Physics 237, Spring 2008 Homework #8Due in class, Thursday April 3, 2008

Homework is from Eisberg and Resnick (E&R) unless otherwise indicated. Please note that "questions" can be answered briefly; "problems" may be more involved. Each question usually counts 5 points; each problem counts 10 points.

- 1. E&R, chapter 7, question 9.
- 2. E&R, chapter 7, problem 5.
- 3. E&R, chapter 7, problem 7.
- 4. E&R, chapter 7, problem 8.
- 5. E&R, chapter 7, problem 10.

no: (day n) even though the same quarture nos corresponding to my. The water is said to be degenerate in the quartum no: n.

$$1 \rightarrow E_{\lambda} = -13.6 \qquad l = 0,1, \dots = -1$$

$$m = -1.6 \quad 0$$

The Energy eigenvalues are degenerate.

$$7_{100} = \frac{1}{\sqrt{\pi}} \left(\frac{7}{4_0} \right)^{3/2} e^{-\frac{7}{4} n/4_0}.$$

Lmg=0

The eyn (7-17) of the text becomes,

$$\frac{1}{10^2}\frac{d}{dx}\left(10^2\frac{dR}{dx}\right) + \frac{24}{4^2}\left[E - V(u)\right]R = 0$$

R= \$100

$$\frac{-\frac{1}{2}}{2m} \sqrt[3]{\frac{1}{100}} = \frac{1}{N^2} \frac{d}{dx} \left(N^2 \frac{d}{dx} \right) \frac{e^{-\lambda/e_0}}{\sqrt{\pi e_0^2}}$$

) from the diagram on pg 245 of the text (RAE), one can see that the probability is maximum for it = 100 where a is the Bohr radius and I is the abonic number.

$$\vec{n} = \int_{0}^{40} \int_{0}^{\pi} \int_{0}^{2\pi} dx \operatorname{Seodod} dy \quad \begin{cases} \frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} \left(\frac{\pi}{4_{0}}\right)^{2} n^{2} & e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \\ \times \sqrt{1 - \frac{\pi}{4_{0}}} \left(\frac{\pi}{4_{0}}\right)^{3} \left(\frac{\pi}{4_{0}}\right)^{2} n^{2} & e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \end{cases}$$

$$= \int_{0}^{40} \int_{0}^{\pi} \int_{0}^{2\pi} dx \operatorname{Seodod} dy \left(\frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} \left(\frac{\pi}{4_{0}}\right)^{2} n^{2} e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \right) dx \operatorname{Seodod} dy \left(\frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \right) dx \operatorname{Seodod} dy \left(\frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \right) dx \operatorname{Seodod} dy \left(\frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \right) dx \operatorname{Seodod} dy \left(\frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \right) dx \operatorname{Seodod} dy \left(\frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \right) dx \operatorname{Seodod} dy \left(\frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \right) dx \operatorname{Seodod} dy \operatorname{Seodod} dy \left(\frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \right) dx \operatorname{Seodod} dy \operatorname{Seodod} dy \operatorname{Seodod} dy \left(\frac{1}{64\pi} \left(\frac{\pi}{4_{0}}\right)^{3} e^{-\frac{\pi}{2} \ln \left| a_{0} \right|} \right) dx \operatorname{Seodod} dy \operatorname{Seodo$$

$$=\frac{1}{64\pi}\left(\frac{7}{a_0}\right)^5 \int_0^{\infty} e^{-\frac{7}{4}a/a_0} dx \int_0^{\infty} Sn^3 dx \int_0^{\infty} dx \int_0^$$

$$\frac{2\pi}{64\pi} \left(\frac{z}{a_0}\right)^5 \int_0^{\infty} e^{-\frac{z}{2}\omega a_0} u^5 dx \left(-\cos\theta \left|_0^{\pi} + \frac{G^3 a}{3}\right|_0^{\pi}\right)$$

$$= \frac{2\pi}{69\pi} \left(\frac{2}{0_6}\right)^5 \int_0^{\infty} e^{-\frac{2\pi}{3}} ds \left(2 - \frac{2}{3}\right)$$

$$\frac{7}{64\pi} \times \left(\frac{7}{60}\right)^{5} \times \frac{4}{3} \times \int_{0}^{\infty} e^{-\frac{7}{4} \ln \left(\frac{7}{60}\right)} dx$$

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=
$$\frac{2\pi}{64\pi} \times \left(\frac{7}{a_6}\right)^5 \times \frac{4}{3} \times \left(\frac{1}{7}\right)^6 \times \int_0^{\infty} e^{-\frac{1}{7}} y^5 dy$$

$$\frac{37}{328} \times \left(\frac{7}{4}\right) \times \frac{1}{2} \times \left(\frac{9}{4}\right) \times \frac$$

- (a) & (b) lies in the fact that the frist result which also coincides with the classical results for Bohr's Correspondence Principle is just the point at which the frohability is naminum. The second vesult is the average position weighted by the frohability. The first vesult was the classical result, the frohability is maximum at the classical result, but because of the waveneture of the particle there is no definite possibility of the fasticle at the same pt but only a probability. The mean value not coincident with the maximum value of frohability is because of the waveneture of the particle.
- 3) The expectation value of V for the fotential energy in the ground state is given by;

$$\bar{V} = \int_{0}^{\infty} \int_{0}^{2\pi} u^{2} dx \text{ anode } d\phi \left(\frac{1}{\pi} \left(\frac{7}{\sigma_{o}} \right)^{3} e^{-\frac{97}{4} \frac{7}{\sigma_{o}} \frac{7}{4} \frac{7}{\sigma_{o}} \frac{7}{4} \frac{7}{\sigma_{o}} \frac{7}{4} \frac{7}{\sigma_{o}} \frac{7}{\sigma_{o}} \frac{7}{4} \frac{7}{\sigma_{o}} \frac{7}{\sigma$$

$$= \frac{-Ze^2}{\pi\epsilon_0} \int_0^{\infty} x \left(\frac{Z}{a_0}\right)^3 e^{-2Zu/a_0} dx$$

$$= \frac{-Ze^2}{\pi\epsilon_0} \times \left(\frac{Z}{a_0}\right)^3 \times \int_0^{\infty} x e^{-2Zu/a_0} dx$$

$$= \frac{-Ze^2}{7\epsilon_0} \times \left(\frac{Z}{a_0}\right)^3 \times \left(\frac{a_0}{Z}\right)^2 \int_0^{\infty} y e^{-2y} dy$$

$$= -Ze^2 \cdot (Z) \cdot (Z)$$

$$= \frac{-2e^2}{\pi\epsilon_0} \times \left(\frac{z}{a_0}\right) \times \frac{1}{4}$$

$$= 2\left(-\frac{2^{2}e^{4}u}{(4\pi\epsilon_{0})^{2}h^{2}}\right) = 2i$$

$$v = \overline{x} + \overline{v}$$

$$-\frac{\sqrt{2}}{2} = K$$

$$\frac{1}{N^2} \frac{d}{dx} \left(N^2 \frac{dR}{dx} \right) + \frac{2u}{4^2} \left[E - V(N) \right] = \frac{2(1+i)R}{N^2}$$

We are attempting to find a robution for 1000

Multiplying throughout with 1,2 one has;

Inv for 12->0

 $\frac{d}{dx}\left(u^2\frac{dR}{dx}\right) + \frac{2\mu}{t^2}\left[E - V(u)\right]u^2 = \ell(\ell+1)R \text{ leads } dx$

 $\frac{l}{l}\left(u^{2}dR\right) = l(l+1)R. \quad \text{if } R(u) = usl \quad \text{one ha},$

 $\frac{d}{dr}\left(N^{2}lN^{l-1}\right) = \frac{d}{dr}\left(Lr^{l+1}\right) = l(l+1)N^{l}$ = l(l+1)R

di (u2-dR) = L(1+1)R.