# Physics 403 Common Probability Distributions

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- Gaussian Distribution
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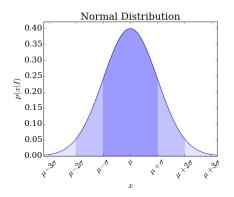
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Last Time

The 68-95-99 Rule

In physics we tend to express rare events in terms of the tails of the Gaussian  $\mathsf{PDF}$ 

$$p(x|I) = rac{1}{\sqrt{2\pi\sigma}} \exp\left\{-rac{1}{2}\left(rac{x-\mu}{\sigma}
ight)^2
ight\}$$



The "68-95-99" quantile rule:

- 68.27% of the data are within 1σ of the mean.
- 95.45% of the data are within 2σ of the mean.
- 99.73% of the data are within  $3\sigma$  of the mean.

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Probability Distributions

- Binomial Distribution
- Poisson Distribution
- Gaussian Distribution
- Uniform Distribution
- $\chi^2$  Distribution
- Other Distributions

### **Binomial Distribution**

- ▶ Bernoulli trials i.e, binary measurements which result in "success" with probability p and "failure" with probability 1 p are described by the binomial distribution.
- ▶ In *n* trials, like a coin toss, the probability of *m* "heads" is

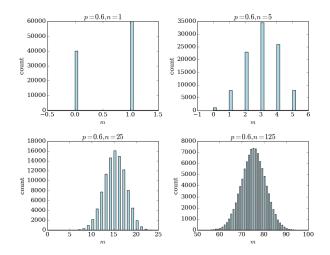
$$p^m(1-p)^{n-m}$$

► If we don't care about the order of the successes, then there are <sub>n</sub>C<sub>m</sub> ways to get m successes in m trials. Therefore,

$$p(m|n,p) = \frac{n!}{m!(n-m)!}p^m(1-p)^{n-m}$$

## **Binomial Distribution**

#### The binomial PDF is a discrete distribution:



Note how the binomial looks increasingly Gaussian as  $n \rightarrow$  large.

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# Binomial Distribution Mean

The mean of the binomial distribution is

$$\begin{split} \langle m \rangle &= \sum_{m=0}^{n} m \cdot \frac{n!}{m!(n-m)!} p^{m} (1-p)^{n-m} \\ &= np \sum_{m=1}^{n} \frac{(n-1)!}{(m-1)!(n-m)!} p^{m-1} (1-p)^{n-m} \\ &= np \sum_{m'=0}^{n'} \frac{n'!}{m'!(n'-m')!} p^{m'} (1-p)^{n'-m'} \\ &= np \end{split}$$

where we simply used the fact that p(m|n, p) is normalized over the sum from m = 0 to n.

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## **Binomial Distribution**

Variance

To find the variance V(m), note that

$$\langle m(m-1) \rangle = \sum_{m=0}^{n} m(m-1) \cdot \frac{n!}{m!(n-m)!} p^{m} (1-p)^{n-m}$$
  
=  $n(n-1)p^2 \sum_{m'=0}^{n'} \frac{n'!}{m'!(n'-m')!} p^{m'} (1-p)^{n'-m'}$   
 $\langle m^2 - m \rangle = n(n-1)p^2$ 

where m' = m - 2, n' = n - 2, and the sum is 1. Therefore,

$$V(m) = \langle m^2 \rangle - \langle m \rangle^2 = \langle m^2 - m \rangle + \langle m \rangle - \langle m \rangle^2$$
$$= n(n-1)p^2 + np - (np)^2$$
$$= np(1-p)$$

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# **Binomial Distribution**

**Detector Efficiencies** 

#### Example

You measure the tracks of cosmic ray particles using a stack of silicon detectors which are 95% efficient. You decide that 3 points are needed to define a track. How efficient is a stack of 3 layers? What about 4, or 5?

$$P(3|p = 0.95, n = 3) = 0.95^{3} = 0.857$$

$$P(3+4|p = 0.95, n = 4) = P(3|...) + P(4|...)$$

$$= \frac{4!}{3!1!} 0.95^{3} 0.05 + 0.95^{4} = 0.986$$

$$P(3+4+5|p = 0.95, n = 5) = P(3|...) + P(4|...) + P(5|...)$$

$$= \frac{5!}{3!2!} 0.95^{3} 0.05^{2} + \frac{4!}{3!1!} 0.95^{3} 0.05 + 0.95^{4}$$

$$= 0.999$$

#### Multinomial Distribution

Generalization of the Binomial Distribution

If instead of two outcomes we have k, we can generalize the binomial distribution to the multinomial distribution:

$$p(m_1, m_2, \ldots, m_k | n, p_1, p_2, \ldots, p_k) = \frac{n!}{\prod_i m_i!} \prod_{i=1}^k p_i^{m_i}$$

where

$$\sum_{i=1}^k p_i = 1, \qquad \sum_{i=1}^k m_i = n$$

• The multinomial is a joint probability distribution over the  $\{m_i\}$ .

#### Example

Example: binned data. If you sample trials from a PDF and bin the results, the predicted counts in each bin will follow a multinomial distribution.

- The Poisson distribution is a limiting case of the binomial distribution (n → ∞, p → 0, ⟨m⟩ → finite).
- It applies when we observe particular outcomes but without knowledge of the number of trials. For example:
  - Number of lightning strikes in a thunderstorm
  - Number of supernova explosions in the Galaxy per century
- Suppose that on average λ events are expected to occur in some interval of length *T*. I.e., the events occur at constant rate *R* such that λ = *RT*.
- If we split the interval up into n sections so that in each section we observe 0 or 1 events, the probability of observing an event in a section is p = λ/n, and the total number of events in the interval follows a binomial distribution:

$$p(m|p = \lambda/n, n) = \frac{n!}{m!(n-m)!}p^m(1-p)^{n-m}$$

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Letting  $n \to \infty$  we find that

$$p(m|p = \lambda/n, n) = \lim_{n \to \infty} \frac{n!}{m!(n-m)!} \left(\frac{\lambda}{n}\right)^m \left(1 - \frac{\lambda}{n}\right)^{n-m}$$

The factorials reduct to a power of n in the large n limit:

$$\lim_{n\to\infty}\frac{n!}{(n-m)!}=\lim_{n\to\infty}n(n-1)(n-2)\dots(n-m+1)\to n^m$$

And we use the definition of the exponential:

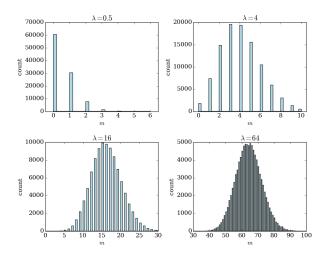
$$\lim_{n\to\infty} \left(1-\frac{\lambda}{n}\right)^{n-m} \to \left(1-\frac{\lambda}{n}\right)^n \to e^{-\lambda}$$

Combining the terms, we get the Poisson distribution:

$$p(m|\lambda) = \frac{e^{-\lambda}\lambda^m}{m!}$$

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#### The Poisson PDF is also discrete distribution:



Note how the Poisson distribution looks increasingly Gaussian as  $\lambda \rightarrow$  large.

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The mean of the Poisson distribution is

$$\begin{split} m \rangle &= \sum_{m=0}^{\infty} m \frac{\lambda^m e^{-\lambda}}{m!} \\ &= \lambda e^{-\lambda} \sum_{m=1}^{\infty} \frac{\lambda^{m-1}}{(m-1)!} \\ &= \lambda e^{-\lambda} \sum_{m'=0}^{\infty} \frac{m' \lambda^{m'}}{m'!} = \lambda e^{-\lambda} e^{\lambda} \\ &= \lambda \end{split}$$

where we used the fact that the sum is the expansion of  $e^{\lambda}$ .

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Variance

To find the variance V(m), we start with

$$\langle m(m-1)\rangle = \sum_{m=0}^{\infty} m(m-1) \cdot \frac{\lambda^m e^{-\lambda}}{m!}$$

As with the binomial distribution, drop the first two terms and set  $m^\prime=m-2$  to get

$$\langle m^2 - m \rangle = \lambda^2 e^{-\lambda} \sum_{m'=0}^{\infty} \frac{\lambda^{m'}}{m'!} = \lambda^2$$

Therefore, the variance is

$$V(m) = \langle m^2 \rangle - \langle m \rangle^2 = \langle m^2 - m \rangle + \langle m \rangle - \langle m \rangle^2$$
$$= \lambda^2 + \lambda - \lambda^2$$
$$= \lambda$$

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#### Poisson Distribution HEP Example

#### Example

Suppose you try to measure a cross-section  $\sigma$  for a process.

- You observe *n* events for an integrated luminosity of  $\mathcal{L}$ .
- For this luminosity, the expected number of events is  $\nu = \sigma \mathcal{L}$ .
- The observed number of events will be Poisson-distributed according to ν.

Our best estimate of  $\nu$  is the number of observed events:  $\hat{\nu} = n$ . For a Poisson distribution, the variance is equal to the mean, so uncertainty on our estimate is given by

$$\hat{\nu} = n \pm \sqrt{n} \implies \hat{\sigma} = \hat{\nu}/\mathcal{L} = (n \pm \sqrt{n})/\mathcal{L}$$

**Note**:  $\sqrt{n}$  is the *estimated* uncertainty of the underlying Poisson mean, not the uncertainty on *n*. There is no "error" on *n*, unless you miscounted!

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Neutrino Counts in Short Time Intervals

#### Example

From Barlow [1]: the number of neutrinos detected in 10-second intervals by the IMB detector on 23 February 1987 was:

No. events	0	1	2	3	4	5	6	7	8	9
No. intervals	1042	860	307	78	15	3	0	0	0	1

The prediction comes from a Poisson distribution with  $\lambda$  obtained by calculating the weighted average

$$\bar{m} = \hat{\lambda} = \sum_{i=0}^{8} w_i c_i / \sum_{i=0}^{8} w_i = \frac{0 \cdot 1042 + 1 \cdot 860 + \dots}{1042 + 860 + \dots} = 0.77$$

Given this mean, the expected Poisson counts are given by

Prediction         1064         823         318         82         16         2         0.3         0.03         0.03	0.0003 0.0003
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# Combining Poisson Variables

#### Sum

The sum of two independent Poisson-distributed variables x and y is itself a Poisson variable z. To see this, first consider the joint probability of xand y:

$$p(x, y|\lambda_x, \lambda_y) = p(x|\lambda_x)p(y|\lambda_y) = \frac{e^{-\lambda_x}\lambda_x^x}{x!} \frac{e^{-\lambda_y}\lambda_y^y}{y!} = \frac{e^{-(\lambda_x+\lambda_y)}\lambda_x^x\lambda_y^y}{x!y!}$$

Now, to find  $p(z|\lambda_z)$ , sum p(x, y) over all (x, y) satisfying x + y = z:

$$p(z|\lambda_z) = \sum_{x=0}^{z} \frac{e^{-(\lambda_x + \lambda_y)} \lambda_x^x \lambda_y^{z-x}}{x!(z-x)!}$$
$$= \frac{e^{-(\lambda_x + \lambda_y)}}{z!} \sum_{x=0}^{z} \frac{z! \lambda_x^x \lambda_y^{z-x}}{x!(z-x)!}$$
$$= \frac{e^{-(\lambda_x + \lambda_y)}}{z!} (\lambda_x + \lambda_y)^z, \text{ by the binomial theorem}$$

## Combining Other Variables

Rules of the road:

- The sum of two Poisson variables is also a Poisson variable, even if the means are different.
- The sum of two Gaussian variables is a Gaussian, even if the means and variances are different.

This is not true for the binomial distribution. In this case:

mean = 
$$np_1 + Np_2$$
, variance =  $np_1(1 - p_1) + Np_2(1 - p_2)$ 

This does not have the general form of the binomial distribution unless  $p_1 = p_2$ . Also note:

- The difference of two Poissons is not Poisson; it follows a Skellam distribution.
- Beware of other false assumptions. E.g., the ratio of two Gaussians is not another Gaussian!

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#### Gaussian Distribution

> You are already familiar with the Gaussian PDF:

$$p(x|\mu,\sigma) = rac{1}{\sqrt{2\pi\sigma}} \exp\left(-rac{(x-\mu)^2}{2\sigma^2}
ight)$$

- The Gaussian is the limiting case of the Poisson distribution (λ → ∞) and the binomial distribution (n → ∞).
- Rules of thumb:
  - Poisson is a good approximation of binomial if  $n \ge 20$  and  $p \le 0.05$ .
  - Gaussian is a good approximation of Poisson if  $\lambda \geq 20$ .
  - Gaussian is a good approximation of binomial if np(1-p) > 9.
- So basically the Gaussian is usually "safe" for large numbers, but beware of using it in the wrong situation.
- The Gaussian has smaller tails than many other distributions and misusing it can cause you to overestimate the significance of rare events.

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## Central Limit Theorem

- Why is the Gaussian so important? Because of the Central Limit Theorem.
- Theorem: the sum of n independent continous random variables x<sub>i</sub> with means μ<sub>i</sub> and variances σ<sup>2</sup><sub>i</sub> becomes a Gaussian with mean and variance

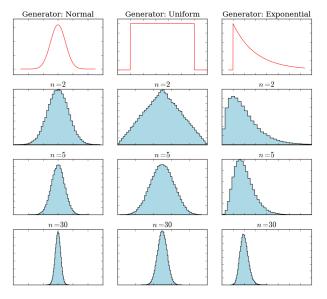
$$\mu = \sum_{i=1}^{n} \mu_i \qquad \sigma^2 = \sum_{i=1}^{n} \sigma_i^2$$

in the limit  $n \to \infty$ .

- See Cowan [2] for a proof based on characteristic functions
- Generally, this is true independent of the individual forms of the PDFs of the x<sub>i</sub> (see next slide).
- Since it is common for many measurements to add together in experiment, the Central Limit Theorem justifies the use of the Gaussian in many cases.

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## Central Limit Theorem



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## Multidimensional Gaussian

▶ The *k*-dimensional generalization of the Gaussian is

$$p(\mathbf{x}|\boldsymbol{\mu}, \boldsymbol{\Sigma}) = rac{1}{\sqrt{(2\pi)^k |\boldsymbol{\Sigma}|}} \exp\left(-rac{1}{2}(\mathbf{x}-\boldsymbol{\mu})^\top \boldsymbol{\Sigma}^{-1}(\mathbf{x}-\boldsymbol{\mu})
ight)$$

- In this expression,  $\mathbf{x} = (x_1, x_2, \dots, x_k)$  is a vector with mean  $\boldsymbol{\mu} = (\mu_1, \mu_2, \dots, \mu_k)$ .
- ➤ is the covariance matrix of the Gaussian. Its diagonal elements are the variances of the x<sub>i</sub>, and its off-diagonal elements are the covariances cov (x<sub>i</sub>, x<sub>j</sub>).

#### Example

Binormal distribution: for k = 2,  $\Sigma$  is a 2  $\times$  2 real symmetric matrix:

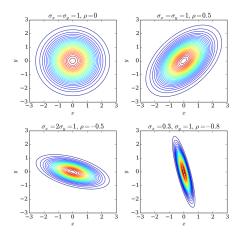
$$\mathbf{x} = \begin{pmatrix} x \\ y \end{pmatrix}, \quad \mathbf{\mu} = \begin{pmatrix} \mu_x \\ \mu_y \end{pmatrix}, \quad \mathbf{\Sigma} = \begin{pmatrix} \sigma_x^2 & \sigma_{xy} \\ \sigma_{xy} & \sigma_y^2 \end{pmatrix}$$

## Change of Variables

- The covariance matrix fully specifies any correlations or anti-correlations between the elements of x.
- If all of the elements of x are independent, then the covariance matrix is diagonal.
- If correlations exist, then there is a unitary matrix *U* that we can identify to diagonalize Σ. I.e.,

 $\Sigma' = U\Sigma U^{\top}.$ 

It is often convenient to change variables to ∑'.



## Uniform Distribution

The uniform (a.k.a. the "top hat" distribution) has a probability which is constant inside some range [a, b] and zero outside:

$$p(x|a,b) = egin{cases} rac{1}{b-a} & a \leq x \leq b, \ 0 & ext{else} \end{cases}$$

• Mean: 
$$\langle x \rangle = (a+b)/2$$

- Variance:  $V(x) = (b a)^2/12$
- Standard deviation:  $\sigma_x = (b a)/\sqrt{12}$
- The uniform distribution is important for two reasons:
  - 1. It is the basis for a large number of pseudorandom number generators.
  - Its constant probability indicates no preferred values inside the range [a, b], making it a popular "objective" prior probability density in Bayesian calculations.

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# $\chi^2$ Distribution

• The  $\chi^2$  distribution of the continuous variable z is

$$p(z|n) = \frac{1}{2^{n/2}\Gamma(n/2)} z^{n/2-1} e^{-z/2},$$

where  $\Gamma$  is the gamma function:

$$\Gamma(x) = \int_0^\infty e^{-t} t^{x-1} dt$$

- ► Note:  $\Gamma(x+1) = x\Gamma(x)$ , and  $\Gamma 1/2 = \sqrt{\pi}$ . For integer *x*,  $\Gamma(x+1) = x!$ .
- Mean: E(x) = n
- Variance: V(x) = 2n
- The simple variance and mean of the  $\chi^2$  distribution make its tail probabilities easy to estimate.

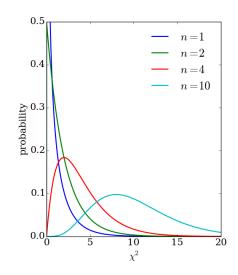
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# $\chi^2$ Distribution

 For n independent Gaussian x<sub>i</sub> with means μ<sub>i</sub> and variances σ<sup>2</sup><sub>i</sub>, the quantity

$$z = \sum_{i=1}^{n} \frac{(x_i - \mu_i)^2}{\sigma_i^2}$$

- follows a  $\chi^2$  with n degrees of freedom.
- Notice that z looks like a least-squares estimator for a fit.
- Physicists often use the tail probability of χ<sup>2</sup> as a measure of goodness of fit.



# Using the $\chi^2$ Distribution

#### Example

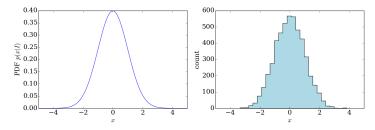
You are shown a fit and told that  $\chi^2$  is 70 for 50 degrees of freedom. Is the fit any good? In other words, how likely is it that  $\chi^2$  could be this large by chance?

Roughly: we expect the mean to be n = 50, and the variance is 2n = 100 with RMS  $\sqrt{100} = 10$ . So this is a  $2\sigma$  effect, which happens  $\sim 2.5\%$  of the time if we approximate using the Gaussian definition of  $\sigma$ .

- If χ<sup>2</sup> ≫ n, then either your model is not a good fit to the data or you badly understimated your uncertainties σ<sub>i</sub>.
- If χ<sup>2</sup> ≪ n, you should also be suspicious. You might have overestimated your uncertainties.

# A Warning about Using the $\chi^2$ Distribution

- Recall the warning given at the end of the last class: the χ<sup>2</sup> statistic z is only asymptotically distributed like a χ<sup>2</sup> distribution if the uncertainties on each x<sub>i</sub> are Gaussian.
- Where this can hurt you: fitting binned data.



- Remember that if your histogram bins are relatively full the uncertainties on the counts in each bin will be Gaussian
- But if the bins are empty or close to empty, the uncertainties in the counts will be Poisson, and z will not follow the χ<sup>2</sup> distribution!

#### Exponential Distribution

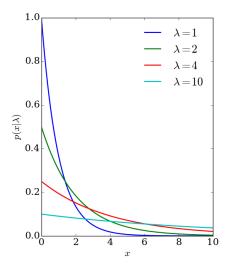
► The exponential PDF is

$$p(x|\lambda) = \frac{1}{\lambda}e^{-x/\lambda}, \quad x \ge 0$$

• Mean:  $E(x) = \lambda$ .

• Variance: 
$$V(x) = \lambda^2$$
, RMS:  $\lambda$ 

- Lack of memory:  $p(t - t_0 | t \ge t_0, \lambda) = p(t | \lambda).$
- Decay time of unstable particle with lifetime  $\lambda \to \tau$
- Lifetime of electrical components, such as lightbulbs

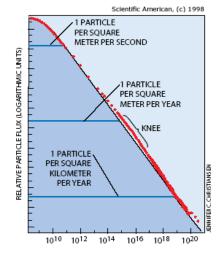


## Power Law (Pareto) Distribution

Power law:

$$p(x|\alpha) = Cx^{-\alpha}$$

- The power law shows up all over physics, and is characteristic of scale invariance, hierarchy, or stochastic generating processes.
- Examples: populations of cities, sizes of lunar impact craters, energies of cosmic rays, sizes of interstellar dust particles, magnitudes of earthquakes, ...



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## Further Reading I

- [1] R.J. Barlow. Statistics: A Guide to the Use of Statistical Methods in the Physical Sciences. New York: Wiley, 1989.
- [2] Glen Cowan. Statistical Data Analysis. New York: Oxford University Press, 1998.

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