

Capacitance & Capacitors

An isolated conductor with charge Q has a ~~pot~~ potential ϕ_0 with respect to infinity.

For example: a conducting sphere of radius a charge Q has:

$$\begin{aligned} \text{Electric field } \vec{E}(\vec{r}) &= \begin{cases} \frac{Q}{r^2} \hat{r} & r > a \\ 0 & r < a \end{cases} \\ \text{potential } \phi(\vec{r}) &= \begin{cases} \frac{Q}{r} & r > a \\ \frac{Q}{a} & r < a \end{cases} \end{aligned}$$

So the sphere is at potential $\phi_0 = \phi(a) = \frac{Q}{a}$

We see that $\phi_0 \propto Q$. This is in general the case no matter what the shape of the conductor.

One defines the capacitance of the conductor by

$$C = \frac{Q}{\phi_0}$$

capacitance of conductor with net charge Q

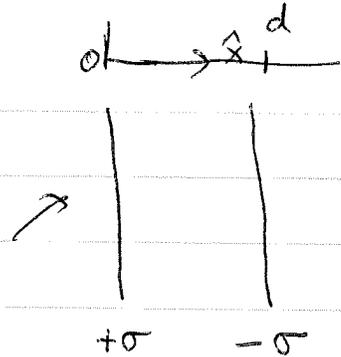
Since $Q \propto \phi_0$, C is independent of the total charge on the conductor. It is determined only by the shape of the conductor

units of $C = \frac{\text{esu}}{\text{stat volt}}$

$$\text{stat volt} = \frac{\text{esu}}{\text{cm}}$$

$$\Rightarrow C = \frac{\text{esu}}{\text{esu/cm}} = \text{cm}$$

area A
where
 $\sqrt{A} \gg d$



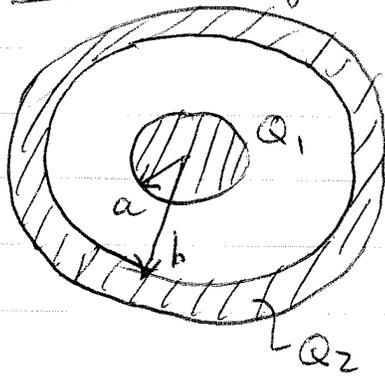
$\vec{E} = 0 \quad x < 0, x > d$
 $\vec{E} = 4\pi\sigma \hat{x} \quad 0 < x < d$

potential $\phi(d) - \phi(0) = -\int_0^d \vec{E} \cdot d\vec{x}$
 $= -4\pi\sigma d$

Capacitance is

$$C = \frac{\sigma A}{\phi(0) - \phi(d)} = \frac{\sigma A}{4\pi\sigma d} = \frac{A}{4\pi d}$$

spherical capacitor



conductor 1: sphere radius a
 conductor 2: spherical shell
 of inner radius b , outer
 radius c ,

Q_1 on conductor 1

Q_2 on conductor 2

The charge on the inner surface of conductor 2 must be $-Q_1$. This is because $\vec{E} = 0$ inside conductor 2 \Rightarrow charge enclosed by a surface of radius $b < r < c$ must contain zero charge.

Charge on outside surface of conductor 2 must therefore be $Q_2 + Q_1$. (If $Q_2 = -Q_1$ then charge on outer wall is zero).

The capacitance of this configuration is defined with respect to conductor 1 and inner surface of conductor 2.

$$C = \frac{Q_1}{\phi(a) - \phi(b)}$$

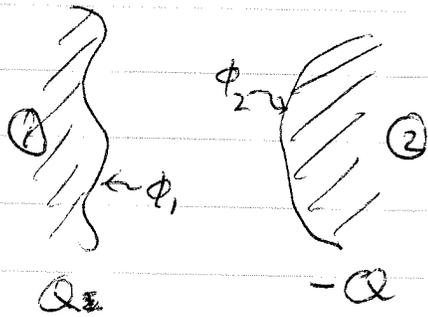
To get ϕ we first solve for $\vec{E}(\vec{r})$ for $a < r < b$

Gauss Law $\oint \vec{E} \cdot d\vec{a} = 4\pi r^2 E(r) = 4\pi Q_1$
↑ surface radii r $\vec{E}(\vec{r}) = \frac{Q_1}{r^2} \hat{r}$

$$\begin{aligned} \Rightarrow \phi(a) - \phi(b) &= - \int_b^a dr E(r) \\ &= \int_a^b dr E(r) \\ &= \int_a^b dr \frac{Q_1}{r^2} \\ &= Q_1 \left[\frac{1}{a} - \frac{1}{b} \right] \\ &= Q_1 \left(\frac{b-a}{ab} \right) \end{aligned}$$

$$C = \frac{Q_1}{Q_1 \left(\frac{b-a}{ab} \right)} = \frac{ab}{b-a}$$

Energy stored in a capacitor



Suppose we take a charge dQ from conductor (2) and move it to conductor (1). The work done to move the charge is

$$dW = dQ(\phi_1 - \phi_2)$$

$$\text{but } \phi_1 - \phi_2 = \frac{Q}{C}$$

$$dW = \frac{dQ Q}{C} \quad \text{where } C \text{ is indep of the charge on capacitor } Q.$$

What is the work done, starting from an uncharged capacitor ($Q=0$) to charge it up to charge Q ?

$$W = \int dW = \int_0^Q \frac{dQ' Q'}{C} = \frac{1}{2} \frac{Q^2}{C}$$

This is the electrostatic energy stored in the capacitor

$$W = \frac{1}{2} \frac{Q^2}{C}$$

$$\text{Since } Q = C \phi_{12}$$

$$W = \frac{1}{2} \left(\frac{C \phi_{12}}{C} \right)^2 = \frac{1}{2} C \phi_{12}^2$$

$$\phi_{12} = \phi_1 - \phi_2$$

↑ potential difference from one conductor to the other

For a parallel plate capacitor $C = \frac{A}{4\pi d}$

$$U = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} \frac{(\sigma A)^2}{A} 4\pi d = 2\pi \sigma^2 A d$$

Compare to

$$U = \frac{1}{8\pi} \int dV |\vec{E}|^2 \quad \text{between plates } |\vec{E}| = 4\pi\sigma$$

$$= \frac{1}{8\pi} (4\pi\sigma)^2 \cdot A d = 2\pi\sigma^2 A d$$

or

$$U = \frac{1}{2} \int dV \rho \phi = \frac{1}{2} \sigma A (\phi_1 - \phi_2)$$

$$\phi_1 - \phi_2 = - \int_2^1 \vec{E} \cdot d\vec{s} = \int_1^2 E \cdot ds = \frac{Q}{4\pi\sigma d}$$

$$U = \frac{1}{2} \sigma A (4\pi\sigma d) = 2\pi\sigma^2 A d$$

Same result holds for an isolated conductor -
For a ^{conducting} sphere of radius a , $C = a$

$$\Rightarrow U = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} \frac{Q^2}{a} \quad \text{energy stored in a charged conducting sphere}$$

$$\text{Compare to } U = \frac{1}{2} \int dV \rho \phi = \frac{1}{2} 4\pi a^2 \sigma \phi(a)$$

$$= 2\pi a^2 \sigma \left(\frac{Q}{a}\right) = 2\pi a^2 \sigma \frac{Q}{a}$$

$$= \frac{1}{2} (4\pi a^2 \sigma) \frac{Q}{a} = \frac{1}{2} \frac{Q^2}{a}$$

$$\text{Then } - \int_{\vec{r}_1}^{\vec{r}_2} \vec{F} \cdot d\vec{s} = \int_{\vec{r}_1}^{\vec{r}_2} \vec{\nabla} U \cdot d\vec{s} = U(\vec{r}_2) - U(\vec{r}_1)$$

line integral of a gradient is always independent of the path.

For the electrostatic electric field \vec{E} we have

$$\vec{E} = \frac{\vec{F}}{q} \quad \text{and} \quad - \int_{\vec{r}_1}^{\vec{r}_2} q \vec{E} \cdot d\vec{s} \text{ is independent of the path}$$

\Rightarrow electrostatic force is conservative

\rightarrow there exists potential energy U

$$U(\vec{r}_2) - U(\vec{r}_1) = - \int_{\vec{r}_1}^{\vec{r}_2} q \vec{E} \cdot d\vec{s}$$

$$\text{or } \frac{U(\vec{r}_2)}{q} - \frac{U(\vec{r}_1)}{q} = - \int_{\vec{r}_1}^{\vec{r}_2} \vec{E} \cdot d\vec{s}$$

$\frac{U}{q}$ is therefore just the electrostatic potential ϕ

$$\phi = \frac{U}{q} \quad \text{and} \quad \text{just like } \vec{F} = -\vec{\nabla} U \\ \text{we have } \vec{E} = -\vec{\nabla} \phi$$

ϕ is the potential energy per unit charge of a charged particle in an electrostatic field due to other charges.

Electric Currents

An electric current is due to charges in motion
Units of current is charge/time
in CGS it is esu/s. In MKS it is coul/sec = "amp"

For current flowing down a wire, the current is steady if the same amount of charge flows past each point in the wire in the same amount of time. Every charge that flows out of a particular segment of the wire is replaced by another charge flowing into that segment of the wire.



For a wire with a steady current, the moving charges have a constant average speed that does not change with time.

Consider a more general situation of charges moving in space. The current I flowing through some ^{small} area $\Delta \vec{a}$ is the flux of charge through $\Delta \vec{a}$

$$I = nq\vec{v} \cdot \Delta \vec{a}$$

↑ density of charges

↑ charge per particle

↑ average ~~speed~~ velocity of particles

See notes when we introduced concept of flux.

For many different types of particles i
 particle of type i has charge q_i and velocity \vec{v}_i
 and density n_i

$$I = \sum_i n_i q_i \vec{v}_i \cdot \Delta \vec{a}$$

If just as many $+q$ as $-q$ pass through $\Delta \vec{a}$
 in same time, there is no net current flow
 (Ex - water flowing in pipe has no electric current.
 just as many electrons with $-e$ pass through
 as do protons with $+e$ - so no net charge
 flowing)

Look at one type of charged particle, say electrons.
 If these have different velocities \vec{v}_k , then

$$I = \sum_k (-e) n_k \vec{v}_k \cdot \Delta \vec{a}$$

$$I = -e n \langle \vec{v} \rangle \cdot \Delta \vec{a}$$

↙ average over all electron velocities.

↑ density of all electrons

- 1) For current of negatively charged particles,
 current flows in opposite direction to
 the particles average velocity