

Physics 142 - November 15, 2005

Implications of Maxwell's Equations

$$\oint_S \vec{E} \cdot d\vec{A} = \frac{Q_{\text{encl}}}{\epsilon_0}$$

$$\vec{\nabla} \cdot \vec{E} = \rho / \epsilon_0$$

$$\oint_S \vec{B} \cdot d\vec{A} = 0$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\oint_C \vec{B} \cdot d\vec{\ell} = \mu_0 I_{\text{encl}} + \mu_0 \epsilon_0 \frac{d}{dt} \int_S \vec{E} \cdot d\vec{A}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\oint_C \vec{E} \cdot d\vec{\ell} = -\frac{d}{dt} \Phi_M = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{A}$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Integral form

Differential form

Maxwell's Equations

in region of no currents

$$\vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

mult by $\vec{\nabla} \times \dots$
Vector Calc Identities

$$\nabla^2 \vec{B} = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

Vector
wave
equation

3 scalar
wave
equations

$$\nabla^2 B_x = -\mu_0 \epsilon_0 \frac{\partial^2 B_x}{\partial t^2}$$

$$\nabla^2 B_y = -\mu_0 \epsilon_0 \frac{\partial^2 B_y}{\partial t^2}$$

$$\nabla^2 B_z = -\mu_0 \epsilon_0 \frac{\partial^2 B_z}{\partial t^2}$$

Laplacian

Laplacian of Scalar

$$\nabla^2 T = \vec{\nabla} \cdot (\vec{\nabla} T) = \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) \cdot \left(\hat{i} \frac{\partial T}{\partial x} + \hat{j} \frac{\partial T}{\partial y} + \hat{k} \frac{\partial T}{\partial z} \right)$$

$$\nabla^2 T = \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2}$$

Laplacian of a vector $\nabla^2 \vec{V}$

$$\nabla^2 \vec{V} = (\nabla^2 V_x) \hat{i} + (\nabla^2 V_y) \hat{j} + (\nabla^2 V_z) \hat{k}$$

$$\nabla^2 B_x = -\mu_0 \epsilon_0 \frac{\partial^2 B_x}{\partial t^2}$$

$$\rightarrow \frac{\partial^2 B_x}{\partial x^2} + \frac{\partial^2 B_x}{\partial y^2} + \frac{\partial^2 B_x}{\partial z^2} = -\mu_0 \epsilon_0 \frac{\partial^2 B_x}{\partial t^2}$$

Similarly $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = - \frac{\partial (\vec{\nabla} \times \vec{B})}{\partial t} = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\boxed{\nabla^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}}$$

wave eqn for
 $\vec{E}$ prop.

w/ velocity

$$\sqrt{\frac{1}{\mu_0 \epsilon_0}} = c$$

$\nearrow \rho=0$ $\nearrow \vec{j}=0$
region no charge no current

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\left[\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0 \right]$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\left[\frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} = 0 \right]$$

$$\vec{\nabla} \times \vec{E} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & E_z \end{vmatrix}$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$= \hat{i} \left(\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} \right) - \hat{j} \left(\frac{\partial E_z}{\partial x} - \frac{\partial E_x}{\partial z} \right) + \hat{k} \left(\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right)$$

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = -\frac{\partial B_x}{\partial t}$$

$\hat{i}$ part of $\vec{\nabla} \times \vec{E}$ eqn

$$-\left(\frac{\partial E_z}{\partial x} - \frac{\partial E_x}{\partial z} \right) = -\frac{\partial B_y}{\partial t}$$

$\hat{j}$ part

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = -\frac{\partial B_z}{\partial t}$$

$\hat{k}$ part

Similarly $\vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$ implies

$$\frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z} = \mu_0 \epsilon_0 \frac{\partial E_x}{\partial t}$$

$$-\left(\frac{\partial B_z}{\partial x} - \frac{\partial B_x}{\partial z} \right) = \mu_0 \epsilon_0 \frac{\partial E_y}{\partial t}$$

$$\frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} = \mu_0 \epsilon_0 \frac{\partial B_z}{\partial t}$$



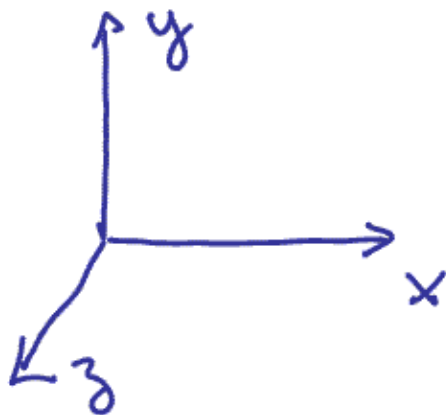
Source



large
Distance



observer



$$\vec{E} = \vec{E}(x, t)$$

Boundary
Conditions

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$$

$\swarrow 0$
 $\swarrow 0$

E_x const for all $x \rightarrow E_x = 0$

E_y E_z might not be zero

E is transverse to
direction of
propagation

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Impose

$\vec{E}$ is along y direction

$$\vec{E} = E(x, t) \hat{y}$$

"polarization"

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = -\frac{\partial B_x}{\partial t}$$

$$-\left(\frac{\partial E_z}{\partial x} - \frac{\partial E_x}{\partial z}\right) = -\frac{\partial B_y}{\partial t}$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = -\frac{\partial B_z}{\partial t}$$

B_x const

B_y const

IF E_y d'cs w/ x
 B_z d'cs

Time dependent
 B field

$\perp$ to E field

$\vec{E}, \vec{B}$ are TRANSVERSE, mutually $\perp$

$$E_y(x,t) = E_{oy} \cos(kx - \omega t + \phi)$$

$$\frac{2\pi}{\lambda}$$

$$\frac{2\pi}{T}$$

phase
(initial
condition)

$$B_z = - \int \frac{\partial E_y}{\partial x} dt$$

Assume
zero for
now

$$B_z = \int k E_{oy} \sin(kx - \omega t) dt$$

$$\phi = 0$$

$$B_z = \frac{k}{\omega} E_{oy} \cos(kx - \omega t)$$

$$\frac{k}{\omega} = \frac{\frac{2\pi}{\lambda}}{\frac{2\pi}{T}} = \frac{T}{\lambda} = \frac{1}{c} \quad \lambda/T = \text{wave speed}$$

$$B_z = \frac{1}{c} E_y$$