

Solving ablative RT Problem in AstroBEAR

1. Euler Fluid Equations with diffusion in c.g.s.

The Euler equations can be written as:

$$\begin{aligned}\rho_t + (\rho u)_x + (\rho v)_y &= 0 \\ (\rho u)_t + (\rho u^2 + p)_x + (\rho u v)_y &= 0 \\ (\rho v)_t + (\rho u v)_x + (\rho v^2 + p)_y &= -\rho g \\ \left(\frac{\rho v_t^2}{2} + \frac{p}{\gamma-1}\right)_t + \nabla \cdot \left[v_t \left(\frac{\rho v_t^2}{2} + \frac{\gamma p}{\gamma-1}\right) \right] &= \nabla \cdot (\kappa T^n \nabla T)\end{aligned}$$

From the last equation, we have:

$$\left(\frac{\rho v_t^2}{2}\right)_t + \left(\frac{p}{\gamma-1}\right)_t + \left(\frac{\rho v_t^2}{2} + \frac{\gamma p}{\gamma-1}\right) \nabla \cdot v_t + v_t \nabla \cdot \left(\frac{\rho v_t^2}{2} + \frac{\gamma p}{\gamma-1}\right) = \nabla \cdot (\kappa T^n \nabla T)$$

The second term can be written as:

$$\frac{p}{\gamma-1} = \frac{(1+Z)\rho k_B T}{(\gamma-1)m_i}$$

If we write $c_v = \frac{(1+Z)}{(\gamma-1)m_i}$, and measure T in terms of energy, then we have the EOS:

$$\frac{p}{\gamma-1} = c_v \rho T$$

then the energy equation can be written as:

$$\left(\frac{\rho v_t^2}{2}\right)_t + (c_v \rho T)_t + \left(\frac{\rho v_t^2}{2} + \gamma c_v \rho T\right) \nabla \cdot v_t + v_t \cdot \nabla \left(\frac{\rho v_t^2}{2} + \gamma c_v \rho T\right) = \nabla \cdot (\kappa T^n \nabla T)$$

Simplify this equation by using the momentum and mass conservations, we get the energy equation in the form of:

$$\rho c_v (\partial_t T + v \cdot \nabla T) = -p \nabla \cdot v + \nabla \cdot (\kappa T^n \nabla T)$$

Notice that in this form, all quantities are in SI with T measured in Joule. So it's actually:

$$\rho^S c_v^S (\partial_t k^S T^S + v^S \cdot \nabla k^S T^S) = -p^S \nabla^S \cdot v^S + \nabla^S \cdot (\kappa^S (k^S)^n (T^S)^n \nabla^S k^S T^S)$$

To convert to the c.g.s units, we have the following conversion factors:

$$\begin{aligned} \rho^S &= 1000 \rho^G & p^S &= 1/10 p^G \\ v^S &= 1/100 v^G & \nabla^S &= 100 \nabla^G \\ c_v^S &= 1000 c_v^G & k^S &= 10^{-7} k^G \end{aligned}$$

Substitute, we find that:

$$\rho^G c_v^G (\partial_t k^G T^G + v^G \cdot \nabla k^G T^G) = -p^G \nabla^G \cdot v^G + \nabla^G \cdot (10^{(-7n-2)} \kappa^S (k^G)^n (T^G)^n \nabla^G k^G T^G)$$

If we define $\kappa^G = 10^{(-7n-2)} \kappa^S$, then:

$$\rho^G c_v^G (\partial_t k^G T^G + v^G \cdot \nabla k^G T^G) = -p^G \nabla^G \cdot v^G + \nabla^G \cdot (\kappa^G (k^G)^n (T^G)^n \nabla^G k^G T^G)$$

Using operator splitting method, we need to solve:

$$\rho c_v^G (\partial_t k^G T) = \nabla^G \cdot (\kappa^G (k^G)^n T^n \nabla k^G T)$$

Divide by k and cv in c.g.s., we get:

$$\rho (\partial_t T) = \nabla \cdot \left(\frac{\kappa^G (k^G)^n}{c_v^G} T^n \nabla T \right)$$

define:

$$\kappa_0 = \frac{\kappa^G (k^G)^n}{c_v^G} = \frac{10^{(-7n+1)} (k^G)^n \kappa^S}{c_v^S}$$

then we have the equation ready to be scaled:

$$\rho (\partial_t T) = \nabla \cdot (\kappa_0 T^n \nabla T)$$

For the flux equation, we have:

$$q^S = \kappa^S (k^S)^n T^n \nabla^S T$$

we need a form like:

$$q_0 = \kappa_0 T^n \nabla^G T$$

we can divide these two equations and use the relations between SI units and Gauss units. The result is:

$$q_0 = \frac{q^S}{10 c_v^S k^S}$$

The scales can be obtained by writing equation $\rho (\partial_t T) = \nabla \cdot (\kappa_0 T^n \nabla T)$ in the forms of scalings:

$$\frac{r_{scale} T_{scale}}{t_{scale}} = \frac{\kappa_{scale} T_{scale}^{(n+1)}}{l_{scale}^2}$$

simplify this form, we get:

$$\kappa_{scale} = \frac{l_{scale} r_{scale}^{1/2} p_{scale}^{1/2}}{T_{scale}^n}$$

also, we have:

$$q_{scale} = \frac{\kappa_{scale} T_{scale}^{n+1}}{l_{scale}} = r_{scale}^{1/2} p_{scale}^{1/2} T_{scale}$$

2.Quasi-equilibrium

The quasi-equilibrium in RT problem is defined by the balance of the thermal pressure, Ram pressure and the gravity force. In equilibrium, the Euler equations are:

$$\frac{\partial \rho v}{\partial y} = 0$$

$$\rho v \frac{\partial v}{\partial y} = -\frac{\partial p}{\partial y} - \rho g$$

$$\rho v C_v \frac{\partial T}{\partial y} = -p \frac{\partial v}{\partial y} + \frac{\partial}{\partial y} (\kappa T^n \frac{\partial T}{\partial y})$$

from the first equation we know that $\rho v = \rho_0 v_0$, which is just the outflow at the bottom. If we define the heat flux as:

$$q = \kappa T^n \frac{\partial T}{\partial y}$$

we can then write the equations as:

$$\frac{\partial \rho}{\partial y} = \frac{\frac{\rho R q}{\kappa T^n} + \rho g}{\frac{\rho_0^2 v_0^2}{\rho^2} - RT}$$

$$\frac{\partial T}{\partial y} = \frac{q}{\kappa T^n}$$

$$\frac{\partial q}{\partial y} = \frac{\rho_0 v_0 C_v q}{\kappa T^n} - \rho_0 v_0 R T \frac{\frac{R q}{\kappa T^n} + g}{\frac{\rho_0^2 v_0^2}{\rho^2} - RT}$$

with ρ , v , q known on the two boundaries.

Using parameters (in SI units, Temperature is measured in Joule):

$$R = 4.7904 \times 10^{26}$$

$$C_v = 7.186 \times 10^{26}$$

$$\kappa = 3.734 \times 10^{69}$$

$$g = 1.0 \times 10^{14}$$

$$q = -5.876 \times 10^{18}$$

$$\rho_0 = 68.1622919147237$$

$$v_0 = -272144.604867564$$

The solutions to this equation set are posted below:



