

1. A hydrogen atom is subjected to the additional potential

$$V = A(\mathbf{L} \cdot \mathbf{S})^2$$

where A is a constant and \mathbf{L} and \mathbf{S} are the orbital and spin angular momenta.

- Give the exact expression for the total energy of the bound states of the system.
- Give the explicit results for the $n = 1$ and $n = 2$ levels of hydrogen. State the degeneracies of each distinct level.
- Compute the trace of V for the $n = 1$ and $n = 2$ levels.
- Show that the results of b) agree with the traces computed in c).

2. Consider the system with Hamiltonian

$$H = AL^2 + B(L_xL_y + L_yL_x)$$

where $B/A \ll 1$.

- Show that the ground state energy can be computed exactly. What is it?
- What are the $\ell = 1$ energy levels to lowest nonvanishing order in B ?
- Show that the sum of the eigenvalues must be independent of B for *all* ℓ values.
- Add the term

$$iC\epsilon_{rst}L_rL_sL_t$$

to H . How does the result of b) change with this replacement?

3. A hydrogen atom is subjected to the potential

$$V = \lambda xyz$$

- Show that to order λ the $n = 1$ and $n = 2$ levels are unaffected.
- For $n = 3$ write out the matrix of V . That is, indicate all zero elements of the matrix. Denote nonvanishing elements by constants a, b, c, \dots
- Determine all the $n = 3$ energy levels in terms of a, b, c, \dots
- Show that the matrix of V can in principle be expressed in terms of a single parameter (basically, a radial integral).

4. Determine the differential cross section for scattering of mass m particles by the square well potential

$$V = \begin{cases} V_0 & r < R \\ 0 & r > R \end{cases}$$

to lowest nonvanishing order in V_0 .

$$\begin{aligned}
 1) \quad a) \quad L \cdot S &= \frac{1}{2} (J^2 - L^2 - S^2) \\
 &= \frac{\hbar^2}{2} [j(j+1) - l(l+1) - 3/4] \\
 &\quad (j = l \pm \frac{1}{2}, \text{ except for } l=0) \\
 &= \frac{\hbar^2}{2} [(l \pm \frac{1}{2})(l \pm \frac{1}{2} + 1) - l(l+1) - 3/4] \\
 &= \frac{\hbar^2}{2} [\pm 2 \pm \frac{1}{2} - \frac{1}{2}] = \frac{\hbar^2}{2} (l, -l-1)
 \end{aligned}$$

$$E = -\frac{me^4}{2\hbar^2 n^2} + A \frac{\hbar^4}{4} \begin{cases} l^2 \\ (l+1)^2 \end{cases} \quad 0, \text{ if } l=0$$

$$b) \quad \frac{n=1}{n=2} \quad E = -\frac{me^4}{2\hbar^2} \quad \text{degen} = 2$$

$$E = -\frac{me^4}{8\hbar^2} + A \hbar^4 \frac{1}{4} \begin{pmatrix} 0 & l=0 \text{ (degen}=2) \\ 1 & j=3/2 \text{ (4)} \\ 4 & l=1, j=1/2 \text{ (2)} \end{pmatrix}$$

$$c) \quad \text{Tr } V = 0 \quad n=1$$

$$\text{Tr } V = A (\text{Tr } L_i L_j) (\text{Tr } S_i S_j)$$

$$= A \hbar^2 \frac{1}{2} \text{Tr } L^2 = \frac{A \hbar^4}{2} 2 \cdot 3 = 3 A \hbar^4$$

$$d) \quad \text{Tr } V = 0 \quad n=1$$

$$\text{Tr } V = \frac{A \hbar^4}{4} \{ 2 \times 0 + 4 \times 1 + 2 \times 4 \} \quad n=2$$

$$= 3 A \hbar^4$$

which checks

2 → for $l=0$

$$H|0\rangle = 0$$

as E of ground state is zero

$$b) L_x = \frac{L_+ + L_-}{2}, \quad L_y = \frac{L_+ - L_-}{2i}$$

$$H = AL^2 + B \left[\frac{(L_+ + L_-)(L_+ - L_-)}{4i} + \frac{(L_+ - L_-)(L_+ + L_-)}{4} \right]$$

$$= AL^2 + \frac{B}{4i} 2(L_+^2 - L_-^2)$$

matrix of H

$$\hbar^2 \begin{pmatrix} 2A & 0 & -iB \\ 0 & 2A & 0 \\ iB & 0 & 2A \end{pmatrix}$$

$$\Rightarrow (2A\hbar^2 - E)^3 - \hbar^2 B^2 (2A\hbar^2 - E) = 0$$

$$\Rightarrow E = 2A\hbar^2, \quad 2A\hbar^2 \pm B\hbar^2$$

$$c) H - AL^2 = -\frac{iB}{2} (L_+^2 - L_-^2)$$

in l, m representation $\text{Tr } L_{\pm}^2 = 0$

$$\Rightarrow \text{Tr } H = A\hbar^2 l(l+1)(2l+1)$$

independent of B

$$d) i C \epsilon_{rst} L_r L_s L_t = \frac{i}{2} C \epsilon_{rst} \hbar^2 [L_s, L_t]$$

$$= i \frac{\hbar^2}{2} C \epsilon_{rst} \epsilon_{stq} L_r L_q = -\frac{\hbar^2}{2} C L^2$$

$$E \rightarrow (A - \hbar C)\hbar^2, \quad (2A - \hbar C \pm B)\hbar^2$$

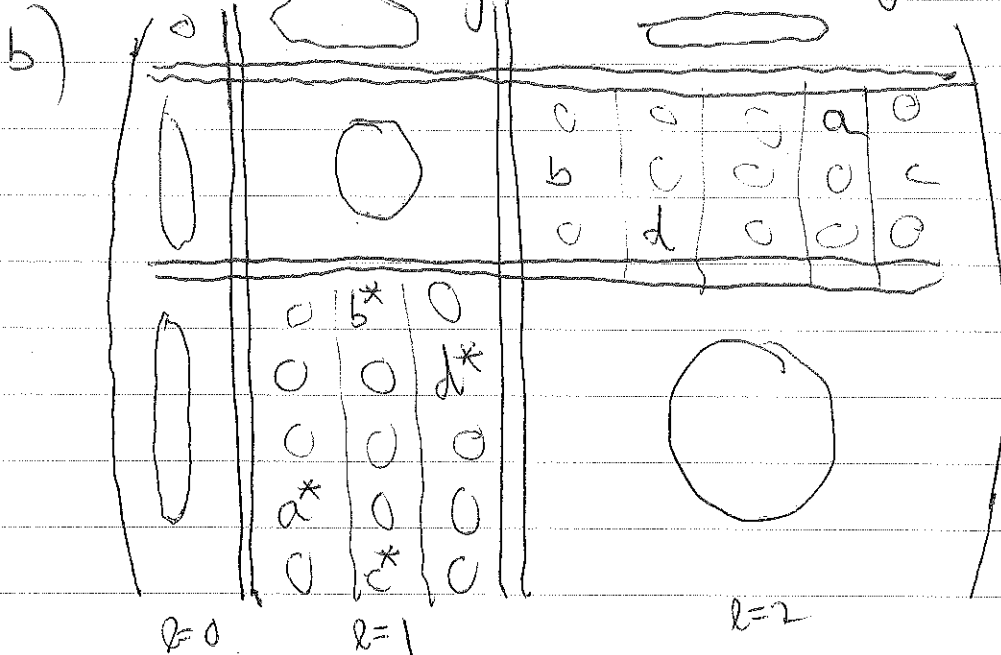
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$$a) \langle n=1 | V | n=1 \rangle = 0 \quad \text{by parity}$$

$$\langle n=2 | V | n=2 \rangle = 0 \quad \text{by parity}$$

$$\langle n=2, l=1 | V | n=2, l=0 \rangle = 0 \quad \text{because } V \propto r^2 [(x+iy)^2 - (x-iy)^2]$$

and m 's differ by two units to be nonzero. Not possible



$$c) E = -\frac{me^4}{18\hbar^2} \quad \text{Twice (} l=0 \text{ and } l=2, m=0)$$

for $l=1, m=0$ and $l=2, m=-1$

$$\begin{vmatrix} E & a \\ a^* & E \end{vmatrix} = 0 \quad \text{or} \quad E = -\frac{me^4}{18\hbar^2} \pm a$$

for $l=1, m=-1$ and $l=2, m=1$

$$\begin{vmatrix} E & d \\ d^* & E \end{vmatrix} = 0 \quad \text{or} \quad E = -\frac{me^4}{18\hbar^2} \pm d$$

(which tells us that $|a| = |d|$)

$$l=1, m=0; \quad l=2, m=\pm 2$$

$$\begin{vmatrix} E & b & c \\ b^* & E & 0 \\ c^* & 0 & E \end{vmatrix} = 0$$

$$E(E^2 - |b|^2 - |c|^2) = 0$$

$$E = -\frac{me^4}{18\hbar^2} \pm \sqrt{|b|^2 + |c|^2} \quad \text{and} \quad -\frac{me^4}{18\hbar^2}$$

so a total of three unchanged levels

d) $V \sim 2[(x+iy)^2 - (x-iy)^2]$
 is an element of a spherical tensor of rank 3; $T_{\pm 2}^{(3)}$
 $\langle l=1, m | V | l=2, m' \rangle = \langle l=1, m=3 || T_{\pm 2}^{(3)} || l=2, m=3 \rangle$
 \uparrow
 $T_{\pm 2}^{(3)} \quad \langle l=1, m=3, \pm 2; 2m' \rangle$

The Clebsch-Gordan coefficients are known so everything can be expressed in terms of a single parameter

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$$\begin{aligned} f(0) &= -\frac{2m}{4\pi\hbar^2} \int V(r) e^{i(k_i - k_f) \cdot r} d^3r \\ &= -\frac{mV_0}{2\pi\hbar^2} 2\pi \frac{1}{i|k_i - k_f|} \int_0^R r dr 2i \sin|k_i - k_f| r \\ &= -\frac{2mV_0}{\hbar^2 |k_i - k_f|} \frac{\partial}{\partial |k_i - k_f|} \left(\int_0^R dr (-\cos|k_i - k_f| r) \right) \\ &= -\frac{2mV_0}{\hbar^2 |k_i - k_f|} \frac{\partial}{\partial |k_i - k_f|} \frac{\sin|k_i - k_f| R}{|k_i - k_f|} \\ &= -\frac{2mV_0}{\hbar^2 |k_i - k_f|} \left[-\frac{\sin|k_i - k_f| R}{|k_i - k_f|^2} + R \frac{\cos|k_i - k_f| R}{|k_i - k_f|} \right] \end{aligned}$$

$$\frac{d\sigma}{d\Omega} = |F|^2 = \left(\frac{2mV_0}{\hbar^2 |k_i - k_f|} \right)^2 \left[\sin |k_i - k_f| R - R |k_i - k_f| \cos |k_i - k_f| R \right]^2$$

$$|k_i - k_f|^2 = 2k^2 (1 - \cos \Theta)$$

$$= 4k^2 \sin^2 \Theta/2$$

$$= \frac{8m^2 V_0^2}{\hbar^4} \sin^2 \Theta/2$$