

HOMEWORK 3

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$$1.1.1. \quad ds^2 = \frac{dx^2 + dy^2}{y^2}$$

$$L = \frac{1}{2} \frac{(x^2 + y^2)}{y^2}$$

$$\frac{\partial L}{\partial x} = \frac{\dot{x}}{y^2}, \quad \frac{\partial L}{\partial y} = \frac{\dot{y}}{y^2}$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) = \frac{\partial L}{\partial x} = 0$$

$$\frac{d}{dt} \left(\frac{\dot{x}}{y^2} \right) = 0 \quad \frac{y^2 \ddot{x} - 2\dot{x}y\dot{y}}{y^4} = 0$$

$$\ddot{x} - 2\frac{\dot{x}}{y}\dot{y} = 0 \quad \text{--- (1)}$$

$$\frac{d}{dt} \left(\frac{\dot{y}}{y^2} \right) = \frac{\partial L}{\partial y}$$

$$\frac{y^2 \ddot{y} - 2y\dot{y}^2}{y^4} = - \frac{(\dot{x}^2 + \dot{y}^2)}{y^3}$$

$$\begin{aligned} y^2 \ddot{y} - 2y\dot{y}^2 &= - \dot{x}^2 y - \dot{y}^2 y \\ \ddot{y} - \frac{\dot{y}^2}{y} + \frac{\dot{x}^2}{y} &= 0 \end{aligned} \quad \text{--- (2)}$$

From (1), (2) we find that -

$$\Gamma_{12}^1 = \Gamma_{21}^1 = -\frac{1}{y}, \quad \Gamma_{22}^2 = -\frac{1}{y}, \quad \Gamma_{11}^2 = +\frac{1}{y}.$$

Only derivatives with respect to the (2) component of these is non zero.

$$R_{22\alpha}^\rho = \partial_2 \Gamma_{2\alpha}^\rho - \partial_2 \Gamma_{2\alpha}^\rho + \Gamma_{2\alpha}^\rho \Gamma_{2\alpha}^\alpha - \Gamma_{2\alpha}^\rho \Gamma_{2\alpha}^\alpha$$

But all these combinations turn out to be zero as the terms cancel out.

$$R_{21\alpha}^\rho = \partial_2 \Gamma_{1\alpha}^\rho - \partial_1 \Gamma_{2\alpha}^\rho + \Gamma_{2\alpha}^\rho \Gamma_{1\alpha}^\alpha - \Gamma_{2\alpha}^\rho \Gamma_{1\alpha}^\alpha$$

$\therefore \Gamma_{12}^1$ is non zero,

$$R_{212}^1 = \partial_2 \Gamma_{12}^1 - \partial_1 \Gamma_{21}^1 + \Gamma_{2\alpha}^1 \Gamma_{12}^\alpha - \Gamma_{2\alpha}^1 \Gamma_{12}^\alpha$$

$$= \partial_y \left(-\frac{1}{y} \right) + 0 + \Gamma_{21}^1 \Gamma_{12}^1 - \Gamma_{21}^1 \Gamma_{12}^1$$

$$= -\frac{1}{y^2}$$

$$\text{Since } R_{121}^1 = -R_{122}^1 = -\frac{1}{y^2}.$$

This is the only non zero component.

$$R_{121}^2 = \partial_1 \Gamma_{21}^2 - \partial_2 \Gamma_{11}^2 + \Gamma_{12}^2 \Gamma_{21}^\alpha - \Gamma_{22}^2 \Gamma_{12}^\alpha$$

$$= 0 - \partial_y \left(\frac{1}{y} \right) + \Gamma_{12}^2 \Gamma_{21}^2 - \Gamma_{21}^2 \Gamma_{12}^2$$

$$R_{121}^2 = -R_{122}^1 = \frac{1}{y^2}$$

$$11.2. \quad \cancel{\frac{D^2 v^\mu}{dz^2}} = -R_{\nu\rho\sigma}^\mu v^\nu \dot{x}^\rho \dot{x}^\sigma$$

$$\frac{D^2 v^1}{dz^2} = -R_{122}^1 v^1 \dot{x}^2 = -\frac{1}{y^2} v^1 \dot{y}^2$$

$$\frac{D^2 v^2}{dz^2} = -R_{211}^2 v^2 \dot{x}^1 = -\frac{1}{y^2} v^2 \dot{x}^1$$

The sign of the curvature determines whether nearby geodesics will come nearer to each other or go far away.

$$1.2.1. \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

$$\partial_\mu F_{\nu\rho} + \partial_\nu F_{\rho\mu} + \partial_\rho F_{\mu\nu}$$

$$= \partial_\mu (\partial_\nu A_\rho - \partial_\rho A_\nu) + \partial_\nu (\partial_\rho A_\mu - \partial_\mu A_\rho) + \partial_\rho (\partial_\mu A_\nu - \partial_\nu A_\mu)$$

$$= (\partial_\mu \partial_\nu - \partial_\nu \partial_\mu) A_\rho + (\partial_\nu \partial_\rho - \partial_\rho \partial_\nu) A_\mu + (\partial_\rho \partial_\mu - \partial_\mu \partial_\rho) A_\nu$$

$$= 0$$

We know A's are continuous functions and the order of derivative does not matter. $\partial_\mu \partial_\gamma A_\rho = \partial_\gamma \partial_\mu A_\rho$.

$$1.2.2. D_\mu R^\alpha_{\gamma\rho\sigma} + D_\gamma R^\alpha_{\mu\nu\sigma} + D_\nu R^\alpha_{\mu\gamma\sigma} = 0$$

$$\begin{aligned} D_\mu R^\alpha_{\gamma\rho\sigma} &= \partial_\mu R^\alpha_{\gamma\rho\sigma} + \Gamma^\alpha_{\mu\beta} R^\beta_{\gamma\rho\sigma} - \Gamma^\beta_{\mu\nu} R^\alpha_{\beta\rho\sigma} - \Gamma^\beta_{\mu\rho} R^\alpha_{\gamma\beta\sigma} \\ &\quad - \Gamma^\beta_{\mu\sigma} R^\alpha_{\gamma\rho\beta} \end{aligned}$$

$$\begin{aligned} D_\gamma R^\alpha_{\mu\nu\sigma} &= \partial_\gamma R^\alpha_{\mu\nu\sigma} + \Gamma^\alpha_{\gamma\beta} R^\beta_{\mu\nu\sigma} - \Gamma^\beta_{\gamma\rho} R^\alpha_{\beta\mu\sigma} - \Gamma^\beta_{\gamma\mu} R^\alpha_{\rho\beta\sigma} \\ &\quad - \Gamma^\beta_{\gamma\sigma} R^\alpha_{\mu\nu\beta} \end{aligned}$$

$$\begin{aligned} D_\nu R^\alpha_{\mu\gamma\sigma} &= \partial_\nu R^\alpha_{\mu\gamma\sigma} + \Gamma^\alpha_{\nu\beta} R^\beta_{\mu\gamma\sigma} - \Gamma^\beta_{\nu\mu} R^\alpha_{\beta\gamma\sigma} - \Gamma^\beta_{\nu\gamma} R^\alpha_{\mu\beta\sigma} \\ &\quad - \Gamma^\beta_{\nu\sigma} R^\alpha_{\mu\gamma\beta} \end{aligned}$$

Adding all of them -

$$\begin{aligned} &= \partial_\mu R^\alpha_{\gamma\rho\sigma} + \partial_\gamma R^\alpha_{\mu\nu\sigma} + \partial_\nu R^\alpha_{\mu\gamma\sigma} + \Gamma^\alpha_{\mu\beta} R^\beta_{\gamma\rho\sigma} - \Gamma^\alpha_{\mu\beta} R^\beta_{\gamma\rho\sigma} - \Gamma^\beta_{\gamma\mu} R^\alpha_{\rho\beta\sigma} \\ &\quad - \Gamma^\beta_{\mu\beta} R^\alpha_{\gamma\rho\sigma} - \Gamma^\beta_{\mu\beta} R^\alpha_{\gamma\rho\sigma} - \Gamma^\beta_{\mu\beta} R^\alpha_{\gamma\rho\sigma} - \Gamma^\beta_{\mu\beta} R^\alpha_{\gamma\rho\sigma} \\ &\quad + \Gamma^\alpha_{\mu\beta} R^\beta_{\gamma\rho\sigma} + \Gamma^\alpha_{\mu\beta} R^\beta_{\gamma\rho\sigma} + \Gamma^\alpha_{\mu\beta} R^\beta_{\gamma\rho\sigma} - \Gamma^\alpha_{\mu\beta} R^\beta_{\gamma\rho\sigma} \\ &\quad - \Gamma^\beta_{\gamma\mu} R^\alpha_{\rho\beta\sigma} - \Gamma^\beta_{\gamma\mu} R^\alpha_{\rho\beta\sigma} \end{aligned}$$

$$\text{using } \Gamma^\alpha_{\mu\sigma} = \Gamma^\alpha_{\sigma\mu} \text{ and } R^\alpha_{\rho\beta\sigma} = -R^\alpha_{\beta\rho\sigma},$$

$$= \partial_\mu [\partial_\gamma \Gamma^\alpha_{\rho\sigma} / -\partial_\rho \Gamma^\alpha_{\gamma\sigma} + \Gamma^\alpha_{\gamma\theta} \Gamma^\theta_{\rho\sigma} - \Gamma^\alpha_{\rho\theta} \Gamma^\theta_{\gamma\sigma}]$$

We also consider an inertial frame now where acceleration is zero. Thus in the local inertial frame, ~~also~~ $\Gamma^\alpha_{\mu\sigma}$ vanish but not its derivative. Since the above identity is covariant, if it is true in some frame it is true in all frames. Thus we have -

$$\begin{aligned} &= \partial_\mu (\partial_\gamma \Gamma^\alpha_{\rho\sigma} - \partial_\rho \Gamma^\alpha_{\gamma\sigma}) + \partial_\gamma (\partial_\rho \Gamma^\alpha_{\mu\sigma} - \partial_\mu \Gamma^\alpha_{\rho\sigma}) \\ &\quad + \partial_\rho (\partial_\mu \Gamma^\alpha_{\gamma\sigma} - \partial_\gamma \Gamma^\alpha_{\mu\sigma}) \end{aligned}$$

$$= (\partial_\mu \partial_\nu - \partial_\nu \partial_\mu) \Gamma_{\rho\sigma}^\alpha + (\partial_\mu \partial_\rho - \partial_\rho \partial_\mu) \Gamma_{\nu\sigma}^\alpha \\ + (\partial_\nu \partial_\rho - \partial_\rho \partial_\nu) \Gamma_{\mu\sigma}^\alpha \\ = 0$$

The Christoffel symbols are associated with the metric which in turn is continuous.

1.2.3. The trace of the curvature tensor gives Ricci tensor and Ricci scalar. From Bianchi identity -

$$D_\mu R_{\nu\rho\sigma}^\alpha + D_\nu R_{\rho\mu\sigma}^\alpha + D_\rho R_{\mu\nu\sigma}^\alpha = 0$$

Multiplying by $g_{\delta\alpha}$, we note that -

$$D_\mu g_{\nu\rho\sigma} = \partial_\mu g_{\nu\rho\sigma} + \Gamma_{\nu\rho}^\theta g_{\theta\sigma} - \Gamma_{\mu\sigma}^\theta g_{\nu\theta} = 0 \\ \text{(covariant derivative of metric is zero)}$$

$$D_\mu R_{\nu\rho\sigma\delta} + D_\nu R_{\rho\mu\sigma\delta} + D_\rho R_{\mu\nu\sigma\delta} = 0$$

Multiplying by $g_{\mu\delta}$ -

$$D_\nu R_{\rho\sigma} + D_\rho R_{\nu\sigma} - D_\sigma R_{\nu\rho} = 0$$

Multiplying by $g_{\nu\delta}$ -

$$D_\nu R_{\rho\delta} + D_\rho R_{\nu\delta} - D_\delta R_{\nu\rho} = 0$$

$$D_\nu R_{\rho\delta} + D_\rho R_{\nu\delta} = D_\rho R$$

$$D_\nu R_{\rho\delta} = \frac{1}{2} D_\rho R = \frac{1}{2} D_\nu g_{\nu\rho} R$$

$$D_\nu (R_{\rho\delta} - \frac{1}{2} g_{\nu\rho} R) = 0$$

$$D_\mu (R_{\nu\delta} - \frac{1}{2} g_{\mu\nu} R) = 0$$

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} g^{\alpha\beta} R_{\alpha\beta}$$

$G_{\mu\nu}$ is the required tensor.

$$\begin{aligned}
 1.3.2. \quad S &= \frac{1}{4} \int F_{\mu\nu} \eta^{\mu\alpha} \eta^{\nu\beta} F_{\alpha\beta} dx \\
 &= \frac{1}{4} \int \eta^{\mu\alpha} \eta^{\nu\beta} F_{\mu\nu} F_{\alpha\beta} dx \\
 SS &= \frac{1}{4} \int \eta^{\mu\alpha} \eta^{\nu\beta} \delta F_{\mu\nu} F_{\alpha\beta} dx + \frac{1}{4} \int \eta^{\mu\alpha} \eta^{\nu\beta} F_{\mu\nu} \delta F_{\alpha\beta} dx \\
 \because \eta^{\mu\alpha} &= \eta^{\alpha\mu} \quad \text{interchanging } \mu, \alpha \text{ and } \nu, \beta \text{ in last term-} \\
 &= \frac{1}{2} \int \eta^{\mu\alpha} \eta^{\nu\beta} \delta F_{\mu\nu} F_{\alpha\beta} dx \\
 \delta F_{\mu\nu} &= \partial_\mu \delta A_\nu - \partial_\nu \delta A_\mu + \delta A_\mu A_\nu + A_\mu \delta A_\nu - \delta A_\nu A_\mu \\
 &\quad - A_\nu \delta A_\mu \\
 &= \partial_\mu (\partial_\rho A_\nu \delta x^\rho) - \partial_\nu (\partial_\rho A_\mu \delta x^\rho) + (\partial_\rho A_\mu) A_\nu \delta x^\rho \\
 &\quad + A_\mu (\partial_\rho A_\nu) \delta x^\rho - (\partial_\rho A_\nu) A_\mu \delta x^\rho - A_\nu (\partial_\rho A_\mu) \delta x^\rho \\
 &= (\partial_\mu \partial_\rho A_\nu) \delta x^\rho + (\partial_\rho A_\nu) (\partial_\mu \delta x^\rho) - (\partial_\nu \partial_\rho A_\mu) \delta x^\rho \\
 &\quad - (\partial_\rho A_\mu) (\partial_\nu \delta x^\rho) + [(\partial_\rho A_\mu) A_\nu - A_\mu \partial_\rho A_\nu - (\partial_\rho A_\nu) A_\mu - A_\nu (\partial_\rho A_\mu)] \delta x^\rho
 \end{aligned}$$

Let us consider integral of this type-

$$\int f(x) \partial_\nu \delta x^\rho dx$$

integration computed with respect to τ component-

$$= \left[f(x) \delta x^\rho \right] - \int \partial_\nu f(x) \delta x^\rho dx$$

Using the above facts we have-

$$\begin{aligned}
 SS &= \frac{1}{2} \int \eta^{\mu\alpha} \eta^{\nu\beta} F_{\alpha\beta} ((\partial_\mu \partial_\rho A_\nu - \partial_\nu \partial_\rho A_\mu + (\partial_\rho A_\mu) A_\nu + A_\mu \partial_\rho A_\nu \\
 &\quad - (\partial_\rho A_\nu) A_\mu - A_\nu (\partial_\rho A_\mu)) \delta x^\rho) dx \\
 &\quad + \frac{1}{2} \int \eta^{\mu\alpha} \eta^{\nu\beta} F_{\alpha\beta} ((\partial_\mu \partial_\rho A_\nu - \cancel{\partial_\nu \partial_\rho A_\mu} - \partial_\nu \partial_\rho A_\mu) \delta x^\rho) dx
 \end{aligned}$$

Since this result $SS = 0$ for any δx^ρ , integrand is zero.

$$\eta^{\alpha\beta} F_{\alpha\beta} (\partial_\mu \partial_\rho A_\gamma + (\partial_\rho A_\mu) A_\gamma + A_\mu (\partial_\rho A_\gamma) - (\partial_\rho A_\gamma) A_\mu - A_\gamma (\partial_\rho A_\mu) - \partial_\gamma \partial_\rho A_\mu + \partial_\mu \partial_\rho A_\gamma - \partial_\gamma \partial_\rho A_\mu) = 0$$

$$F^{M\gamma} (2\partial_\mu \partial_\rho A_\gamma - 2\partial_\gamma \partial_\rho A_\mu + (\partial_\rho A_\mu) A_\gamma - (\partial_\rho A_\gamma) A_\mu + A_\mu \partial_\rho A_\gamma - A_\gamma \partial_\rho A_\mu) = 0$$

$$F^{M\gamma} (\partial_\rho A_\gamma) A_\mu = F^{\gamma\mu} (\partial_\rho A_\mu) A_\gamma = - F^{\mu\gamma} (\partial_\rho A_\mu) A_\gamma$$

Thus we get -

$$F^{K\gamma} (\partial_\mu \partial_\rho A_\gamma - \partial_\gamma \partial_\rho A_\mu + (\partial_\rho A_\mu) A_\gamma + A_\mu \partial_\rho A_\gamma) = 0$$

$$F^{K\gamma} (2\partial_\mu \partial_\rho A_\gamma - \partial_\gamma \partial_\rho A_\mu + \partial_\rho (A_\mu A_\gamma)) = 0$$

$$F^{K\gamma} (\partial_\mu \partial_\rho A_\gamma + \frac{1}{2} \partial_\rho (A_\mu A_\gamma)) = 0$$

$$F_{M\gamma} = \partial_\mu A_\gamma - \partial_\gamma A_\mu + [A_\mu, A_\gamma]$$

Let us define covariant derivative -

$$D_\rho = \partial_\rho + A_\rho$$

$$\begin{aligned} [D_\mu, D_\gamma] &= [\partial_\mu + A_\mu, \partial_\gamma + A_\gamma] \\ &= \partial_\mu \partial_\gamma - \partial_\gamma \partial_\mu + [A_\mu, A_\gamma] + [A_\mu, \partial_\gamma] \\ &\quad + [A_\mu, A_\gamma] \end{aligned}$$

$$[A_\mu, \partial_\gamma] \Psi$$

$$= A_\mu \partial_\gamma \Psi - A_\mu (\partial_\gamma \Psi) - (\partial_\gamma A_\mu) \Psi = -\partial_\gamma A_\mu \Psi$$

$$\begin{aligned} [D_\mu, D_\gamma] &= + [A_\mu, A_\gamma] + (-\partial_\gamma A_\mu) + (-\partial_\mu A_\gamma) \\ &= + (\partial_\mu A_\gamma - \partial_\gamma A_\mu + [A_\mu, A_\gamma]) \\ &= F_{\mu\gamma} \end{aligned}$$

The analogue of Bianchi identity is -

$$D_\mu F_{\gamma p} + D_\gamma F_{p\mu} + D_p F_{\mu\gamma} = 0$$

To arrive at this we use Jacobi identity -

$$[D_\mu, [D_\gamma, D_p]] + [D_\gamma, [D_p, D_\mu]] + [D_p, [D_\mu, D_\gamma]] = 0$$

$$[D_\mu, F_{\gamma p}] + [D_\gamma, F_{p\mu}] + [D_p, F_{\mu\gamma}] = 0$$

We now see its action on a wave function -

$$[D_\mu, F_{\gamma p}] \psi + [D_\gamma, F_{p\mu}] \psi + [D_p, F_{\mu\gamma}] \psi = 0$$

$$D_\mu (F_{\gamma p} \psi) - F_{\gamma p} D_\mu \psi + D_\gamma (F_{p\mu} \psi) - F_{p\mu} D_\gamma \psi + D_p (F_{\mu\gamma} \psi) - F_{\mu\gamma} D_p \psi = 0$$

$$(D_\mu F_{\gamma p}) \psi + F_{\gamma p} D_\mu \psi - F_{\gamma p} D_\mu \psi + (D_\gamma F_{p\mu}) \psi + F_{p\mu} D_\gamma \psi - F_{p\mu} D_\gamma \psi + (D_p F_{\mu\gamma}) \psi + F_{\mu\gamma} D_p \psi - F_{\mu\gamma} D_p \psi = 0$$

$$(D_\mu F_{\gamma p} + D_\gamma F_{p\mu} + D_p F_{\mu\gamma}) \psi = 0$$

\therefore This is true for any ψ , we get -

$$D_\mu F_{\gamma p} + D_\gamma F_{p\mu} + D_p F_{\mu\gamma} = 0$$

