## Observations of Stars

Measurable properties of stars<br>Stellar photometry

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University of Rochester

## Observations of stars

Today's topics:

- Ensemble averages
- Measurable properties of stars:

P Parallax and the distance to the nearest stars

- Magnitude, flux, and luminosity

Blackbodies as approximations to stellar spectra
$>$ Stellar properties from photometry
$>$ Color and temperature

- Bolometric correction and bolometric magnitude
Reading: Kutner Ch. 2, Ryden Sec. 5.7 \& 13.1-13.4


Open cluster M39 in Cygnus

## Averages: Binary star systems

## Example

- If the stars are too close together to observe their orientation, then $i$ cannot be determined.
- Still, an estimate of $\sin i$ would be useful. Could we make a reasonable assumption?
- Let us try an average value, assuming that binary star orbits are in general uniformly distributed. That is, all orientations are equally likely:

Orbital axis


$$
\begin{aligned}
p(i) & =\text { constant }=C, \quad i=0 \rightarrow \frac{\pi}{2} \\
1 & =\int_{0}^{\pi / 2} p(i) d i=C \int_{0}^{\pi / 2} d i=\frac{\pi C}{2} \\
\therefore C & =\frac{2}{\pi}
\end{aligned}
$$

## Averages: Binary star systems

## Example

Now that we know $p(i)$,

$$
\begin{aligned}
\langle\sin i\rangle & =\int_{0}^{\pi / 2} \sin i p(i) d i=\frac{2}{\pi} \int_{0}^{\pi / 2} \sin i d i \\
& =\left[-\frac{2}{\pi} \cos i\right]_{0}^{\pi / 2}=-\frac{2}{\pi}(-1-0)=\frac{2}{\pi}
\end{aligned}
$$

Thus, a reasonable estimate of the true orbital speed $v$ given the measured line-of-sight speed $v_{r}$ is

$$
v=\frac{v_{r}}{\langle\sin i\rangle}=\frac{\pi}{2} v_{r}
$$

Note that a more sophisticated treatment with axes randomly oriented yields $\langle\sin i\rangle=\pi / 4 \ldots$

## Averages: Binary star systems

## Example

Allowing all orientations requires us to average over the full sky $(\Omega)$ defined by $\theta, \phi$ :

$$
\begin{aligned}
p(\theta, \phi) & =\text { constant }=C \\
1 & =\int_{0}^{2 \pi} d \phi \int_{0}^{\pi} \sin \theta d \theta p(\theta, \phi)=2 \pi \cdot 2 \cdot C \quad \therefore p(\theta, \phi)=\frac{1}{4 \pi}
\end{aligned}
$$

$$
\begin{aligned}
\langle\sin i\rangle & =\int_{\Omega} d \Omega p(\theta, \phi) \sin \theta=\frac{1}{4 \pi} \int_{0}^{2 \pi} d \phi \int_{0}^{\pi} \sin ^{2} \theta d \theta \\
& =\frac{1}{2} \int_{0}^{\pi} \frac{1-\cos 2 \theta}{2} d \theta \\
& =\frac{1}{4}\left[\theta-\frac{1}{2} \sin 2 \theta\right]_{0}^{\pi}=\frac{\pi}{4}
\end{aligned}
$$

## What do we want to know about stars?

As a function of mass and age:

- Structure of their interiors
- Structure of their atmospheres
- Chemical composition
- Power output and spectrum
- Spectrum of oscillations
- All the ways they can form
- All the ways they can die
- The details of their lives in between
- The nature of their end states
- Their role in the dynamics and evolution of their host galaxy

For visible stars, measure. .

- Flux at all wavelengths, at high resolution (probes atmosphere)
- Magnetic fields, rotation rate
- Oscillations in power output (probes interior)
- Distance from us and Galactic position
- Orbital parameters (for binaries)
... and match up the measurements with theory.


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... and match up the measurements with theory. We can do what is highlighted this semester.


## Flux and luminosity

- Primary observable quantity: flux, $f$, or power per unit area within some range of wavelengths.
- Surface brightness: flux at the surface of an emitting object.
- Luminosity: $L$, total power output.

Like most astronomical objects, stars emit light isotropically (same in all directions).

- Since the same total power $L$ must pass through all spheres centered on a star and is uniformly distributed on those spheres, the flux a distance $r$ away is

$$
f=\frac{L}{4 \pi r^{2}}
$$

- To obtain luminosity from measured flux, we must (1) add up flux measurements over all wavelengths, and (2) measure the distance.


## Distance from trigonometric parallax



## Trigonometric parallax



How small is the typical stellar parallax? Very small ( $<1$ ")

## Example

Suppose stars were typically 1 ly (light year) away - an extreme underestimate, as we will see. Then

$$
\begin{aligned}
\tan p & =p+\frac{p^{3}}{3}+\frac{2 p^{5}}{15}+\ldots \\
& =\frac{1 \mathrm{AU}}{r} \approx \frac{1.5 \times 10^{13} \mathrm{~cm}}{9.5 \times 10^{17} \mathrm{~cm}} \\
& =1.6 \times 10^{-5} \\
\tan p & \ll 1
\end{aligned}
$$

Therefore, $p \approx 1 \mathrm{AU} / r$ is an excellent approximation.

## Trigonometric parallax

- Note that in expressions such as

$$
\tan p \approx p \quad \text { and } \quad r=\frac{1 \mathrm{AU}}{p}
$$

$p$ is understood to be measured in radians.

- For convenience, parallaxes of nearby stars are usually reported in arcseconds or fractions thereof:

$$
\pi \text { radians }=180^{\circ}=180^{\circ} \times \frac{60 \operatorname{arcmin}}{1^{\circ}} \times \frac{60 \operatorname{arcsec}}{1 \operatorname{arcmin}}=6.48 \times 10^{5} \mathrm{arcsec}
$$

- The distance to an object with a parallax of 1 arcsecond is called a parsec (pc):

$$
1 \mathrm{pc}=3.0857 \times 10^{18} \mathrm{~cm}=3.2616 \mathrm{ly}
$$

- Useful to remember: $206265 \mathrm{arcsec} / \mathrm{rad}$ and $206265 \mathrm{AU} / \mathrm{pc}$.
- Smallest parallax measurements possible: about $1 \times 10^{-3}$ arcsec from the ground, and about $1 \times 10^{-5}$ arcsec (currently) from Gaia.


## Parallax: Mapping the galaxy

What is the largest distance that can be measured with ground-based and space-based telescopes using trigonometric parallax?
We can write the distance-parallax relation as

$$
r=\frac{1 \mathrm{AU}}{p[\mathrm{rad}]}=\frac{1 \mathrm{pc}}{p[\operatorname{arcsec}]}
$$

- Ground-based: $r_{\max }=\frac{1 \mathrm{pc}}{10^{-3^{\prime \prime}}}=1000 \mathrm{pc}=1 \mathrm{kpc}$
- Gaia: $r_{\max }=\frac{1 \mathrm{pc}}{10^{-5 \prime \prime}}=100,000 \mathrm{pc}=100 \mathrm{kpc}$

Note: We live 8 kpc from the center of our Galaxy.

## Parallax from Mars

## If we lived on Mars, what would be the numerical value of the parsec?

Mars's orbital semimajor axis is 1.524 AU , so the Martian "parsec" would be larger by a factor of 1.524:

$$
\begin{aligned}
r & =1 \text { parsec }=\frac{1.524 \mathrm{AU}}{1 \mathrm{arcsec}} \\
& =\frac{(1.524 \mathrm{AU})\left(1.496 \times 10^{13} \mathrm{~cm} / \mathrm{AU}\right)}{1 \mathrm{arcsec}}\left(\frac{3600 \mathrm{arcsec}}{1^{\circ}}\right)\left(\frac{180^{\circ}}{\pi}\right) \\
& =4.703 \times 10^{18} \mathrm{~cm}=4.971 \mathrm{ly}
\end{aligned}
$$

Note: Mars's orbit is twice as eccentric as Earth's, so be cautious about trusting those last decimal places.

## The thirty nearest stars

- There are 35 stars within 12 ly of Earth.
- Nearest: Proxima Centauri, 1.3 pc (4.2 ly), part of a triple system.
- 79\% of stars are low-mass red dwarfs.



## Astronomical units: Magnitudes

## Magnitude $m$, Brightness $f$

Ancient Greek system: first magnitude (and less) are the brightest stars. Sixth magnitude is the faintest the naked eye can see. (Magnitude $\uparrow$, brightness $\downarrow$ )

The eye is approximately logarithmic in its response: perceived brightness is proportional to the logarithm of flux $f$. To match up with the Greek scale,

$$
\begin{aligned}
& 5 \mathrm{mag} \Leftrightarrow \text { a factor of } 100 \text { in flux } \\
& 1 \mathrm{mag} \Leftrightarrow \text { a factor of } 100^{1 / 5} \approx 2.5 \text { in flux }
\end{aligned}
$$

So, for two stars with apparent magnitudes $m_{1}$ and $m_{2}$,

$$
m_{2}-m_{1}=2.5 \log \left(\frac{f_{1}}{f_{2}}\right)
$$

Note: $\log \equiv \log _{10}$ from now on.

## Magnitudes

Most people find magnitudes to be confusing at first, but they turn out to be very useful once you get used to the convention:

- Magnitudes are dimensionless quantities. They are related to the ratios of fluxes or distances.
- Another legacy from Greek astronomy: magnitudes run backwards from the intuitive sense of brightness.
- Brighter objects have smaller magnitudes. Fainter objects have larger magnitudes. Think of them like a "rank."
- Fluxes from a combination of objects measured all at once add up algebraically as usual. The magnitudes of combinations of objects do not (examples in a few slides).


## Magnitudes

The reference point of the apparent magnitude scale is a matter of arbitrary convention.

- As usual! Same as for the meter, second, kilogram, etc.

The practical definition of zero apparent magnitude is Vega ( $\alpha$ Lyrae) which has $m=0.0$ at a wide range of wavelengths (UV, optical, IR).

- This reference is hard to lose; it is the brightest star in the northern hemisphere.
- The apparent magnitude $m$ of a star from which we measure a flux is

$$
m=m-m_{\mathrm{Vega}}=2.5 \log \left(\frac{f_{\mathrm{Vega}}}{f}\right)
$$

- Fluxes from Vega have been measured with exquisite care and are tabulated.


## Magnitudes

Two other magnitude definitions are necessary:
Absolute magnitude is the apparent magnitude of a star if it were placed at a standard distance of 10 pc .
Bolometric magnitude is the magnitude calculated from the flux from all wavelengths, rather than from a small range of wavelengths.
In your homework, you will show that the absolute and apparent bolometric magnitudes $M$ and $m$ of a star are related by

$$
M=m-5 \log \left(\frac{r}{10 \mathrm{pc}}\right)=4.74-2.5 \log \left(\frac{L}{L_{\odot}}\right)
$$

## Magnitude calculations

The absolute bolometric magnitude of the Sun is 4.74 . What is its apparent bolometric magnitude?
The Sun's absolute magnitude is 4.74 if seen from 10 pc , but we see it from 1 AU . Thus,

$$
\begin{aligned}
m_{1} & =m_{2}+2.5 \log \left(\frac{f_{2}}{f_{1}}\right)=m_{2}+2.5 \log \left(\frac{L_{\odot}}{4 \pi r_{2}^{2}} \frac{4 \pi r_{1}^{2}}{L_{\odot}}\right) \\
& =4.74+2.5 \log \left(\frac{(1 \mathrm{AU})^{2}}{(10 \mathrm{pc})^{2}}\right)=4.74+5 \log \left(\frac{1 \mathrm{AU}}{10 \mathrm{pc}}\right) \\
m_{1} & =-26.83
\end{aligned}
$$

## Magnitude calculations

Two objects are measured to have fluxes $f$ and $f+\Delta f$, with $\Delta f \ll f$. What is the difference $\Delta m$ between their magnitudes?

$$
\begin{aligned}
\Delta m=m_{1}-m_{2} & =2.5 \log \left(\frac{f+\Delta f}{f}\right)=2.5 \log \left(1+\frac{\Delta f}{f}\right) \\
& =\frac{2.5}{\ln 10} \ln \left(1+\frac{\Delta f}{f}\right) \\
\Delta m & \approx 1.09 \frac{\Delta f}{f} \text { to first order in } \Delta f
\end{aligned}
$$

Thus, if two objects differ in flux by $1 \%$, then they differ in magnitude by 0.011 . It is very difficult to measure fluxes from astronomical objects more accurately than $1 \%$, so magnitudes are not typically known more accurately than $\pm 0.01$.

