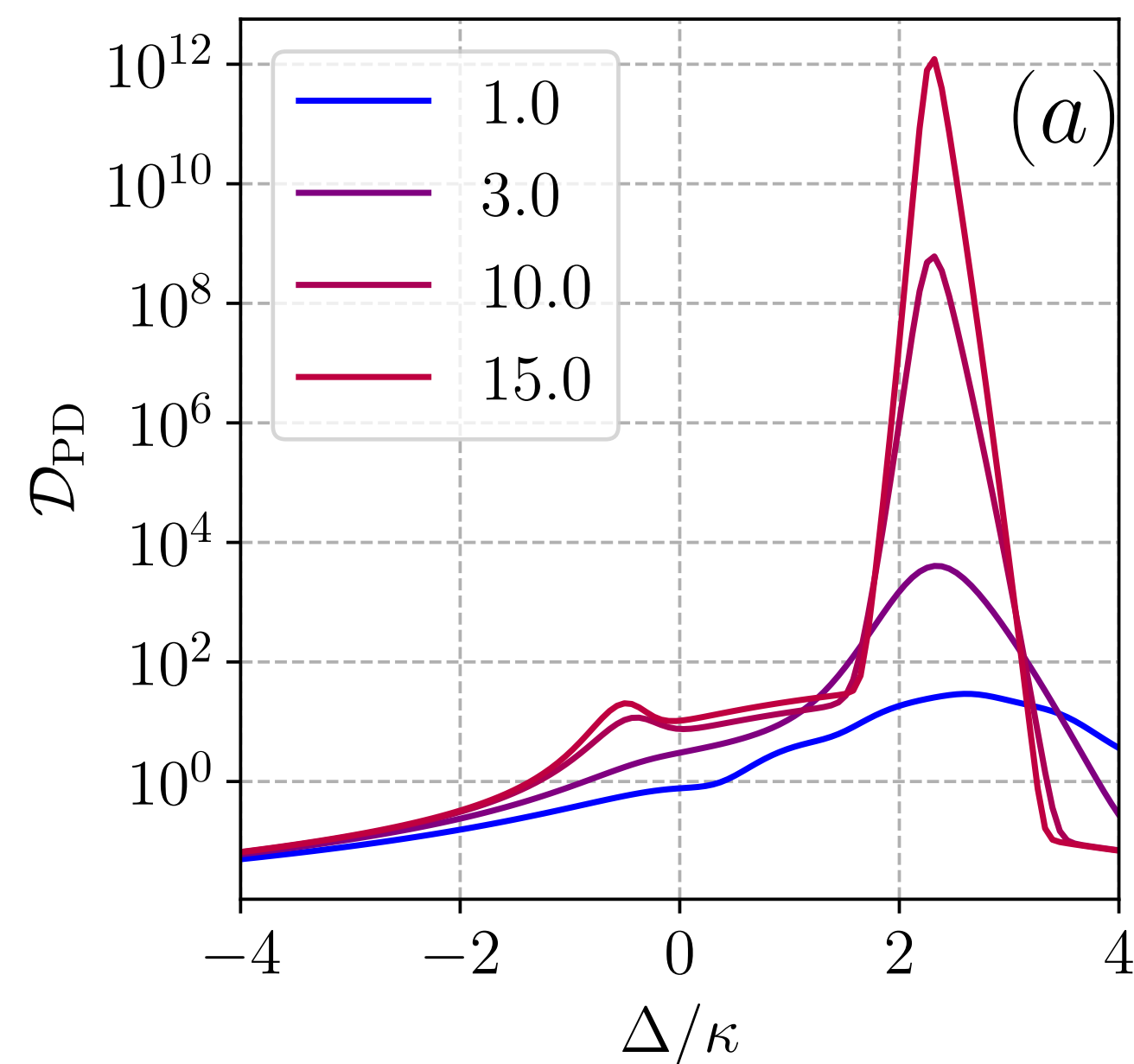
# Divergence of the diffusion coefficient

- “Thermodynamic limit:”  $U \rightarrow 0$
- In the continuous transition ( $\Delta < 0$ )

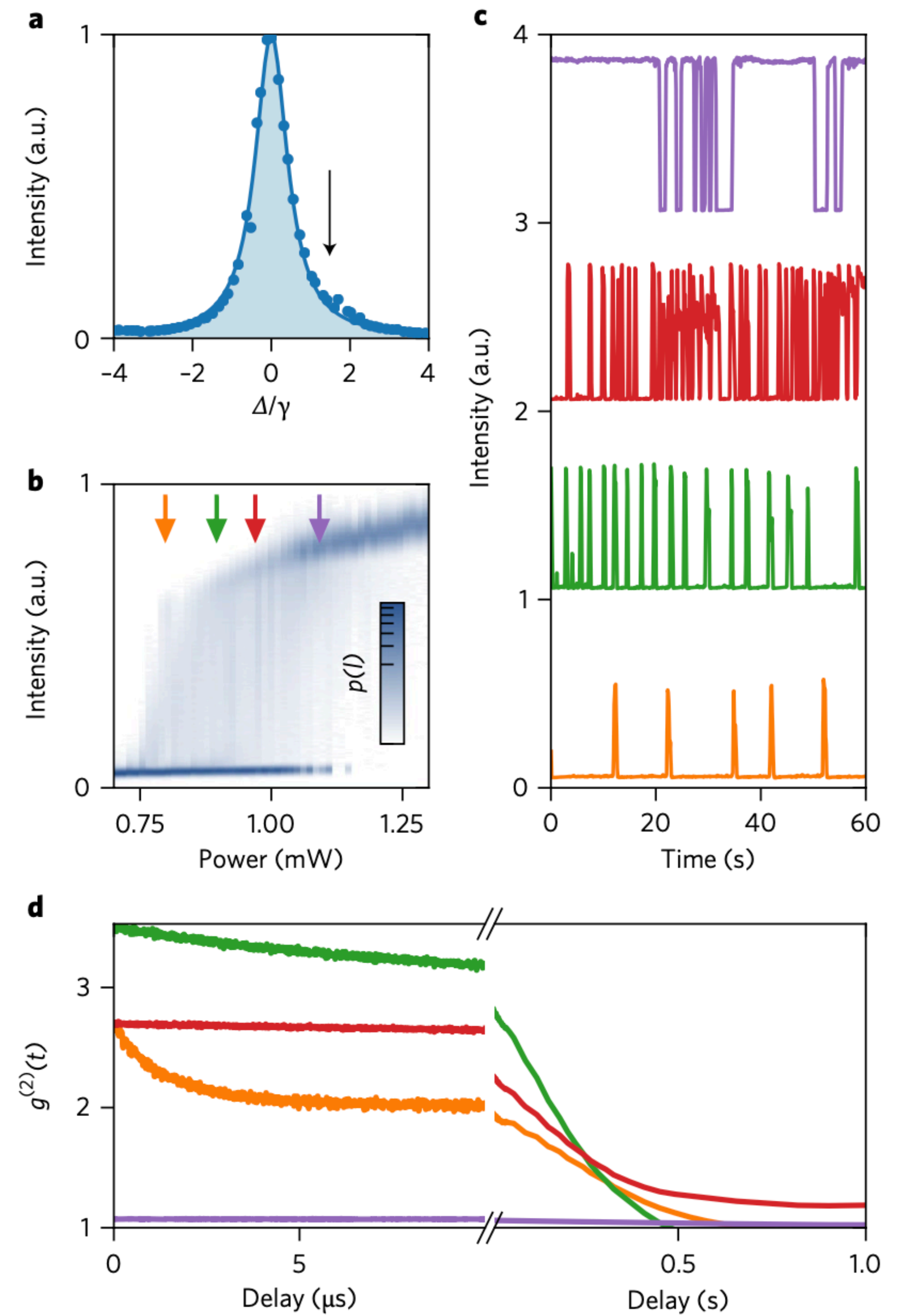
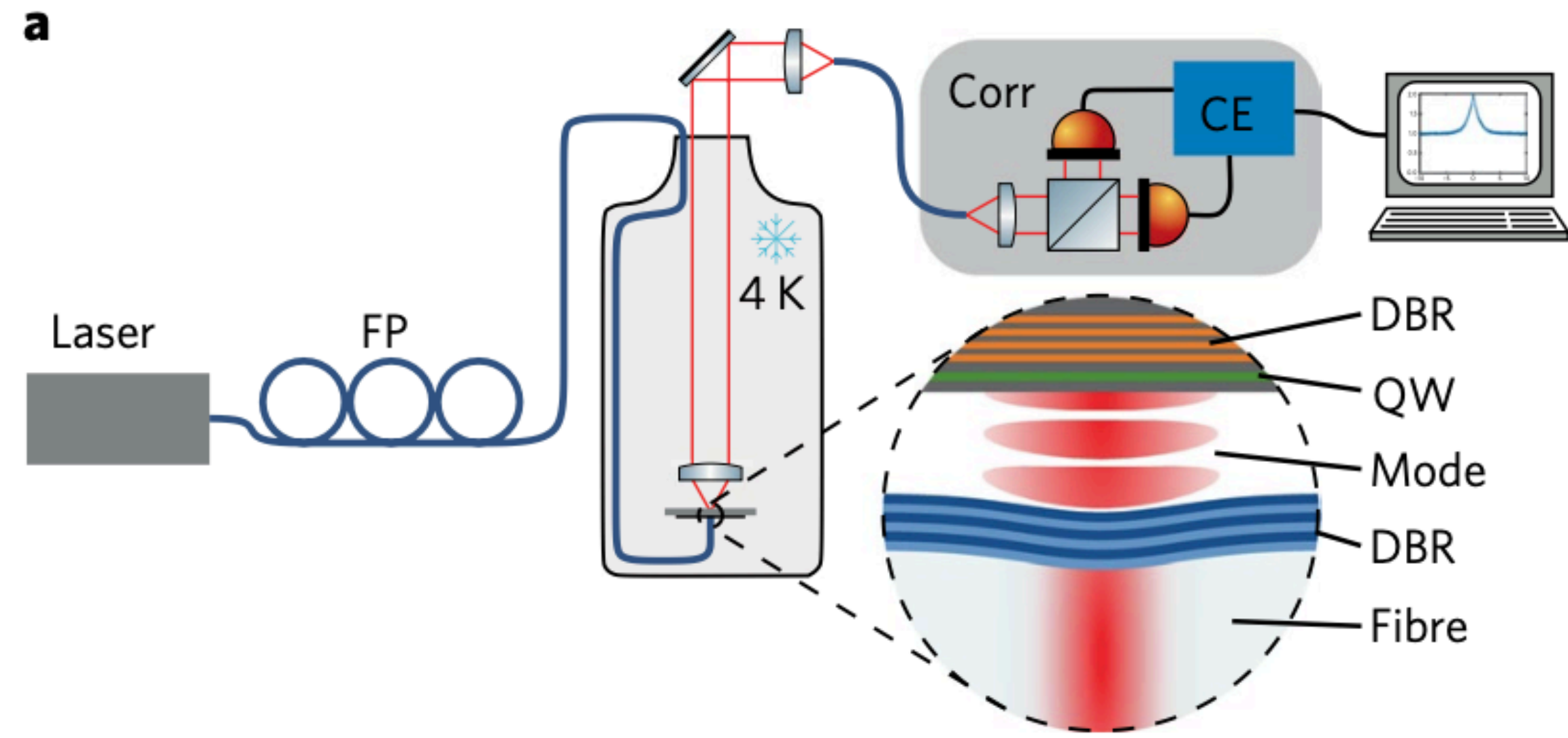
$$D \sim (1/U)^2$$

- In the discontinuous transition ( $\Delta > 0$ )

$$D \sim e^{1/U}$$



GaAs cavity polaritons.





# Homodyne current

- Mix photon output with a strong laser source  $\alpha = |\alpha| e^{i\phi}$ .
- Equivalent to measuring jumps of  $(a + \alpha)\rho(a + \alpha)^\dagger$ , where  $\alpha$  is a large number.

$$J = \kappa \langle (a + \alpha)^\dagger (a + \alpha) \rangle = \kappa \left( |\alpha|^2 + |\alpha| \langle a e^{-i\phi} + a^\dagger e^{i\phi} \rangle + \langle a^\dagger a \rangle \right)$$

- $|\alpha|^2$  is just a constant offset.
- If  $\alpha$  is large, then the current will predominantly

$$x := a e^{-i\phi} + a^\dagger e^{i\phi}$$

instead of  $a^\dagger a$ .

- Quantum diffusion unravelling:

$$d\rho = dt \mathcal{L} \rho + dW [\mathcal{H} \rho - \langle x \rangle \rho],$$

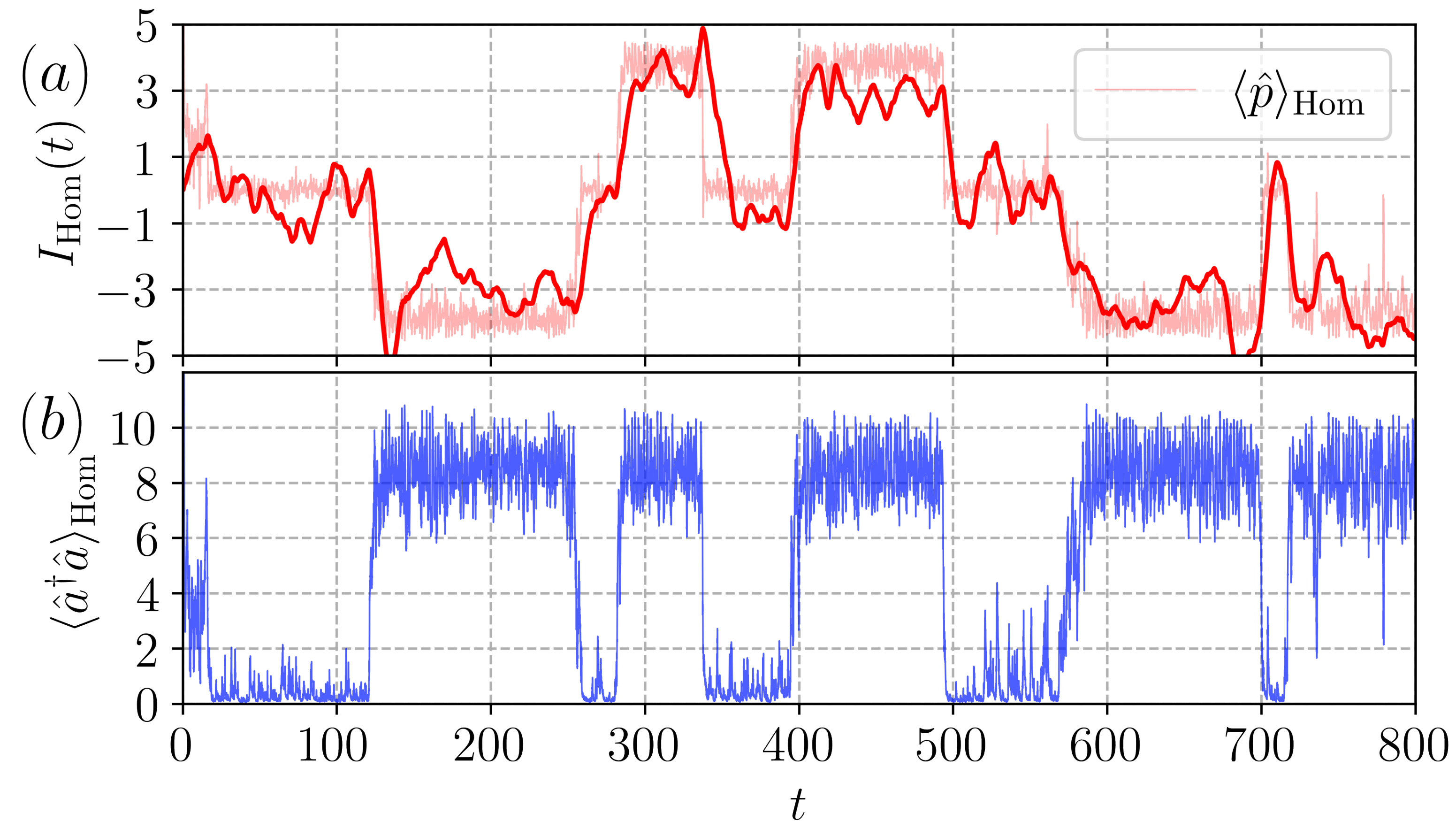
$$\mathcal{H} \rho = \kappa (a \rho + \rho a^\dagger)$$

$dW$  = Wiener increment:

$$E(dW) = 0, \quad dW^2 = dt$$

# Homodyne current (in $p = i(a^\dagger - a)$ )

- The homodyne current switches between 3 values (+,0,-).
- Captures the tunneling between the 3 blobs.



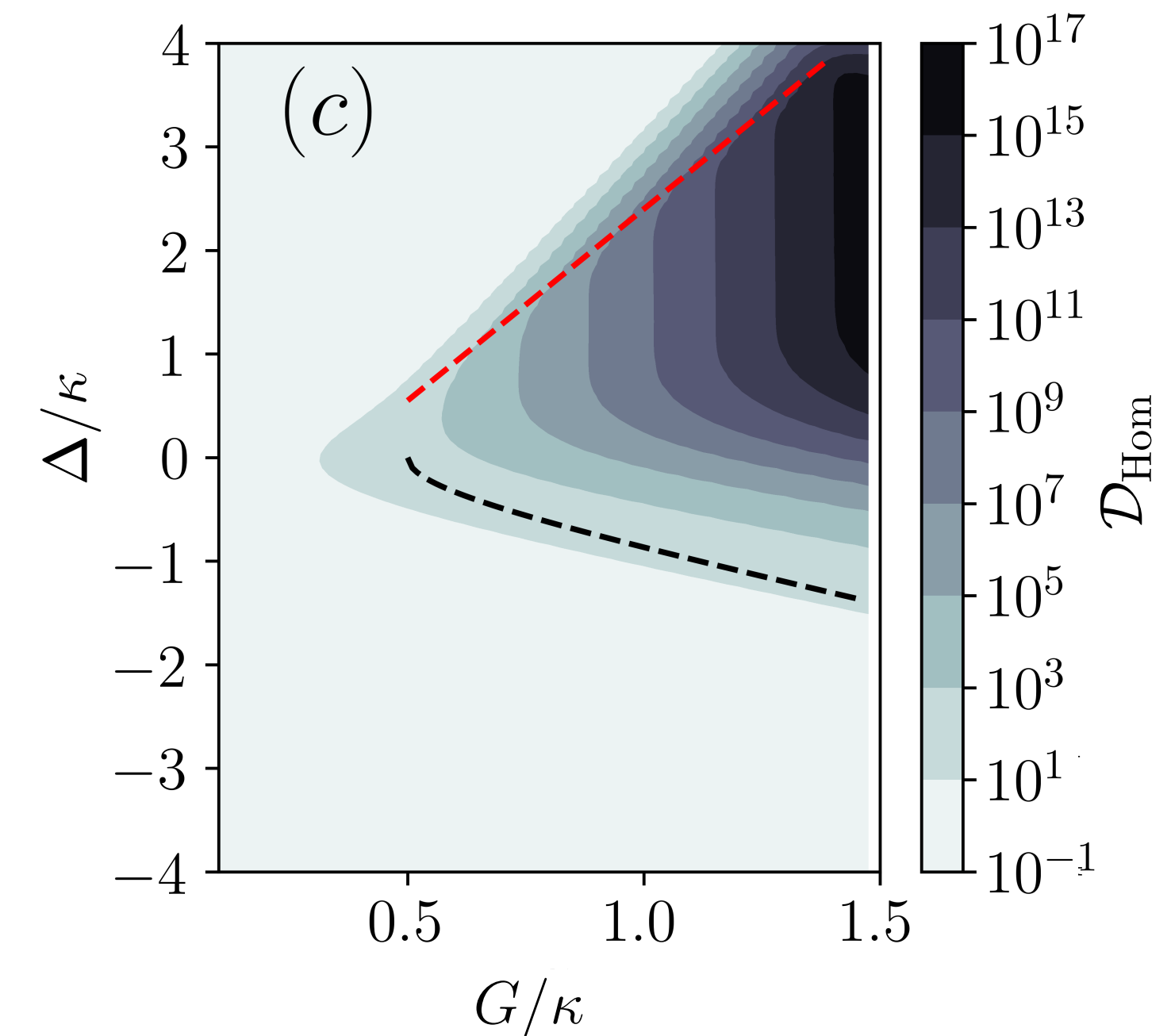
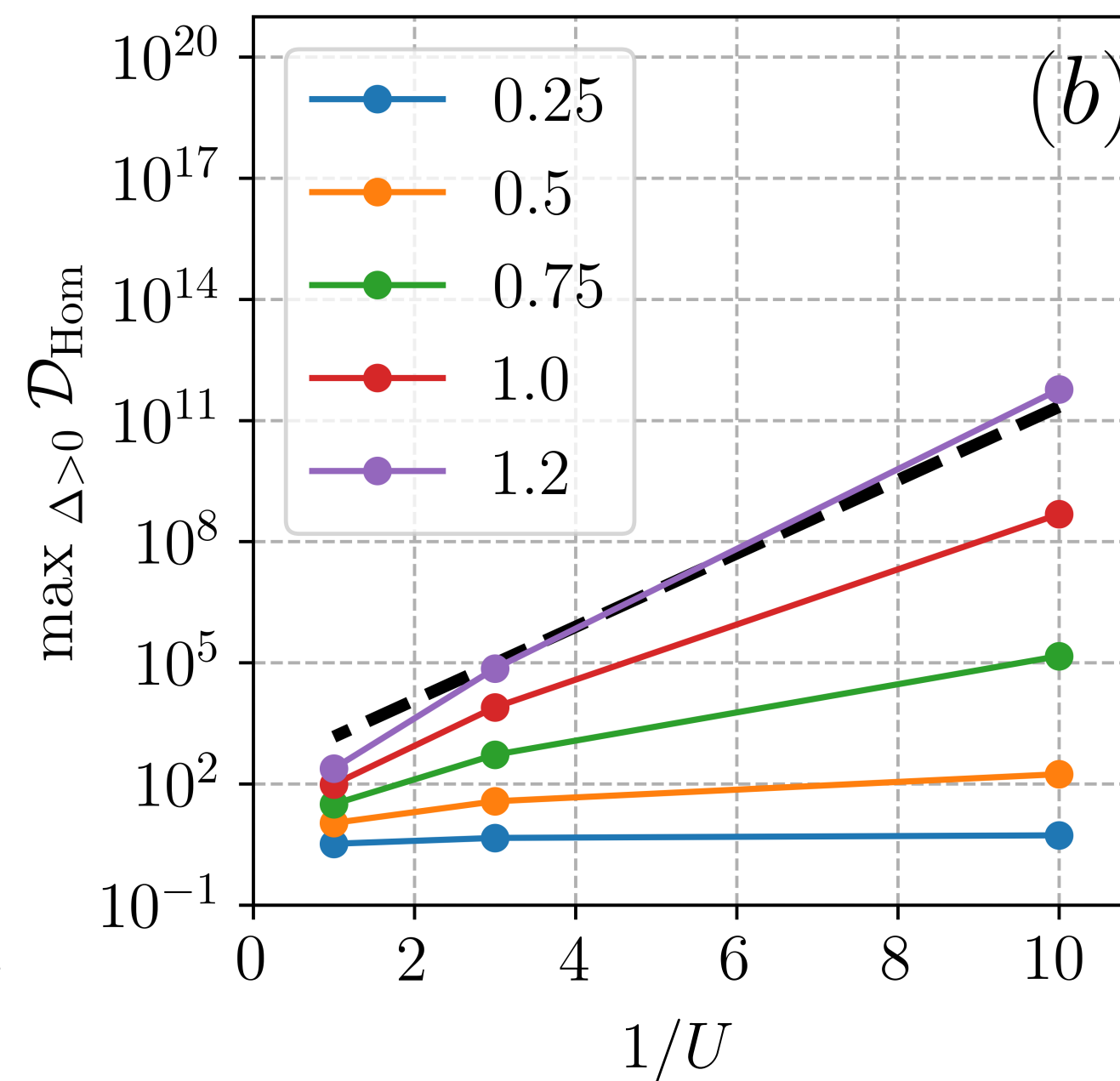
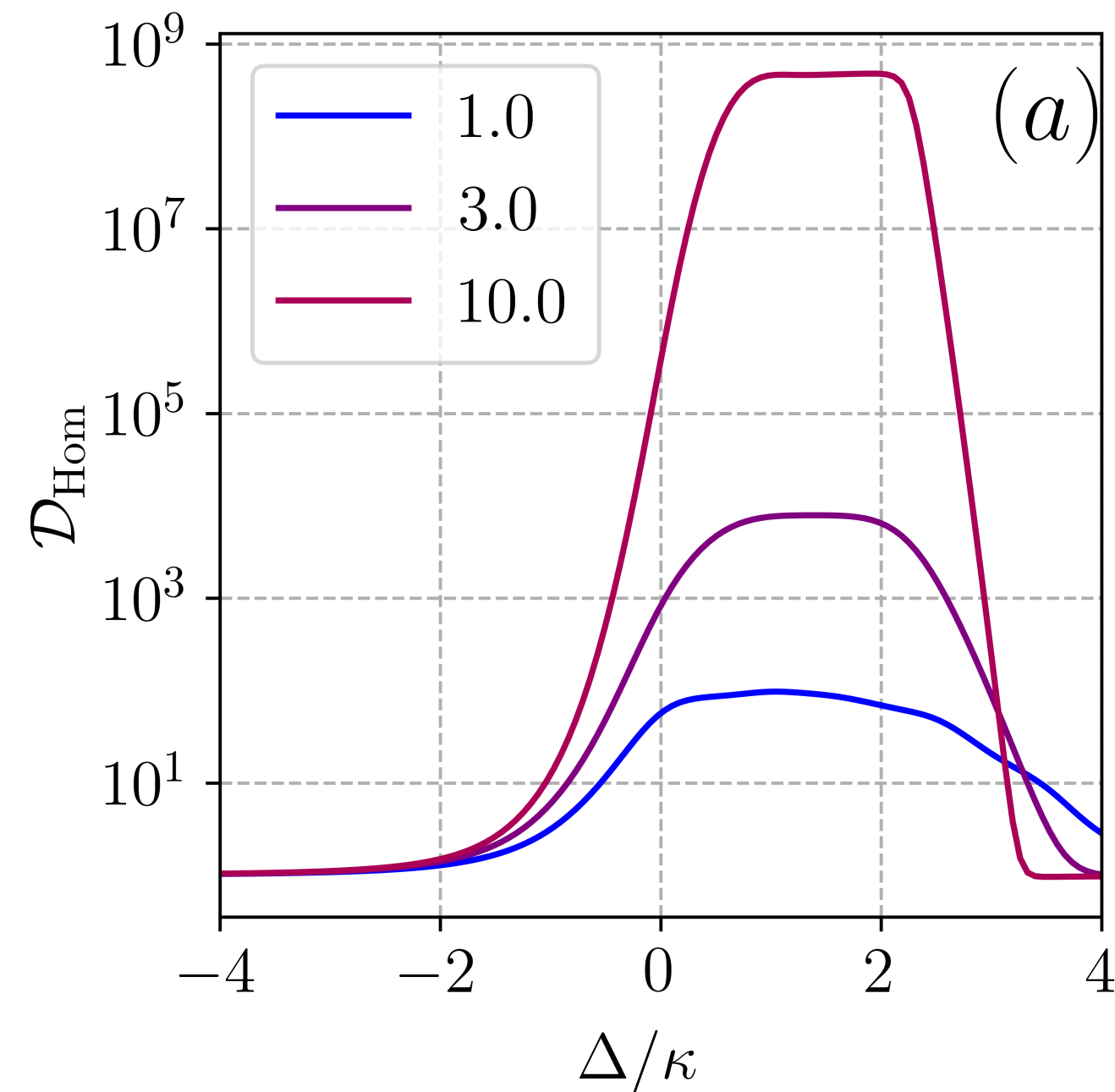
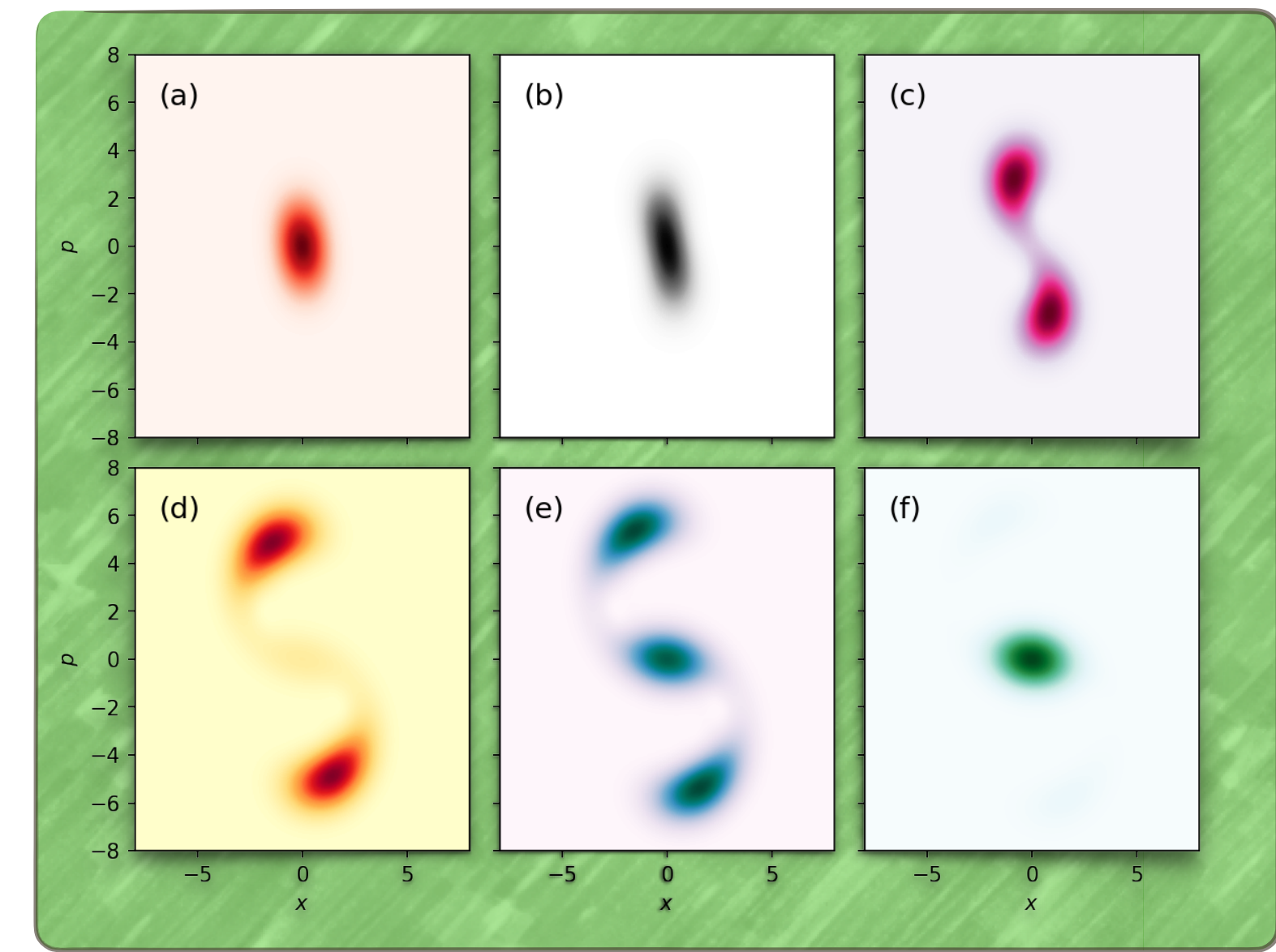


# Divergence of the diffusion coefficient

- Homodyne current noise diverges exponentially in a much broader region.

$$D \sim e^{1/U}$$

- Reflects sensitivity to all 3 blobs.



# What information the current conveys about the system?

Gabriel T. Landi, Mauro Paternostro, and Alessio Belenchia,

**“Informational steady-states and conditional entropy production in continuously monitored systems”**

PRX Quantum 3, 010303 (2022)

Massimiliano Rossi, Luca Mancino, Gabriel T. Landi, Mauro Paternostro, Albert Schliesser, Alessio Belenchia

**“Experimental assessment of entropy production in a continuously measured mechanical resonator”**

Phys. Rev. Lett. 125, 080601 (2020)



Mauro Paternostro

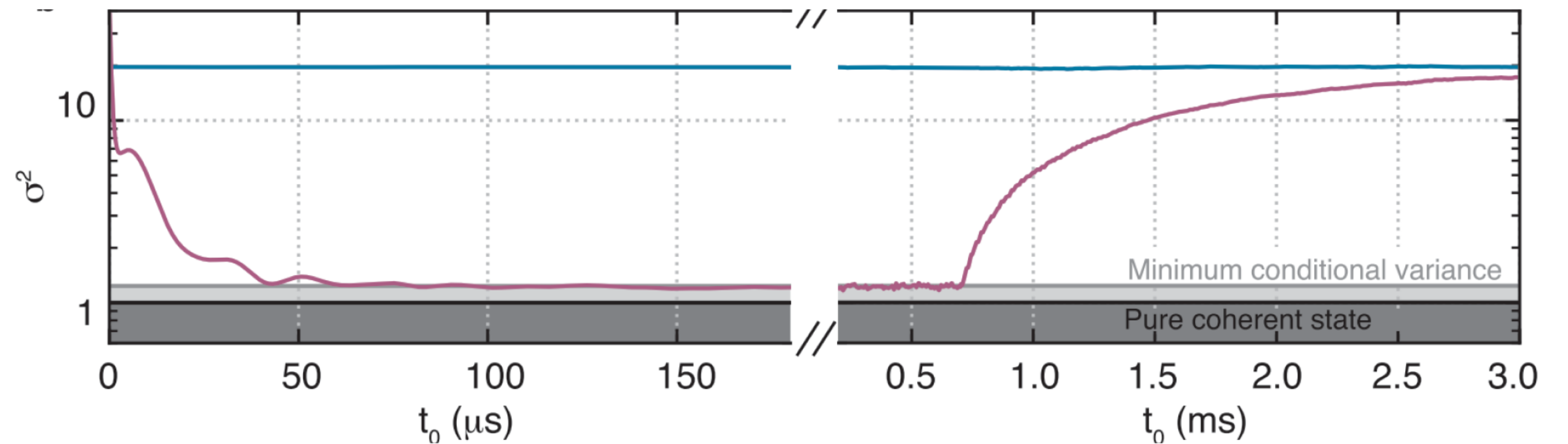
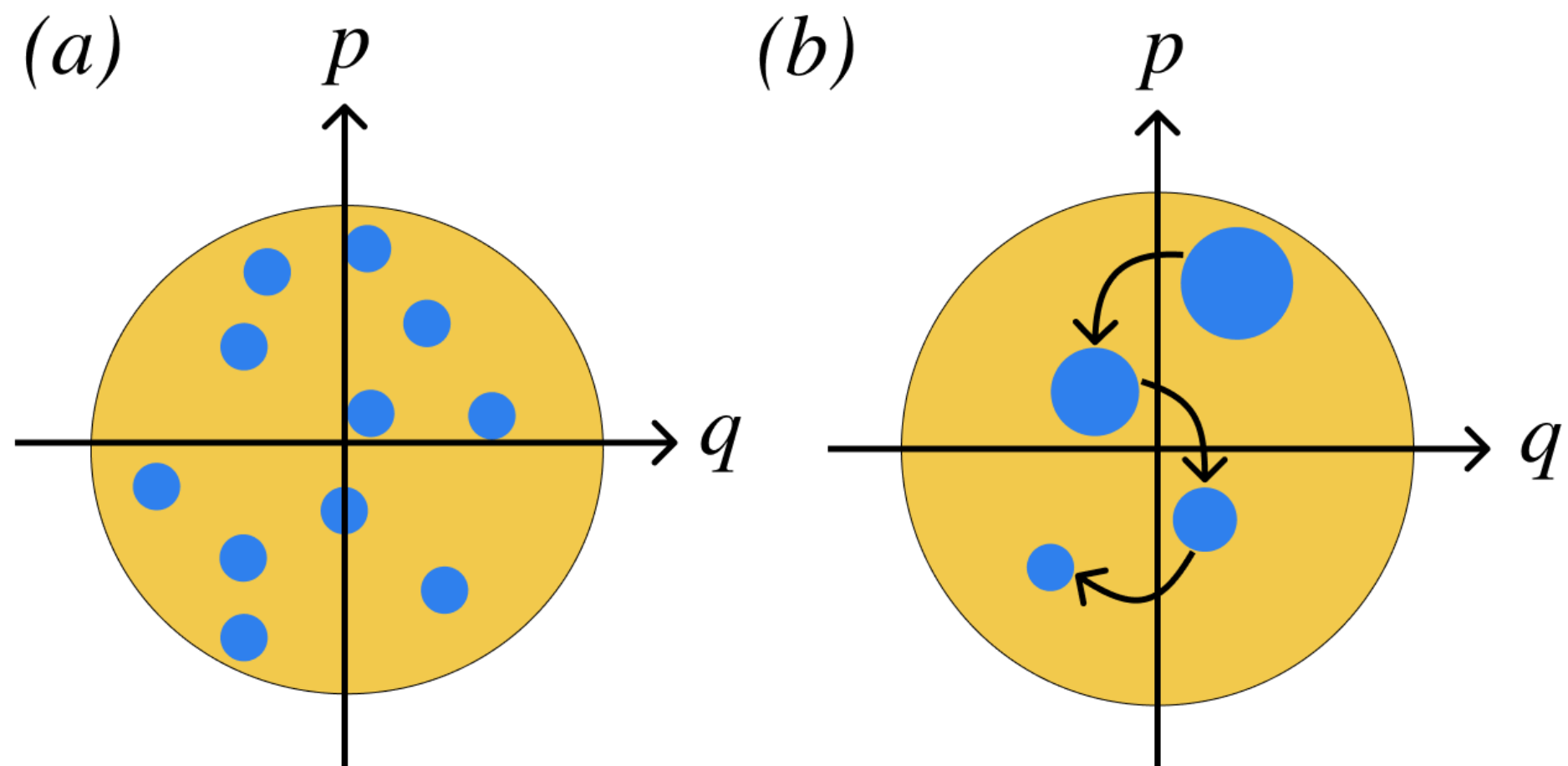
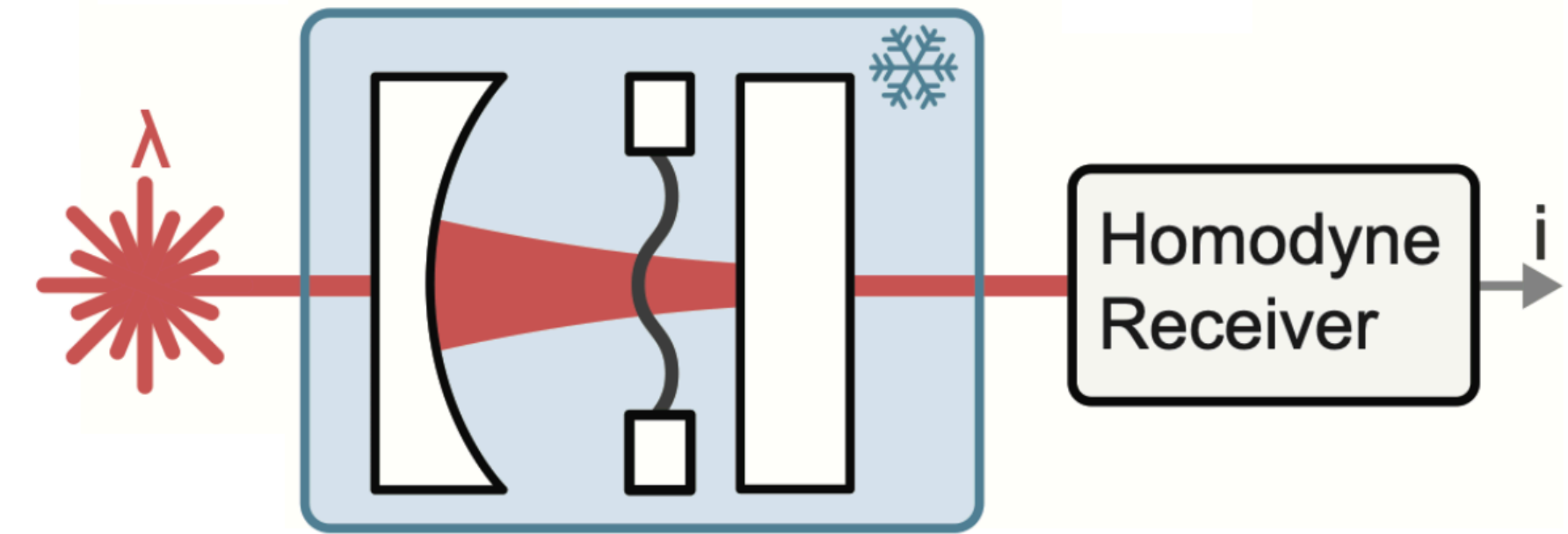


Alessio Belenchia

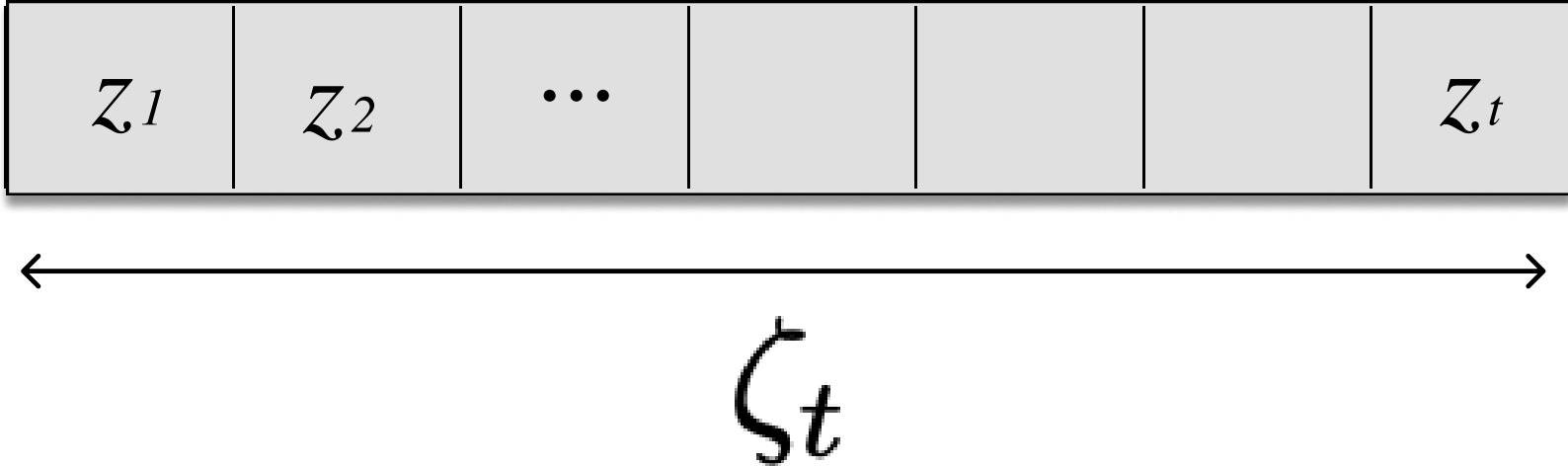
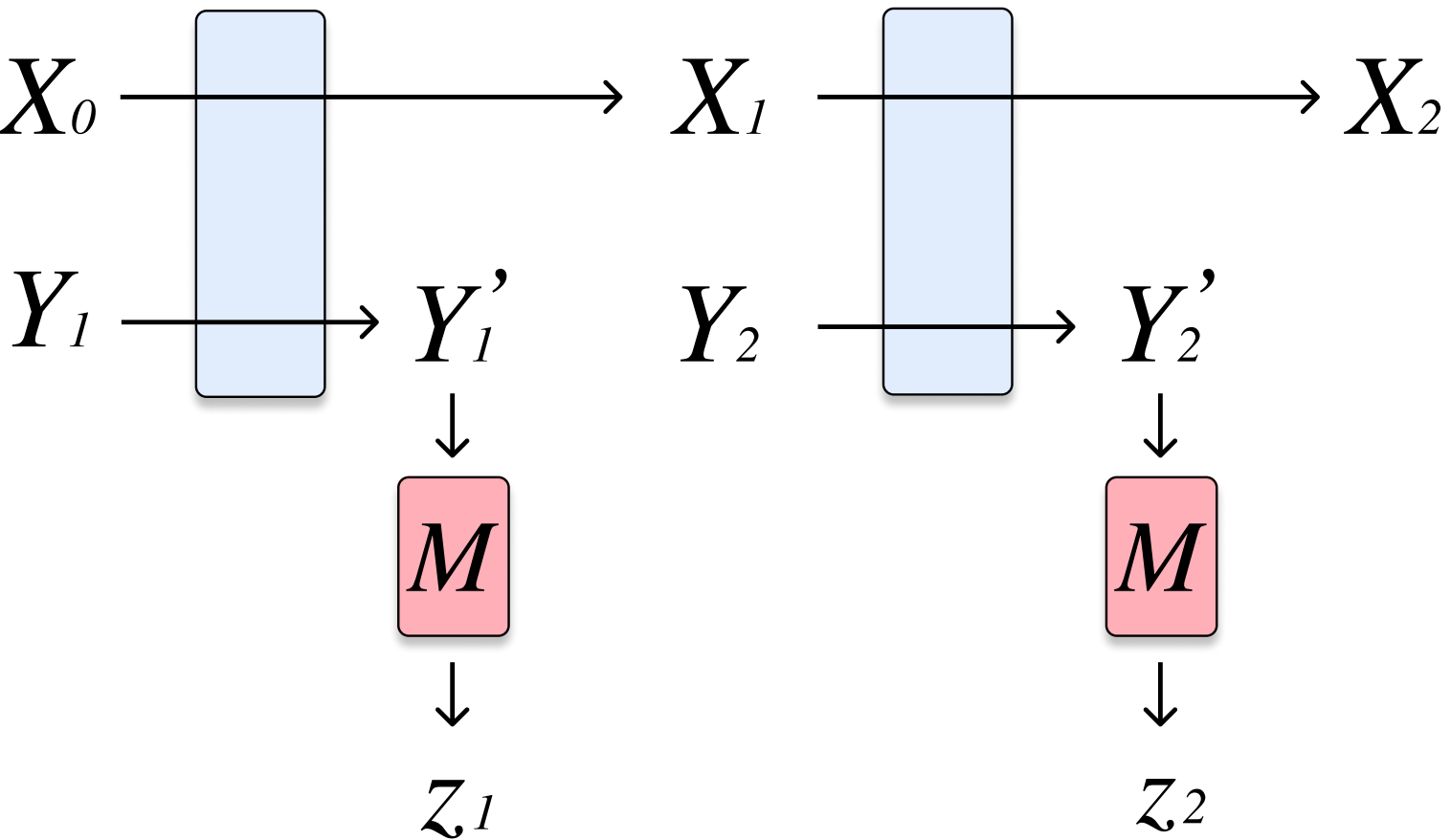
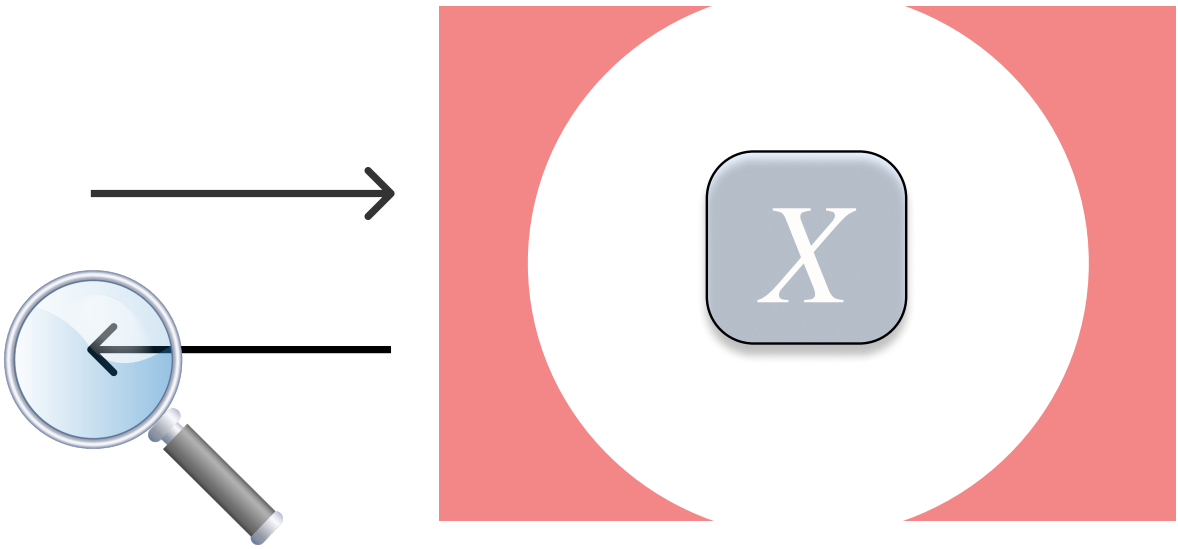


# Optomechanical experiment

- $\sigma(t)$  = variance of the position of a mechanical membrane.
- Start measuring at  $t = 0$ .
  - Use information to push oscillator to the middle.
    - Reduce uncertainty.
      - *Informationally driven cooling.*



# Collision model



- Unconditional dynamics (no measurement):

$$\rho_{X_t} = \mathcal{E}(\rho_{X_{t-1}}) = \text{tr}_{Y_t} \left\{ U_t(\rho_{X_{t-1}} \otimes \rho_{Y_t}) U_t^\dagger \right\}$$

- Entropy:  $S(X_t) = - \text{tr} \{ \rho_{X_t} \ln \rho_{X_t} \}$

- Conditional dynamics, *given the measurement outcomes* (unnormalized)

$$\rho_{X_t | \zeta_t} = \mathcal{E}_{z_t}(\rho_{X_{t-1} | \zeta_{t-1}}) = \text{tr}_{Y_t} \left\{ M_{z_t} U_t(\rho_{X_{t-1} | \zeta_{t-1}} \otimes \rho_{Y_t}) U_t^\dagger M_{z_t}^\dagger \right\}$$

$P(\zeta_t)$  = probability of a trajectory  $\zeta_t$ .

$$\rho_{X_t} = \sum_{\zeta_t} P(\zeta_t) \rho_{X_t | \zeta_t}$$

- Quantum-classical conditional entropy:

$$S(X_t | \zeta_t) = \sum_{\zeta_t} P(\zeta_t) S(\rho_{X_t | \zeta_t})$$

# Holevo information

$$I(X_t : \zeta_t) = S(X_t) - S(X_t | \zeta_t) = \sum_{\zeta_t} P(\zeta_t) D(\rho_{X_t|\zeta_t} || \rho_{X_t}) \geq 0$$

- Represents the *net* amount of information gained about X over the entire trajectory  $\zeta_t$ .
- Change can have any sign:  $\Delta I_t = I(X_t : \zeta_t) - I(X_{t-1}, \zeta_{t-1}) \lesseqgtr 0$ .
- Further split into a Holevo information gain/loss:

$$\Delta I_t = G_t - L_t$$

$$G_t = I(X_t : z_t | \zeta_{t-1}) = I(X_t : \zeta_t) - I(X_t : \zeta_{t-1}) \geq 0$$

$$L_t = I(X_{t-1} : \zeta_{t-1}) - I(X_t : \zeta_{t-1}) \geq 0$$

**Informational steady-state:**

$$\Delta I_{ISS} = 0$$

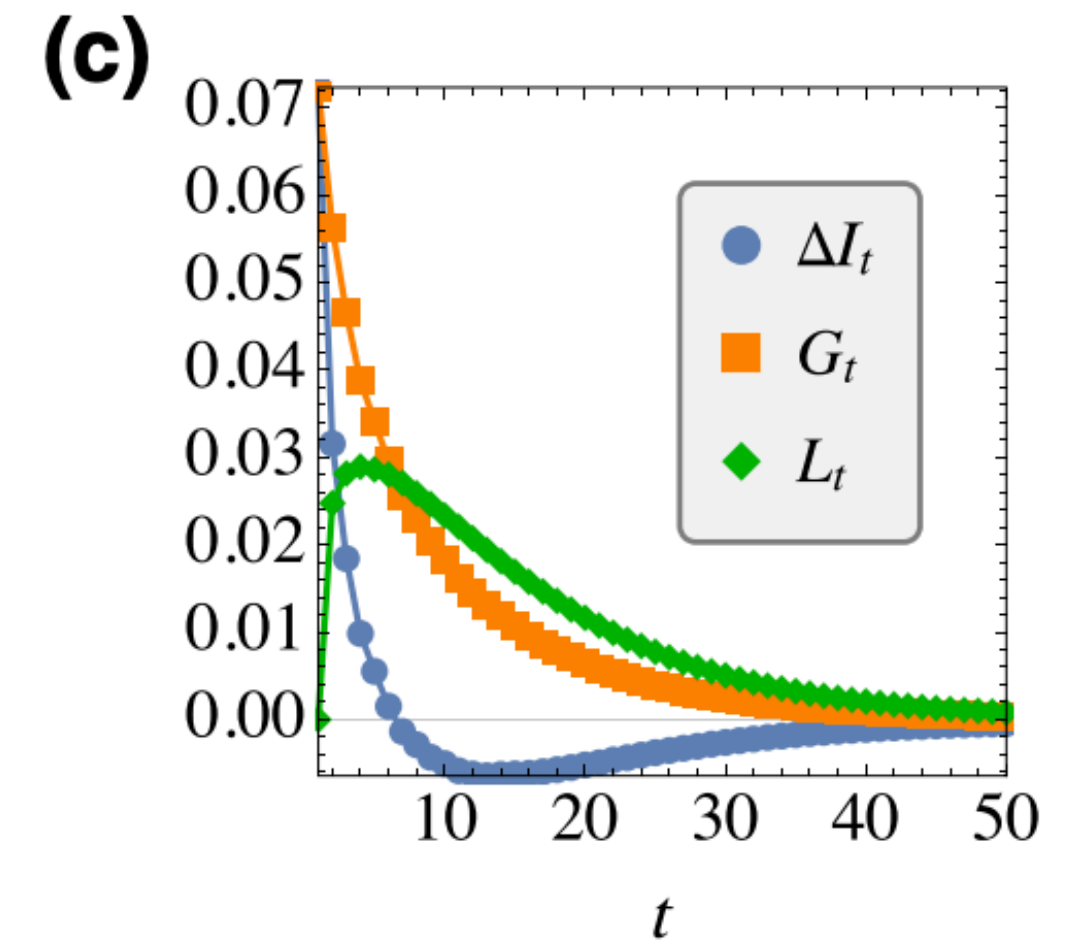
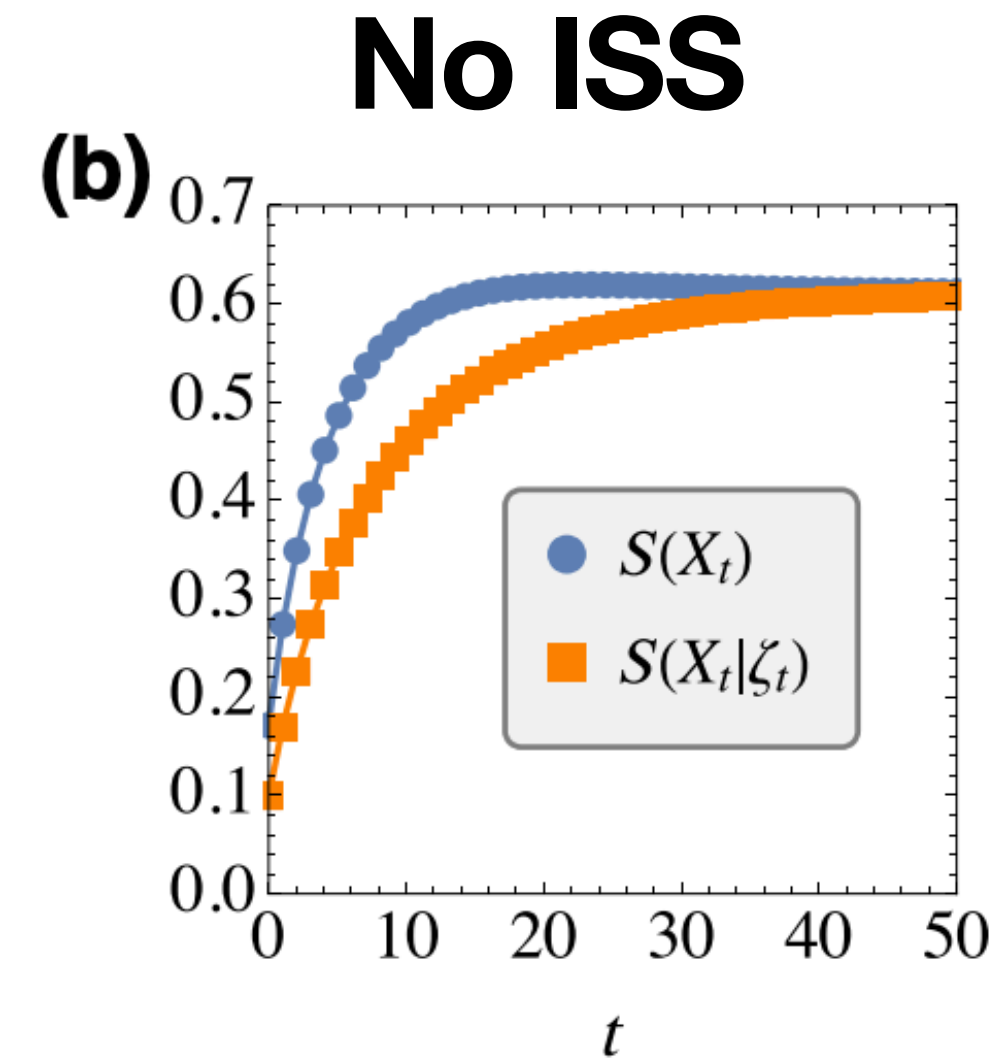
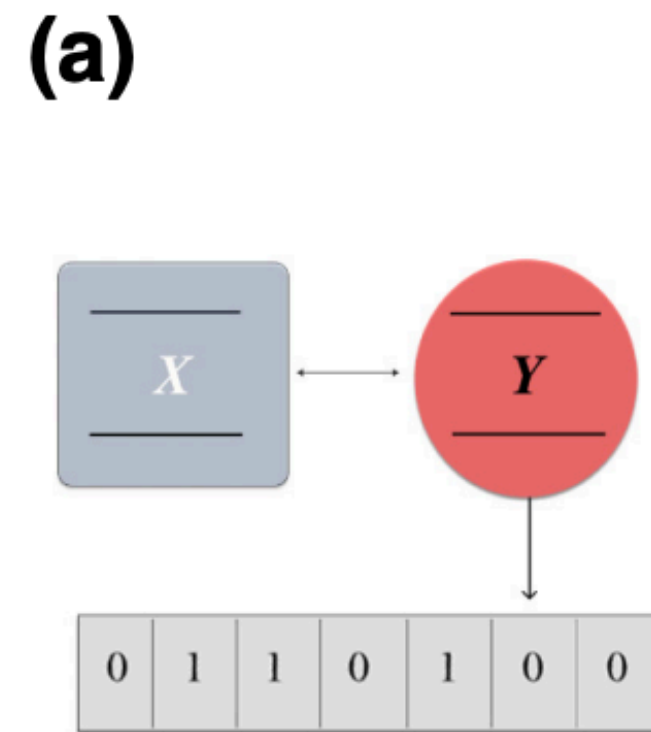
but

$$G_{SS} = L_{SS} \neq 0.$$

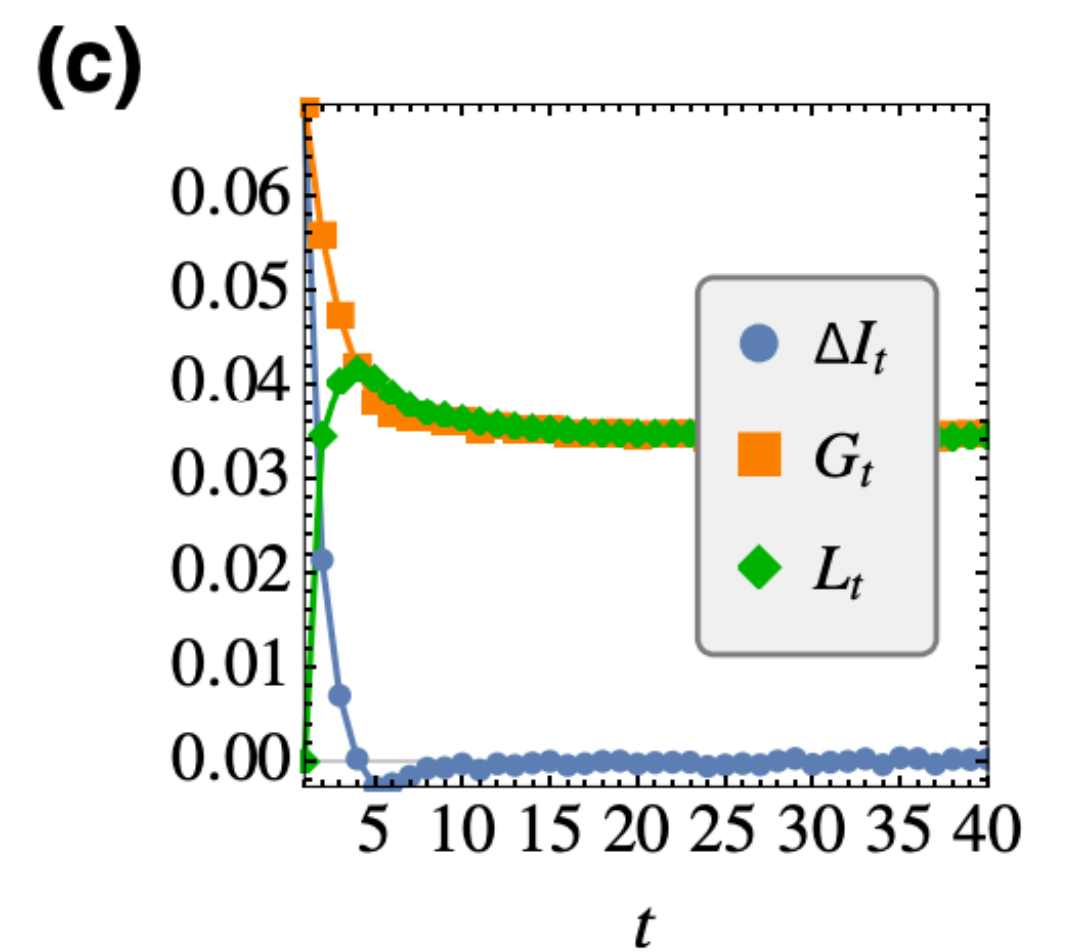
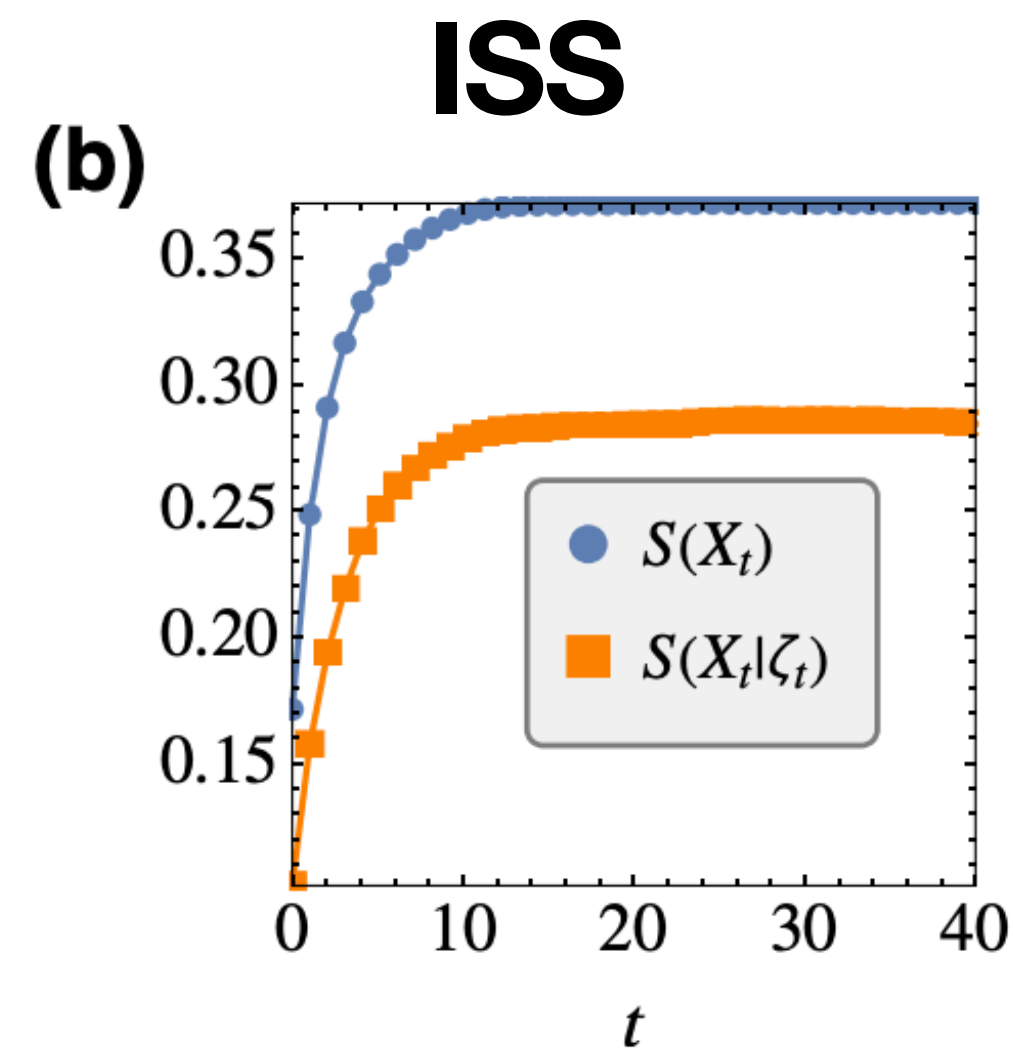
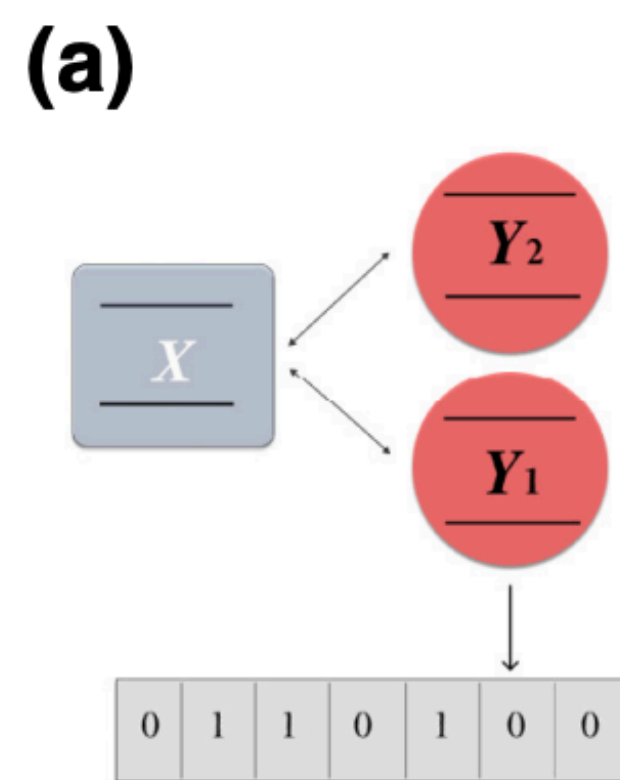


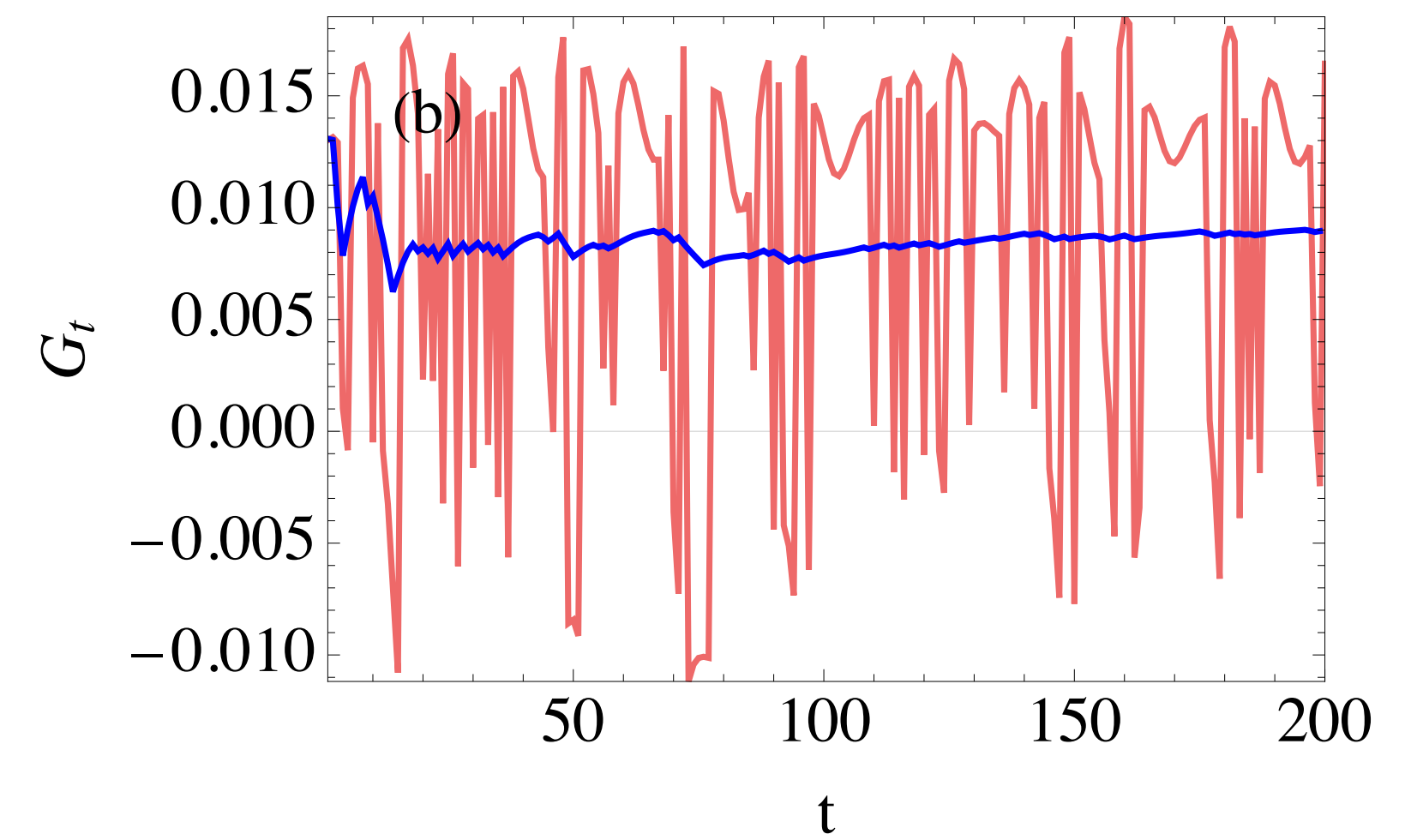
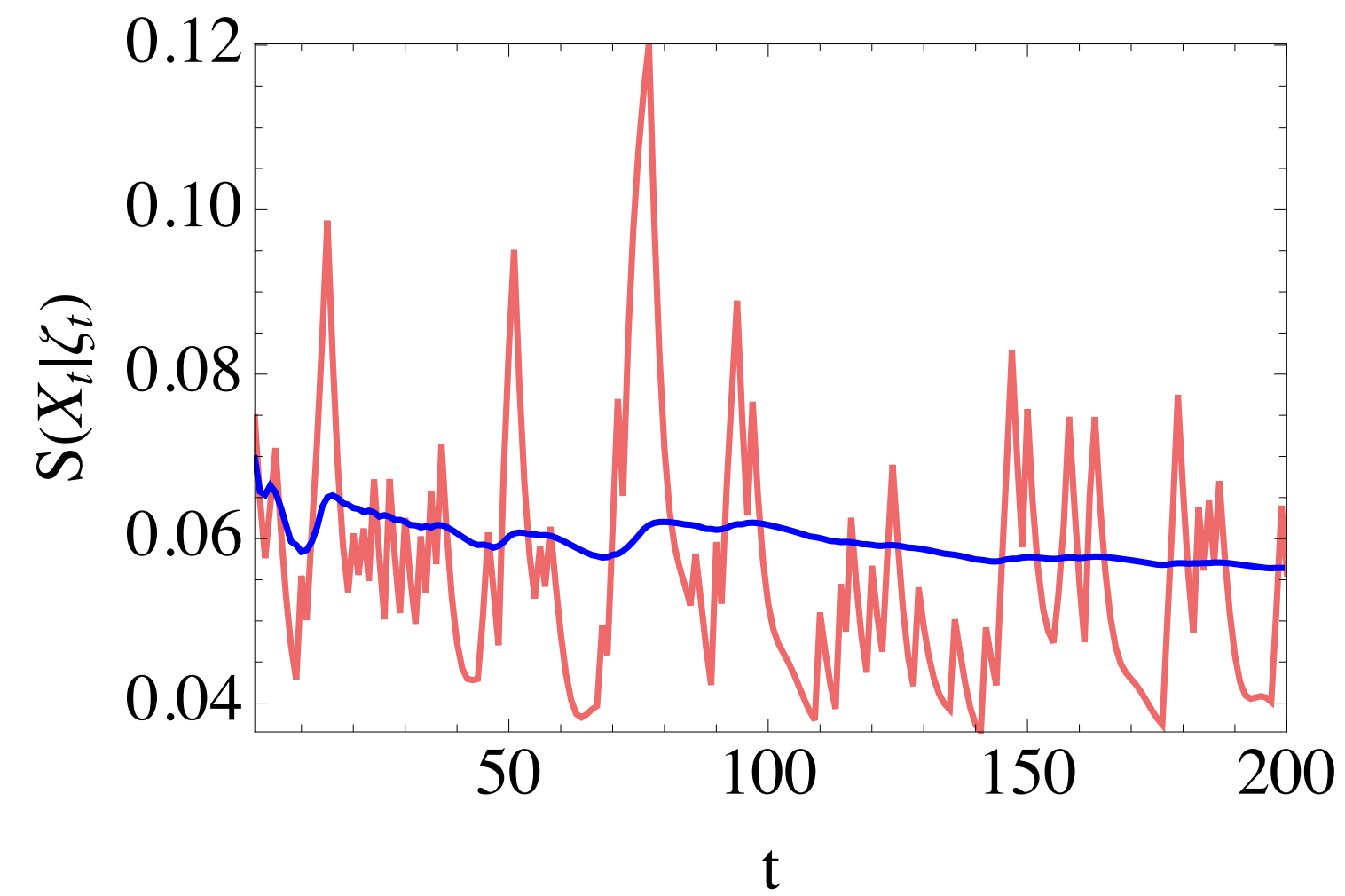
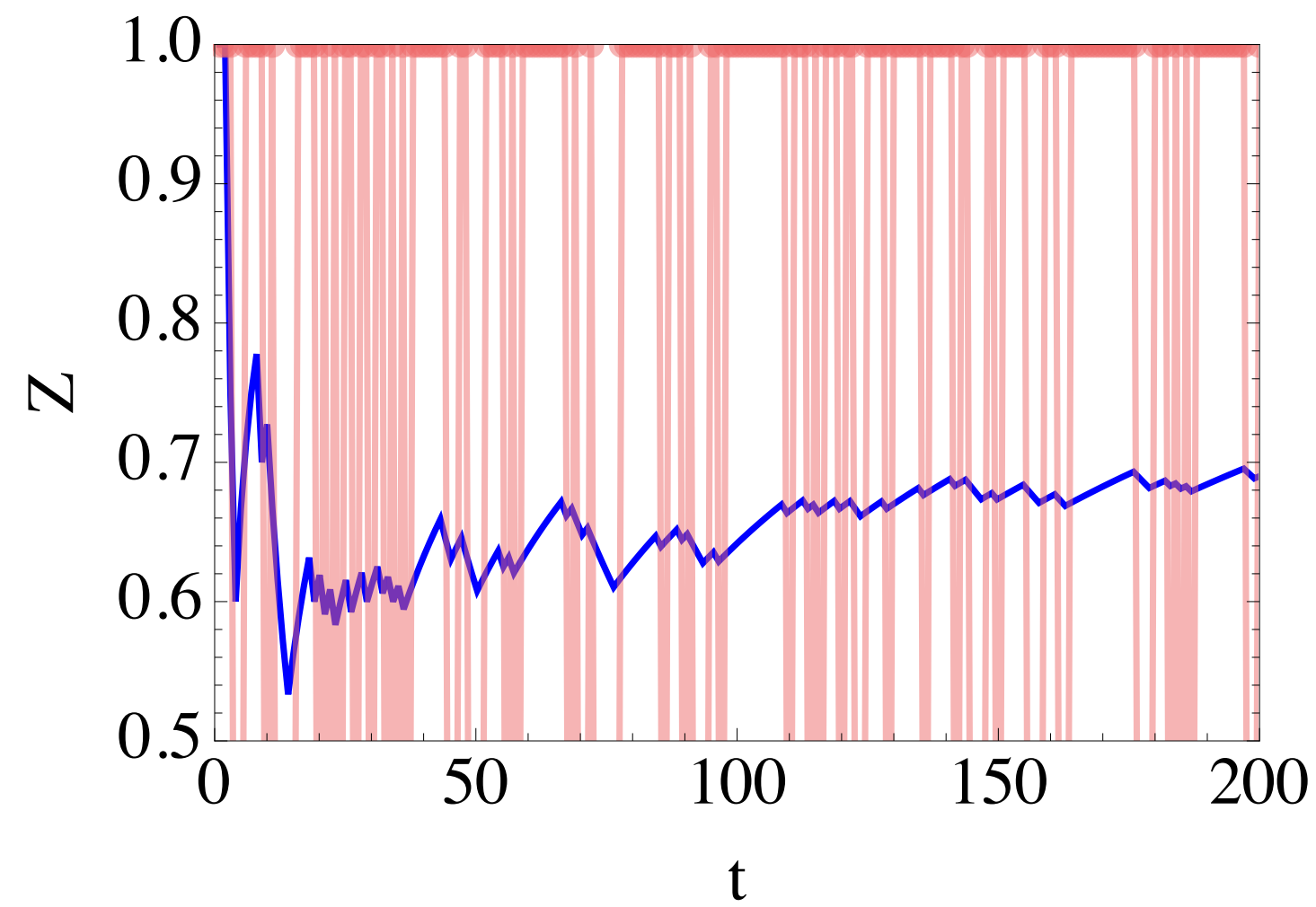
# Examples

- Long-time limit no additional information is acquired.



- Measurements continue to gain information.
- But gain & loss balance out.





- Blue = accumulated average.
- $Z$  is the measurement outcome.
- System reaches an ISS because  $G_t$  tends to a finite value on average.

# Optomechanical experiment

- Unconditional dynamics: interplay between measurement backaction and thermal bath:

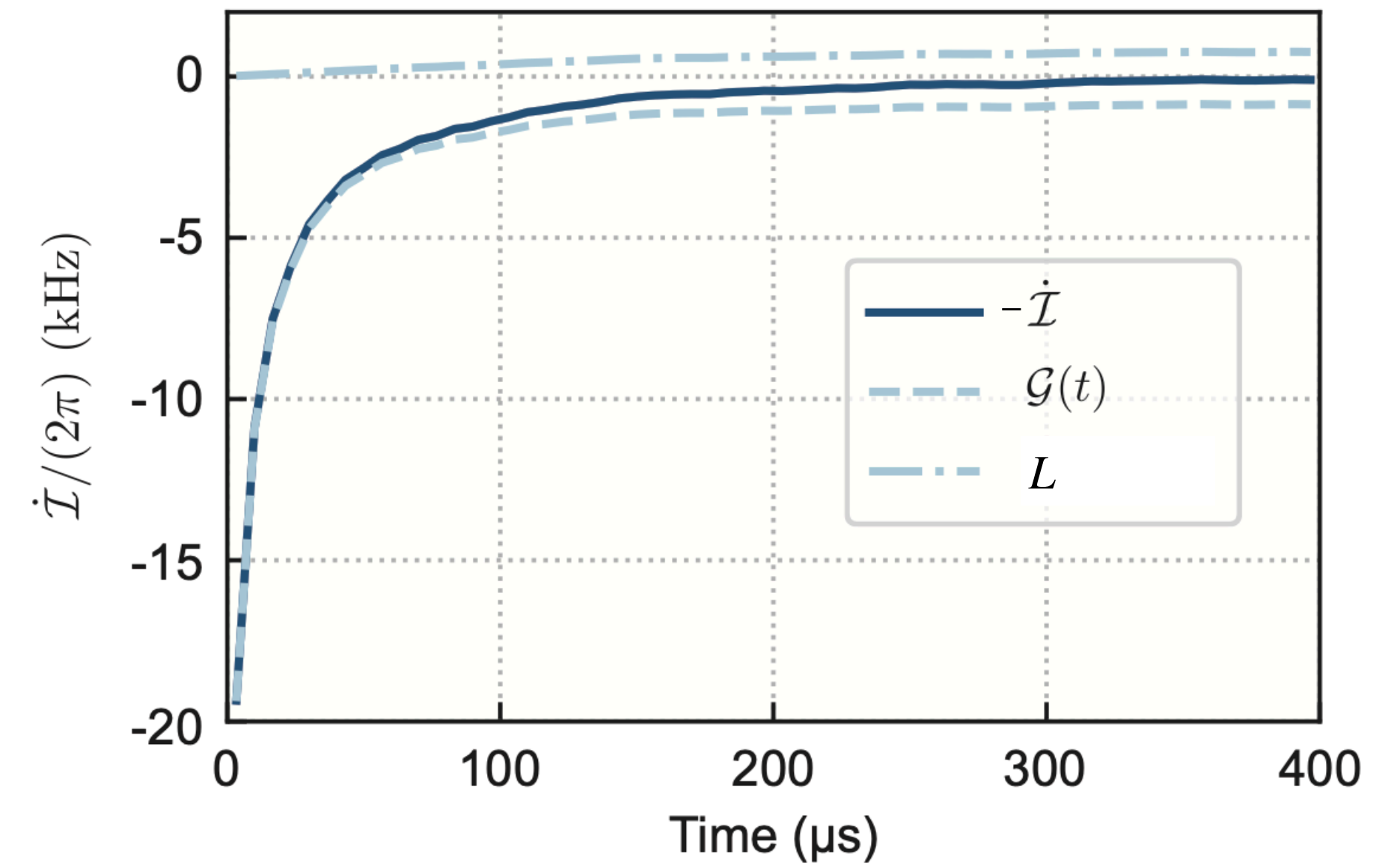
$$\sigma_u := \langle \Delta x^2 \rangle = \bar{n} + 1/2 + \Gamma_{\text{qba}}/\Gamma_m$$

- Stochastic evolution of 1st & 2nd moments:

$$\frac{dx_c}{dt} = -\frac{\Gamma_m}{2}x_c + \sqrt{4\eta\Gamma_{\text{qba}}}\sigma_c(t)\xi(t)$$

$$\frac{d\sigma_c}{dt} = \Gamma_m(\sigma_u - \sigma_c) - 4\eta\Gamma_{\text{qba}}\sigma_c^2$$

Information gain/loss rates  
characterizing the information  
steady-state

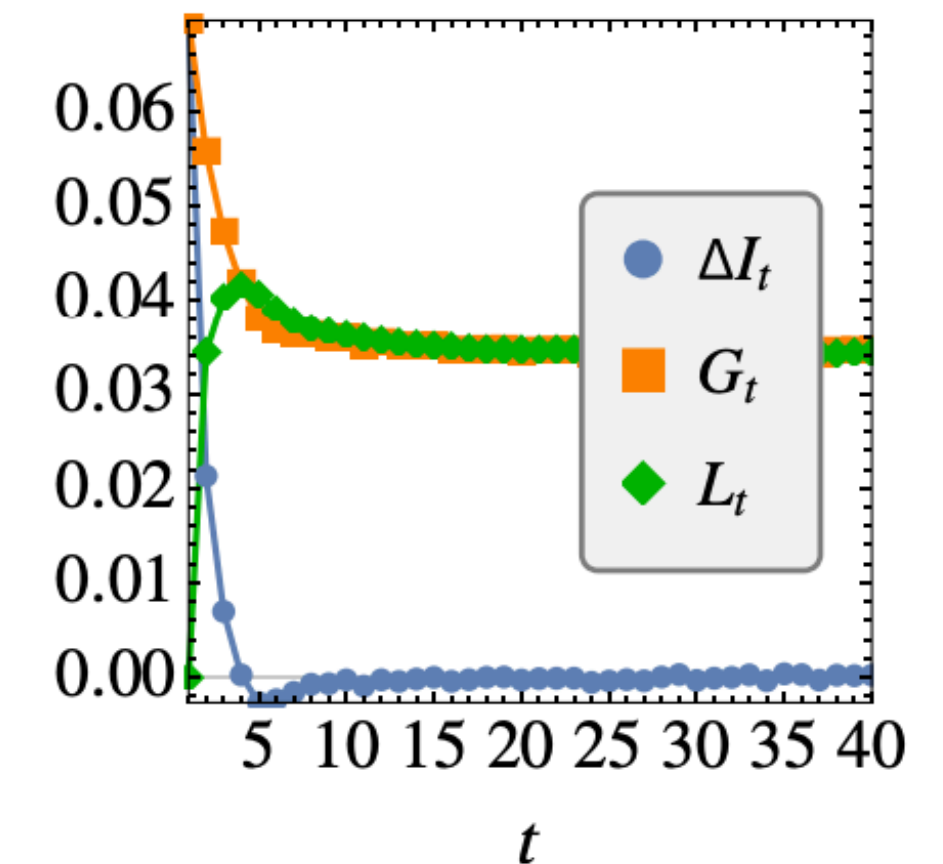
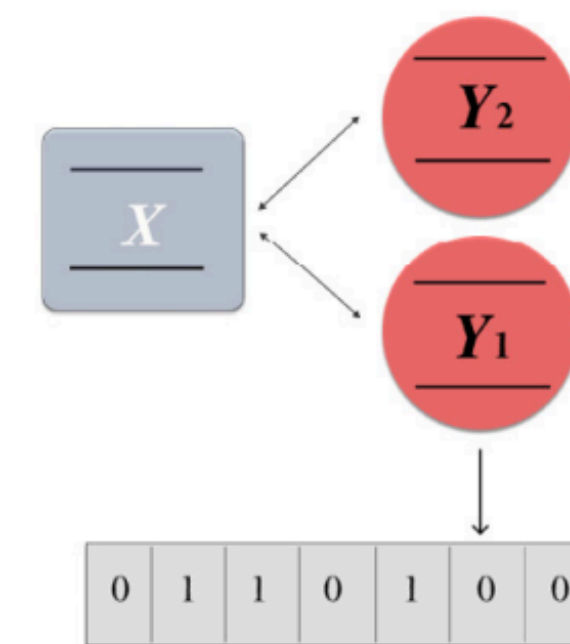
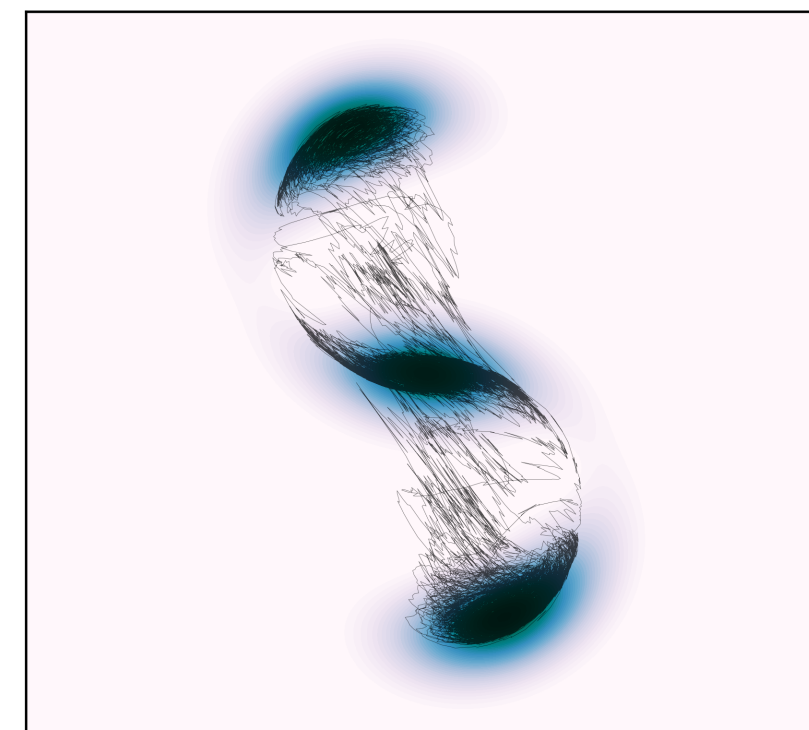
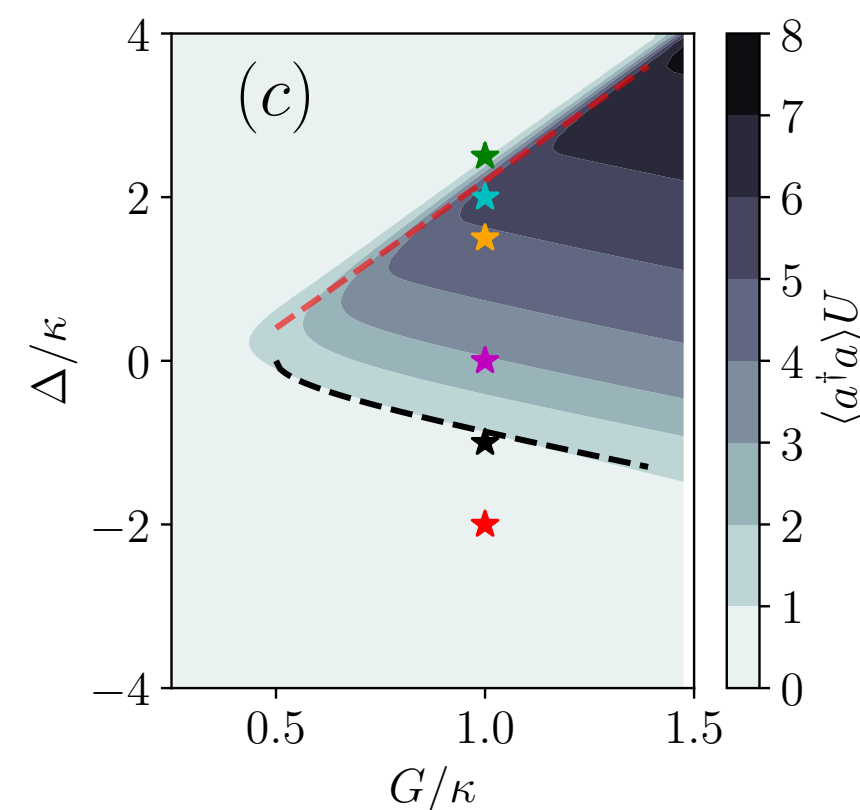




# Conclusions

- Fluctuations in the time domain:
  - NESS beyond averages
- Connection between quantum optics and full counting statistics.
- Parametric Kerr model: critical properties of the fluctuations shed light on the Nature of the transitions
  - Photo-detection: exponential divergence in the discontinuous transition.
  - Homodyne: exponential divergence in the entire critical region.

- Information gain/loss in continuous measurements.
  - Modeled using a Collision Model.
- Holevo information: cumulative knowledge we acquired about the system.
- Decomposable into Gain - Loss.
- Optomechanical experiment.





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# University of Rochester

## Post-doc positions 2023

My group is moving to UofR next year, and we have open post-doc positions to work on

- theory of quantum thermodynamics
- open quantum systems
- quantum information

For more information:

[www.fmt.if.usp.br/~gtlandi](http://www.fmt.if.usp.br/~gtlandi)



Thank you.



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