$\begin{array}{c} \mbox{Physics 237, Spring 2008} \\ \mbox{Homework } \#11 \\ \mbox{Due in class, Thursday April 24, 2008} \end{array}$

Homework is from Eisberg and Resnick (E&R) unless otherwise indicated. Please note that "questions" can be answered briefly; "problems" may be more involved. Each question usually counts 5 points; each problem counts 10 points.

- 1. E&R, chapter 8, problem 14.
- 2. E&R, chapter 8, problem 16.
- 3. E&R, chapter 8, problem 18.
- 4. E&R, chapter 9, problem 4.
- 5. E&R, chapter 9, problem 5.

HW#11 Solutions (Chapter 8)

| (4) Verify that the parities of the one-electron atom eigenfunctions 4300, 4310, 4320, and 4322 are determine by (-1)? |
|---------------------------------------------------------------------------------------------------------------------------------------------------------|
| Under a parity operation all rectangular coordinates change sign. Under spherical coordinates we have |
| $r \rightarrow r \qquad \varphi \rightarrow \pi + \varphi$ |
| So we must show that $\Psi_{nem}(r, T-0, T+\Psi) = (-1)^{\ell} \Psi_{nem}(r, C)$ from table 2-2(pg 243) |
| $\Psi_{300} = \frac{1}{81\sqrt{3}\pi} \left(\frac{2}{a_0}\right)^{3/2} \left(27 - 18\frac{2r}{a_0} + 2\frac{2^2r^2}{a_0^2}\right) e^{-2r/3a_0}$ |
| $P \Psi_{300} = \Psi_{300} = \Psi_{300} (-1)^{\ell=0} \sqrt{2}$ |
| $\Psi_{310} = \frac{\sqrt{2}}{81\sqrt{\pi}} \left(\frac{2}{a_0}\right)^{3/2} \left(6 - \frac{2r}{a_0}\right) \frac{2r}{a_0} e^{-2r/3a_0} \cos^2 \theta$ |
| $P\Psi_{310} = \Psi_{310} \frac{\cos(\pi - G)}{\cos \phi} = \Psi_{310} - \frac{\cos(-\phi)}{\cos(\phi)} = -\Psi_{310} = \Psi_{310}(-)$ |
| $\Psi_{320} = \frac{1}{81\sqrt{cT}} \left(\frac{2}{q_0}\right)^{3/2} \frac{2^2 \Gamma^2}{q_0} e^{-2r/3q_0} \left(3\cos^2(\sigma - 1)\right)$ |
| $P\Psi_{330} = \Psi_{330} \left(\frac{3 \cos^2(\pi - 0) - 1}{(3 \cos^2(0) - 1)} - \Psi_{330} - \frac{3 (-\cos^2)^2 - 1}{3 \cos^2(0 - 1)} \right)$ |
| $= \Psi_{320} = \Psi_{320}(-1)^{\ell=2} V$ |
| |

$$\begin{aligned} \Psi_{3aa} &= \frac{1}{16a\sqrt{\pi}} \left(\frac{2}{a_0} \right)^a \frac{2r^a}{a_a^a} e^{-2r/3a_0} \sin^a \sigma e^{\pm ai\varphi} \\ P \Psi_{3aa} &= \Psi_{3aa} \frac{\sin^a(\pi \cdot \sigma)}{\sin^a \sigma} \frac{e^{\pm ai(\pi + \varphi)}}{e^{\pm 2i\varphi}} \\ &= \Psi_{3aa} \left(-\sin(\cdot \sigma) \right)^a} \frac{e^{\pm 2i\varphi}}{e^{\pm 2i\varphi}} e^{\pm ai\pi} = \Psi_{3aa} \cos(\pm a\pi) \\ &= \Psi_{3aa} = \Psi_{3aa}(-1)^{a=a} \checkmark \end{aligned}$$

$$\begin{aligned} &(G) \text{ Show that the selection rule } \Delta \ell = \pm 1 \text{ is valid} \\ \text{for transitions } n = a \rightarrow n = 1 \text{ of the hydrogen atom.} \\ \text{The electric dipole matrix element has the form} \\ &P_{ei} = \left| \int \Psi_e e^{\frac{\pi}{2}} \Psi_e d^{\frac{\pi}{2}} \right| \\ n = a_i d^{\frac{\pi}{2}} - \frac{\pi}{2} n = a, \ R = 0, \ m_z = 0 \\ &R_{ei} \text{ only deal with anywher part of wave function} \\ &P_{ei} = \left| \int R_{ei}^{\pi}(r) \Theta_{ei}^{\pi}(\sigma) \Phi_{ei}^{\pi}(\varphi) = \overline{R}(r) \Theta_{ei}(\sigma) \Phi_{ei}(\varphi) d^{\frac{\pi}{2}} \right| \\ &e \int R_{ei}^{\pi}(r) R_{i}(r) n^{\frac{\pi}{2}} d^{\frac{\pi}{2}} = \operatorname{Constant} = C \\ &S_{o} we evaluate the rest of the integrals \\ &I = \int_{0}^{\pi} O_{ei}^{\pi}(\sigma) \Theta_{i}(\sigma) \sin\sigma d\sigma \int_{0}^{\pi} \overline{P}_{ei}^{\frac{\pi}{2}} \Psi_{ei}^{\frac{\pi}{2}} \Psi_{ei}^{\frac{\pi}{2}} \\ &= \int_{0}^{\pi} O_{ei}^{\pi}(\sigma) \Theta_{i}(\sigma) \sin\sigma d\sigma \int_{0}^{\pi} \overline{P}_{ei}^{\frac{\pi}{2}} \Psi_{i}^{\frac{\pi}{2}} \\ \end{bmatrix}$$

Ix = S.S. sino do - sino cosp dy $= \int_{0}^{\pi} \sin^{2} \sigma \, d\sigma \int_{0}^{2\pi} \cos \phi \, d\phi = 0 \, v$ $T_y = \int_0^{T} \sin^2 \phi \, d\phi \int_0^{2T} \sin \phi \, d\phi \, \phi = 0 v$ Dl =0 Iz = So sinofcoso St dq = Or for $n=2, l=1, m=0 \rightarrow n=1, l=0, m=0$ Ix= St cosor do dy sino cosy sino = Stososino do Stoos 4 dy = 0 $I_{y} = \int_{0}^{T} \cos \varphi \sin \varphi \, d\varphi = 0$ $I_z = \int_{a}^{T} \cos \sigma \sin \sigma \int_{a}^{M} d\phi = -\cos \sigma \left[\frac{T}{2} - 2T \right]_{a}^{T}$ = 4TT = O transition allowed

for
$$h=a, \lambda=1, M_{z=1} \rightarrow h=1, \lambda=0, M_{z}=0$$

 $I_{x} = \int_{0}^{\pi} \int_{0}^{\pi} \sin^{2} \theta e^{\pm i\theta} \sin \theta \cos \theta \, d\theta \, d\theta$
 $= \int_{0}^{\pi} \sin^{2} \theta \int_{0}^{2\pi} e^{\pm i\theta} \cos \theta$
 $= \left[-\frac{1}{3}\cos(\sin^{2}\theta+2)\right]_{0}^{\pi} \int_{0}^{2\pi} e^{\pm i\theta} \left(\frac{e^{i\theta}+e^{-i\theta}}{2}\right) d\theta$
 $= -\frac{1}{3}\left(-1-(t)\right)(2) \int_{0}^{2\pi} \frac{e^{i\theta}(1\pm 1)}{2} + \frac{e^{i\theta}(1\pm 1)}{2} d\theta$
 $= -\frac{2}{3}\left(2\pi - 0\right) = -\frac{4\pi}{3}$
 $= -\frac{2}{3}\left(2\pi - 0\right) = -\frac{4\pi}{3}$
 $= -\frac{4}{3}\left[\int_{0}^{2\pi} \frac{1+e^{i\theta}}{2} d\theta\right]$
 $= -\frac{4}{3}\left[\int_{0}^{2\pi} \frac{1+e^{i\theta}}{2} d\theta\right] = -\frac{4\pi}{3} \neq 0$ allowed
 $I_{y} = \int_{0}^{\pi} \sin^{2}\theta \, d\theta \int_{0}^{2\pi} e^{\pm i\theta} \sin \theta \, d\theta$
 $= \frac{4\pi\pi}{3} = \frac{4\pi\pi}{3} \neq 0$
 $I_{z} = \int_{0}^{\pi} \sin^{2}\theta \cos \theta \, d\theta \int_{0}^{2\pi} e^{\pm i\theta} d\theta = 0$

| (18) a) Calculate rate for n=1->n=0 transitions in a |
|-----------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|
| simple harmonic oscillator carrying charge e, |
| C= 103 J/ma |
| For a harmonic oscillator |
| Yo= Ao eu²/2 U= [[(em)1/4/1/5]X |
| Y= - LIC du=adx |
| V- att V m |
| Pio = St, extodx |
| $1 = S \#_{a} \#_{a} dx = A_{a}^{2} S e^{-u^{2}} du dx = A_{a}^{2} v_{\overline{m}}$ |
| $A_0 = \alpha^{1/2}$ |
| 1= S 4 4 dx = Ai Suie u = Ai [- Lue 10 - Soft)e du |
| $= A_1^2 \frac{\sqrt{\pi}}{4\alpha} A_1 = \frac{2\alpha'^2}{\pi'^4}$ |
| $P_{io} = \int_{0}^{\infty} \left(\frac{\alpha' a}{\pi' 4}\right) \left(\frac{2\alpha' a}{\pi' 4}\right) e^{-\frac{\alpha' a}{2}} e^{-\frac{\alpha' a}{2}} e^{-\frac{\alpha' a}{2}} dx$ |
| = dea Servidu |
| $= \frac{2e\alpha}{\sqrt{\pi}} \frac{\sqrt{\pi}}{4a^2} = \frac{2e}{\alpha}$ |
| $P_{10}^{2} = \frac{4e^{2}}{a^{2}}$ |

$$R = \frac{16\pi^{3}v^{3}}{3\epsilon_{o}hc^{3}} \frac{p^{2}}{r_{i}} (8-43)$$

$$= \frac{16\pi^{3}(\frac{1}{2\pi}\sqrt{\frac{e}{m}})^{3}}{3\epsilon_{o}hc^{3}} (\frac{4e^{2}h}{(e^{m})^{1/2}})$$

$$= 33\pi^{2}(e^{2}h) = \frac{16\pi}{3\epsilon_{o}hc^{3}} \frac{(10^{3}\sqrt{m^{2}})(1.6r^{-19}c)^{2}}{(3r^{6}m_{s}^{2})^{3}(2r^{10}c^{-19}c)}$$

$$= \frac{16\pi^{2}(e^{2}h)}{3\epsilon_{o}hc^{3}m^{2}} \frac{16\pi}{3(3r^{6}m_{s}^{2})^{3}(2r^{10}c^{-19}c)^{2}} (8.85r^{10}c^{2}n^{1})$$

$$= 4.5r^{10}\frac{5}{second} \text{ varies with molecular mass}$$
b) $\mathcal{I} = \frac{1}{2} = 2.3r^{10}\frac{6}{5}$

R

(4) Verify the eigenfunction of example 9-2 is antisymmetric with respect to a change of labels of two particles.

$$\begin{split} \Psi_{A} &= \int_{3} \left[\left(\Psi_{a}(1) \Psi_{\beta}(a) \Psi_{\beta}(3) + \Psi_{\beta}(1) \Psi_{\beta}(a) \Psi_{a}(3) + \Psi_{\beta}(1) \Psi_{\beta}(a) \Psi_{a}(3) + \Psi_{\beta}(1) \Psi_{\beta}(a) \Psi_{a}(3) + \Psi_{\beta}(1) \Psi_{\beta}(a) \Psi_{a}(3) + \Psi_{\beta}(1) \Psi_{\beta}(a) \Psi_{\alpha}(3) - \Psi_{\beta}(1) \Psi_{\beta}(a) \Psi_{\alpha}(3) - \Psi_{\beta}(1) \Psi_{\beta}(a) \Psi_{\beta}(3) - \Psi_{\beta}(1) \Psi_{\beta}(a) \Psi_{\beta}(3) - \Psi_{\beta}(3) - \Psi_{\beta}(3) \Psi_{\beta}(3) - \Psi_{\beta}(3) \Psi_{\beta}(3) - \Psi$$

All combination orders of a.B. & are present in positive and negative sign terms symmetrically for each indice. So we only have to show that the is antisymmetric under exchange of two indices.

 $\Psi_{A} = 7 \quad \Psi_{A}^{\prime} = \frac{1}{\sqrt{3}} \left[-\Psi_{\beta}(1) \Psi_{2}(2) \Psi_{\beta}(3) - \Psi_{\beta}(1) \Psi_{2}(2) \Psi_{4}(3) \right]$ + $\psi_{a}(1)$ $\psi_{a}(a)$ $\psi_{p}(3)$ + $\psi_{1}(1)$ $\psi_{p}(a)$ $\psi_{a}(3)$ + $\psi_{\beta}(1) \psi_{2}(2) \psi_{3}(3) + \psi_{2}(1) \psi_{3}(2) \psi_{\beta}(3) \int \sqrt{2}$

5 If two particles are in the same state then Ψ_{A} in 9-2 becomes for ex. $(a=\beta)$ $\Psi_{A} = \frac{1}{\sqrt{3}} \left[\Psi_{A}(1) \Psi_{A}(2) \Psi_{A}(3) + \Psi_{A}(1) \Psi_{A}(2) \Psi_{A}(3) + \Psi_{A}(1) \Psi_{A}(2) \Psi_{A}(3) - \Psi_{A}(1) \Psi_{A}(1) - \Psi_{A}(1) \Psi_{A}(1) - \Psi_{A}(1) \Psi_{A}(1) - \Psi$